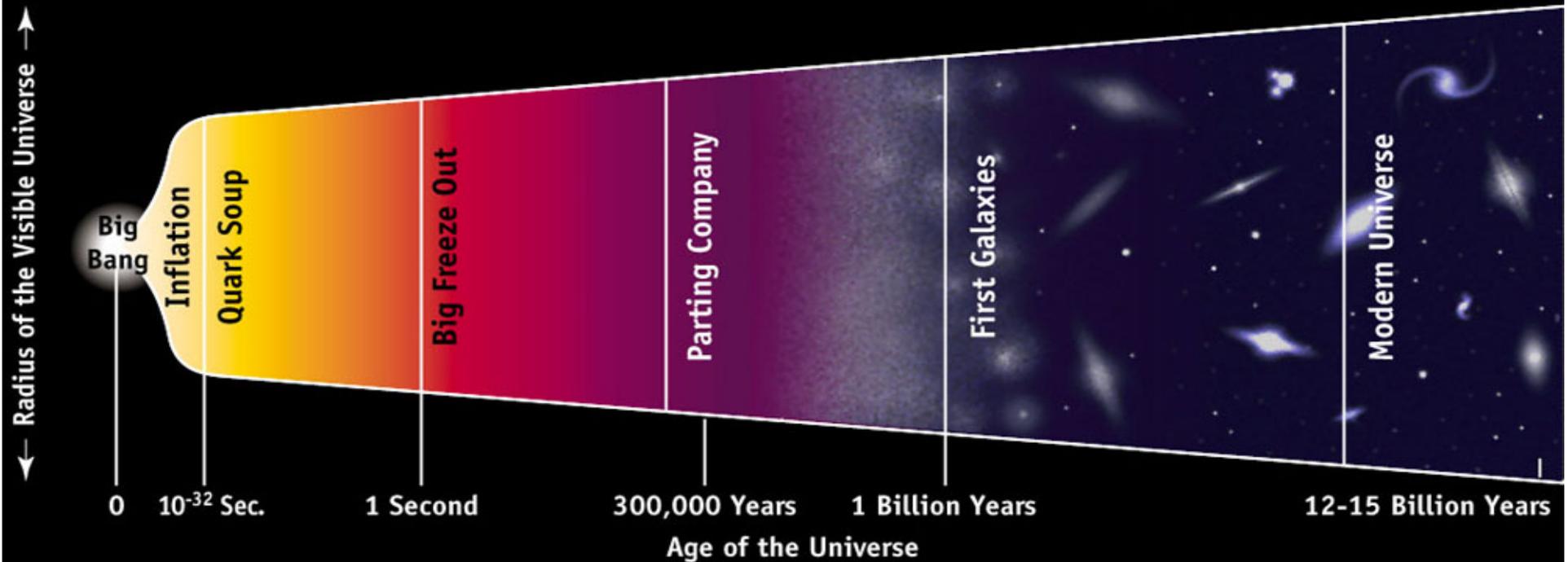
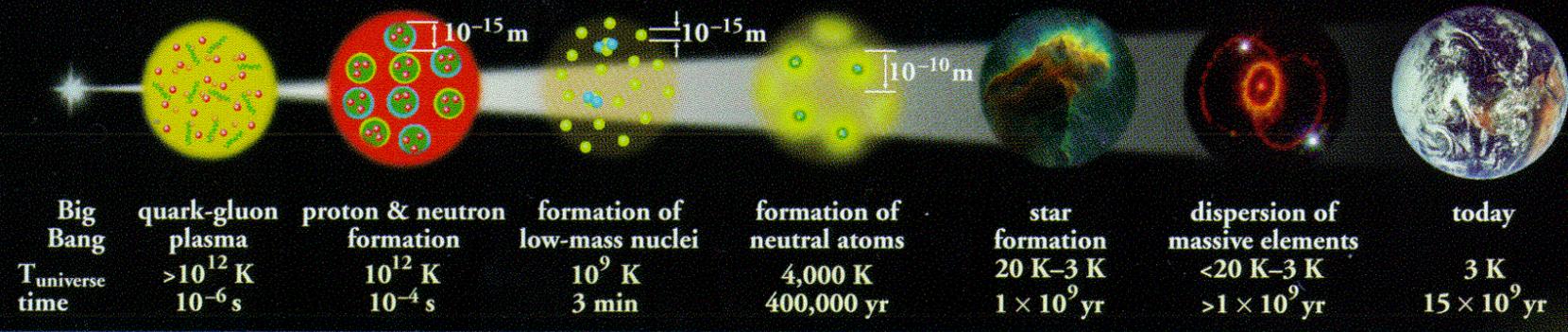
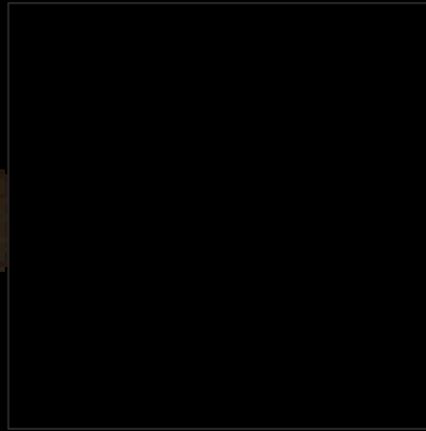
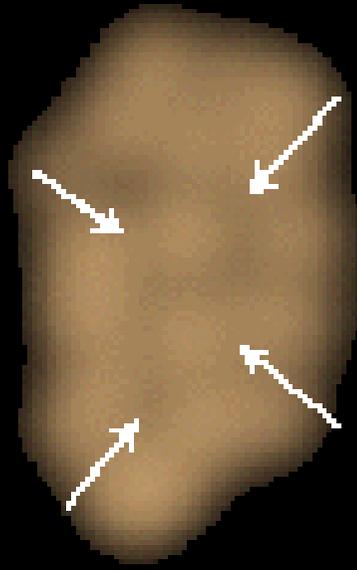


Expansion of the Universe

After the Big Bang, the universe expanded and cooled. At about 10^{-6} second, the universe consisted of a soup of quarks, gluons, electrons, and neutrinos. When the temperature of the Universe, T_{universe} , cooled to about 10^{12} K, this soup coalesced into protons, neutrons, and electrons. As time progressed, some of the protons and neutrons formed deuterium, helium, and lithium nuclei. Still later, electrons combined with protons and these low-mass nuclei to form neutral atoms. Due to gravity, clouds of atoms contracted into stars, where hydrogen and helium fused into more massive chemical elements. Exploding stars (supernovae) form the most massive elements and disperse them into space. Our earth was formed from supernova debris.



nebula



increasing temperature



NGC 6543

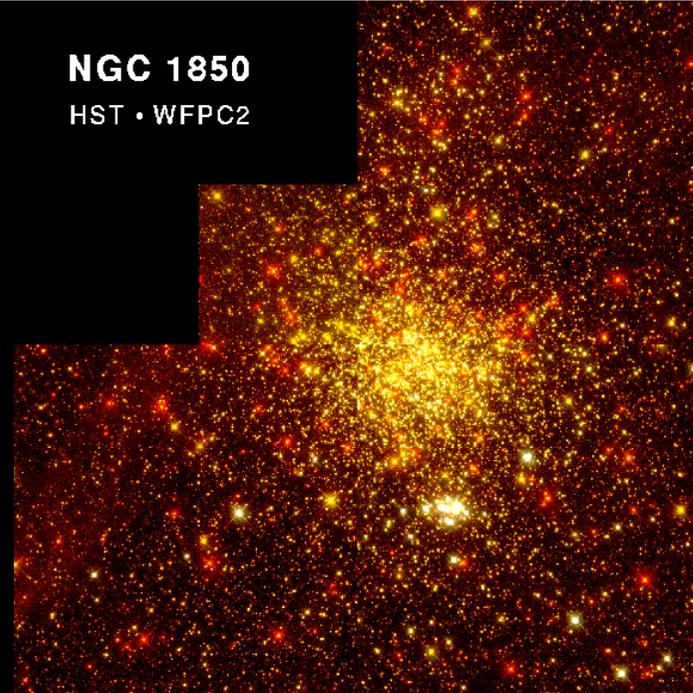
PR95-01a • ST ScI OPO • January 1995 • P. Harrington (U.MD), NASA

HST • WFPC2

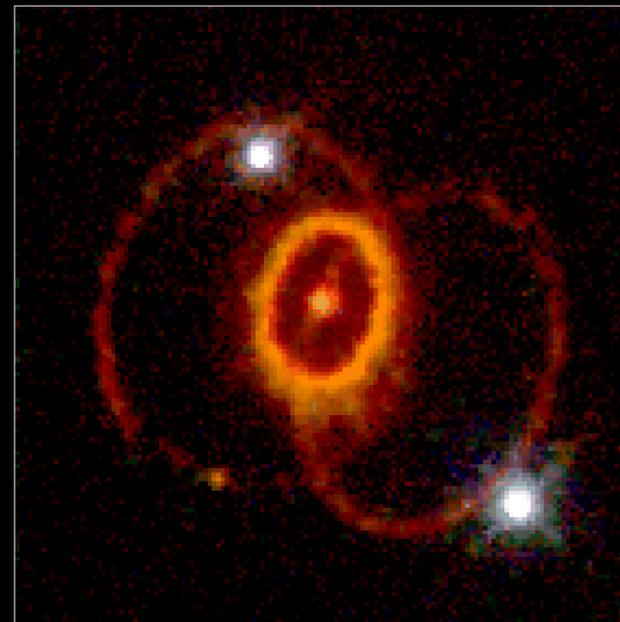
12/13/94 zgl

NGC 1850

HST • WFPC2



Supernova 1987A Rings



Hubble Space Telescope

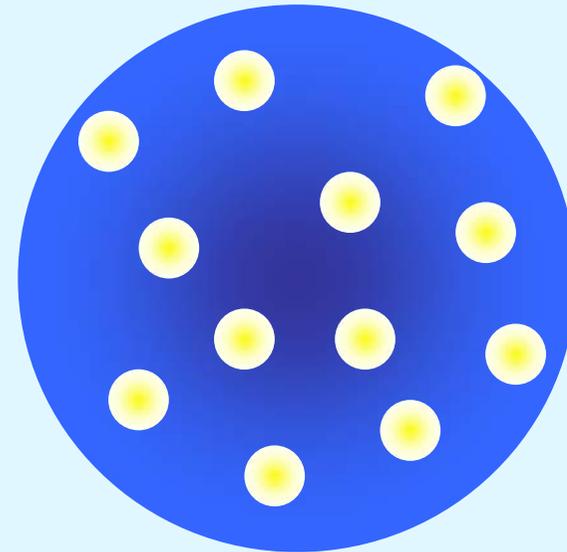


NUCLEO DE THOMSON



ANIMAÇÃO

http://galileoandeinstein.physics.virginia.edu/more_stuff/Applets/rutherford/rutherford2.html

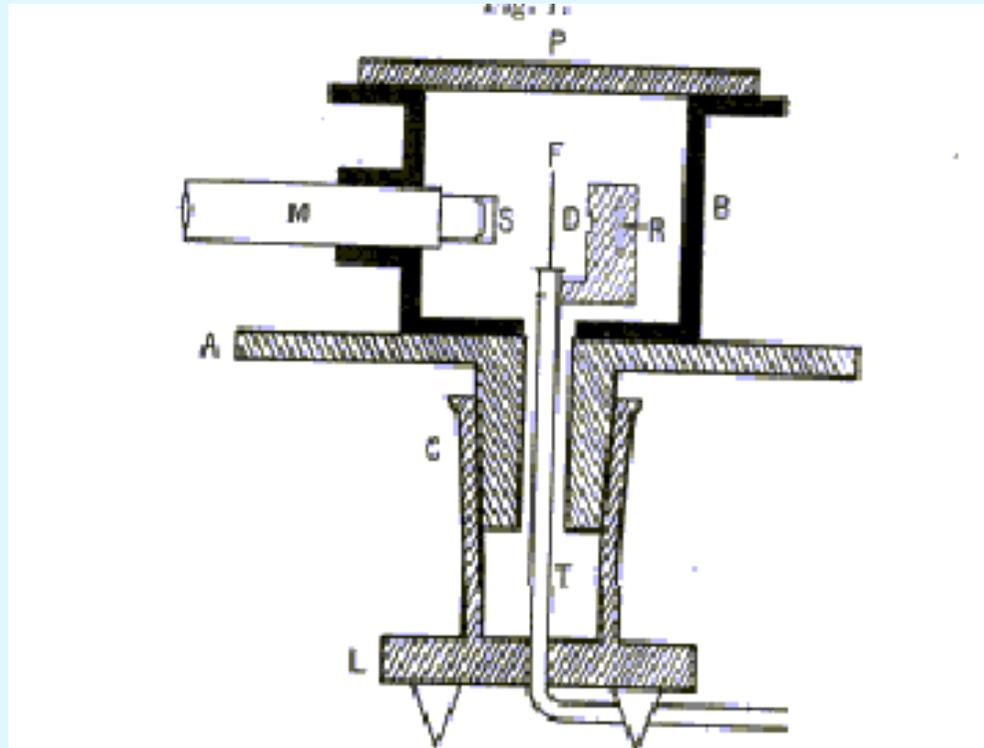


^{24}Mg

NUCLEO DE RUTHERFORD



ANIMAÇÃO



http://galileoandeinstein.physics.virginia.edu/more_stuff/Applets/rutherford/rutherford.html

<http://www.nat.vu.nl/~pwgroen/projects/sdm/applets.htm>

http://www.nat.vu.nl/~pwgroen/sdm/hyper/anim/anim_DI.html

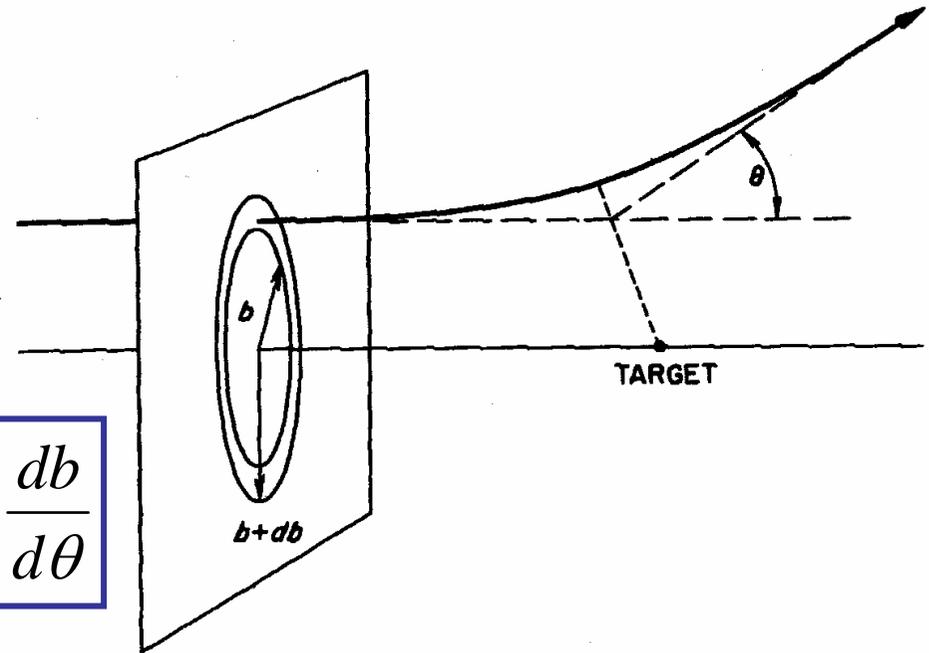
<http://physics.uwstout.edu/physapplets/>

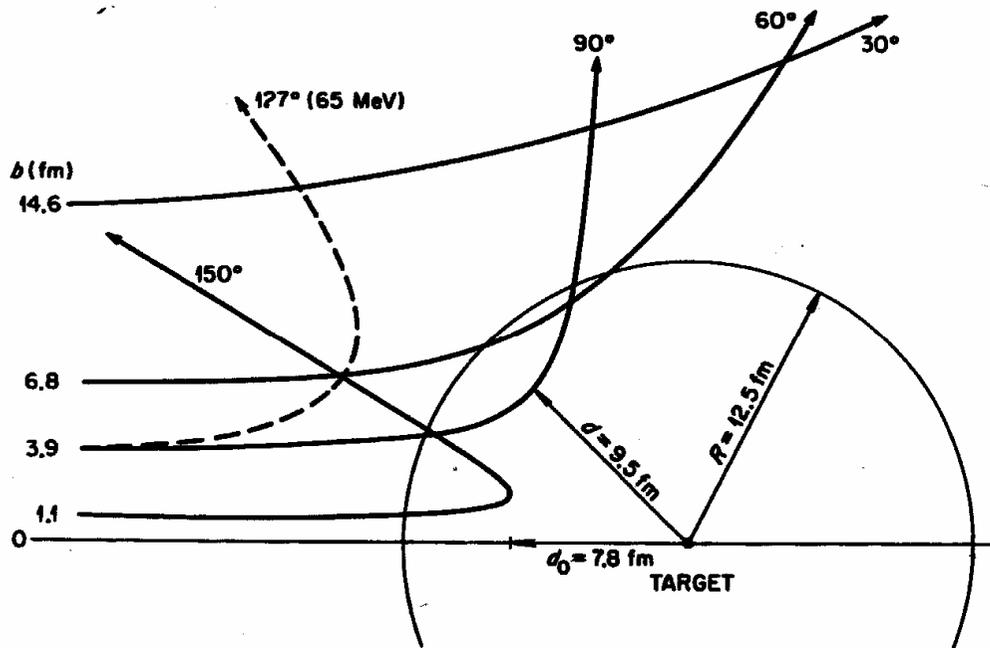
<http://micro.magnet.fsu.edu/electromag/java/rutherford/>

$$\frac{d\sigma}{d\theta} = \frac{d(\pi b^2)}{d\theta} = 2\pi b \frac{db}{d\theta}$$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\theta} \left(\frac{d\theta}{d\Omega} \right)$$

$$\frac{d\sigma}{d\Omega} = \left(\frac{1}{2\pi \sin \theta} \right) \left(2\pi b \frac{db}{d\theta} \right) = \frac{b}{\sin \theta} \frac{db}{d\theta}$$





$$\frac{d\sigma}{d\Omega} = \left[\frac{Z_A Z_a e^2}{2E} \right]^2 \frac{1}{\sin^4 \frac{\theta}{2}}$$

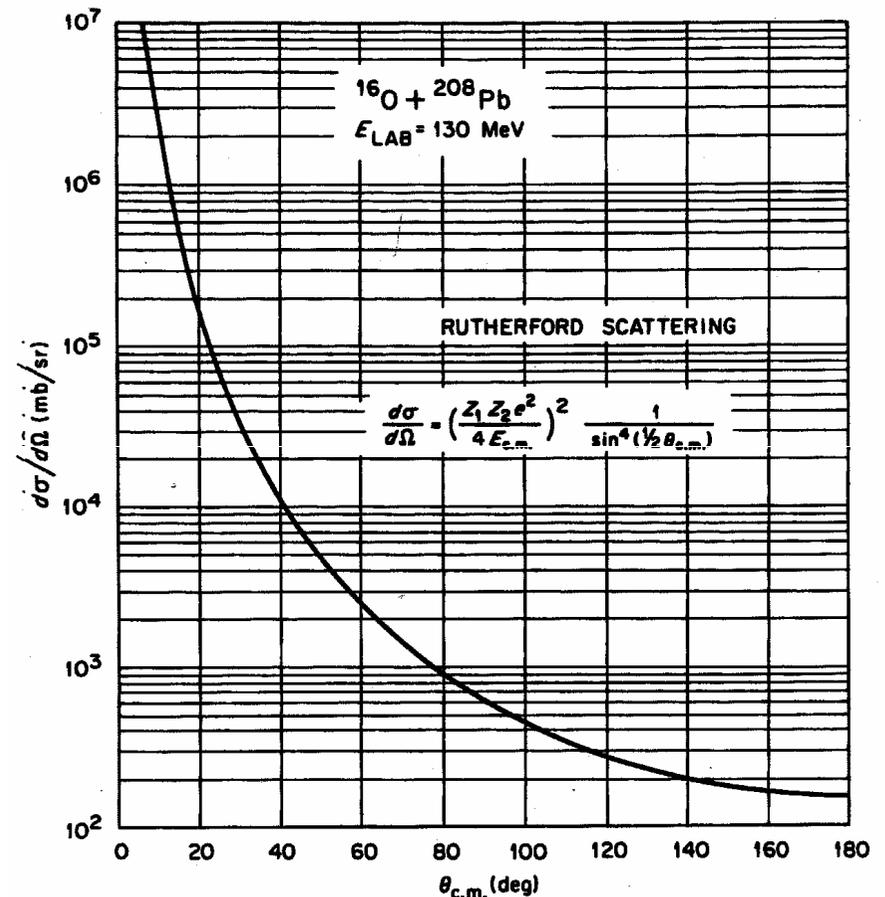
espalhamento:



na energia $E_{^{16}\text{O}} = 130\text{MeV}$

a curva pontilhada

corresponde a $E_{(16\text{O})} = 65\text{MeV}$



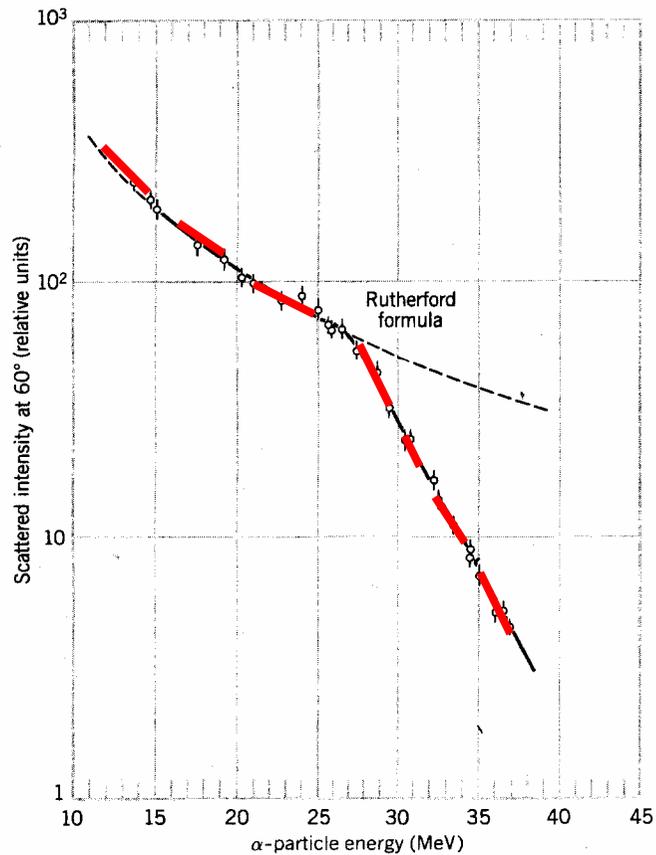


Figure 3.11 The breakdown of the Rutherford scattering formula. When the incident α particle gets close enough to the target Pb nucleus so that they can interact through the nuclear force (in addition to the Coulomb force that acts when they are far apart) the Rutherford formula no longer holds. The point at which this breakdown occurs gives a measure of the size of the nucleus. Adapted from a review of α particle scattering by R. M. Eisberg and C. E. Porter, *Rev. Mod. Phys.* **33**, 190 (1961).

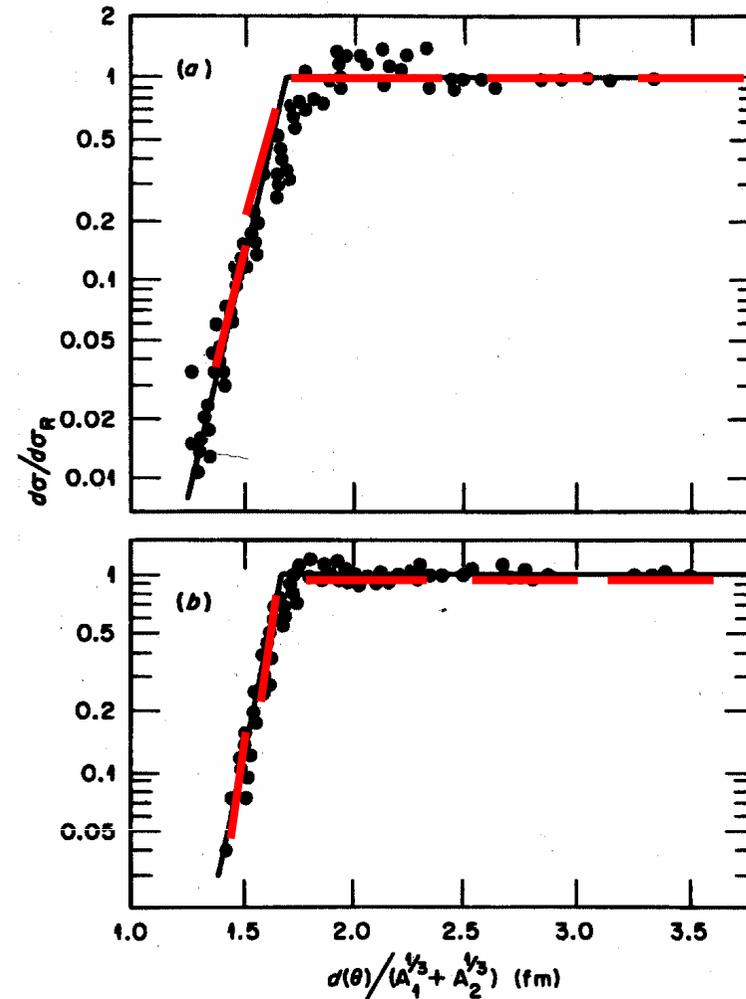
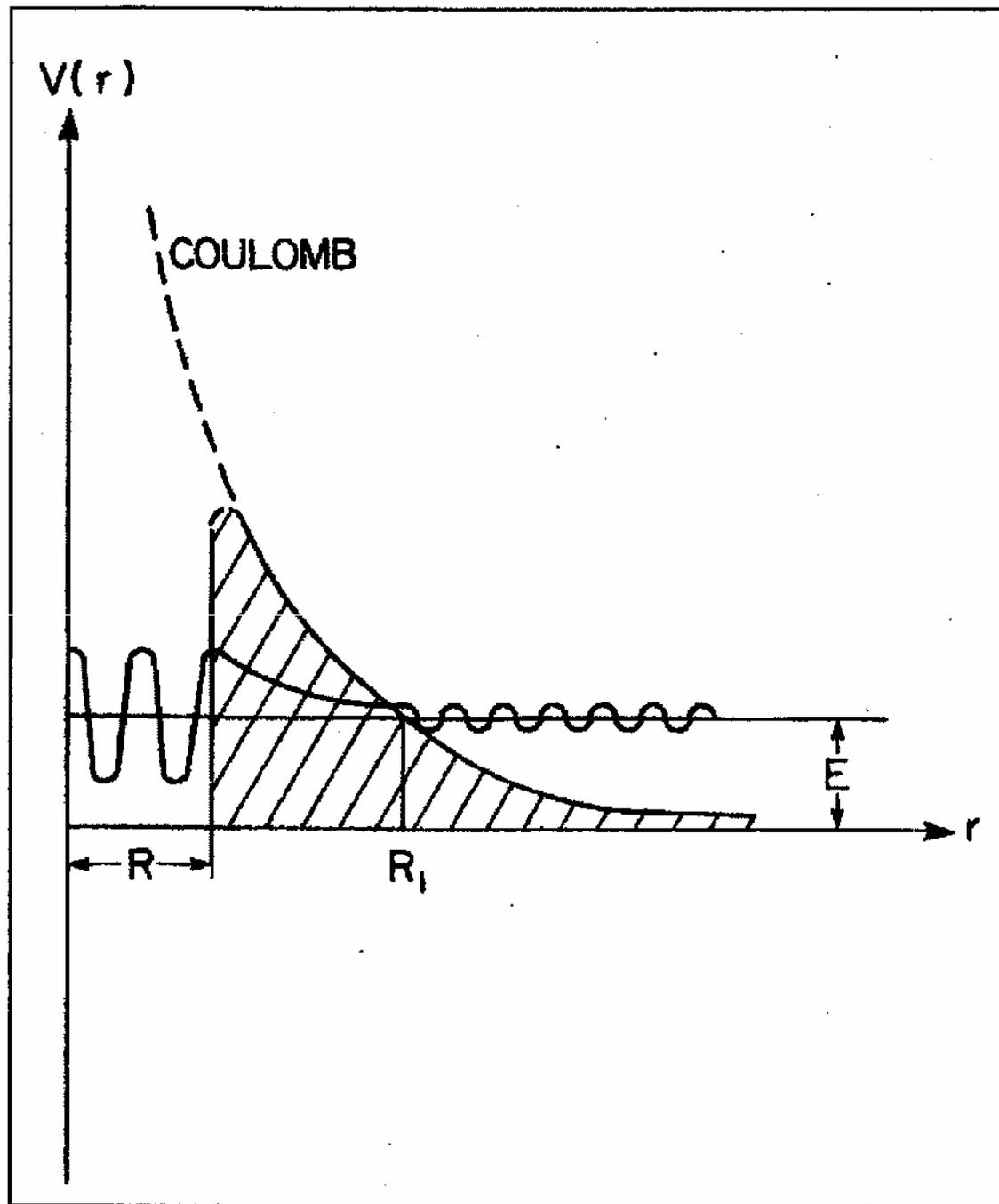
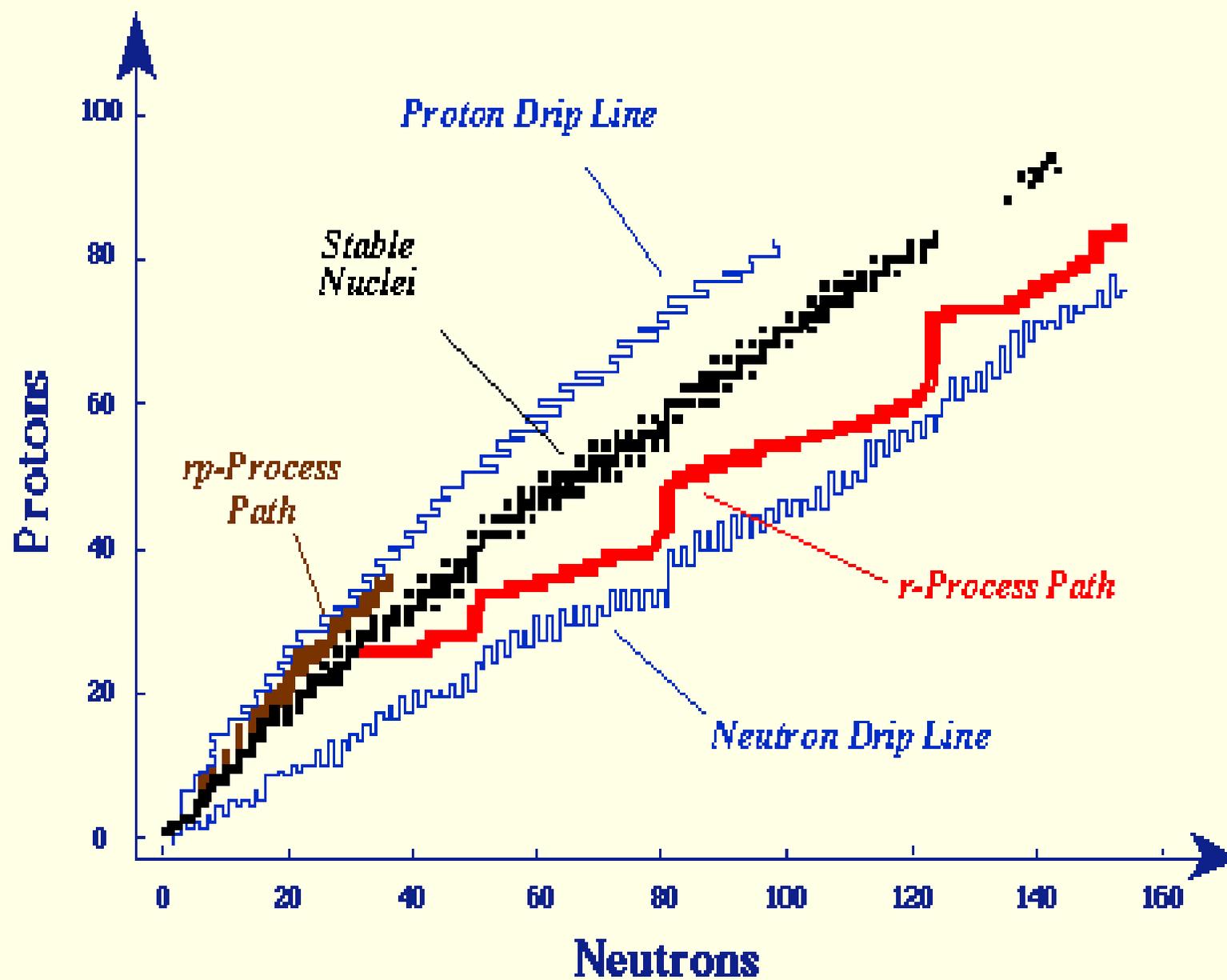
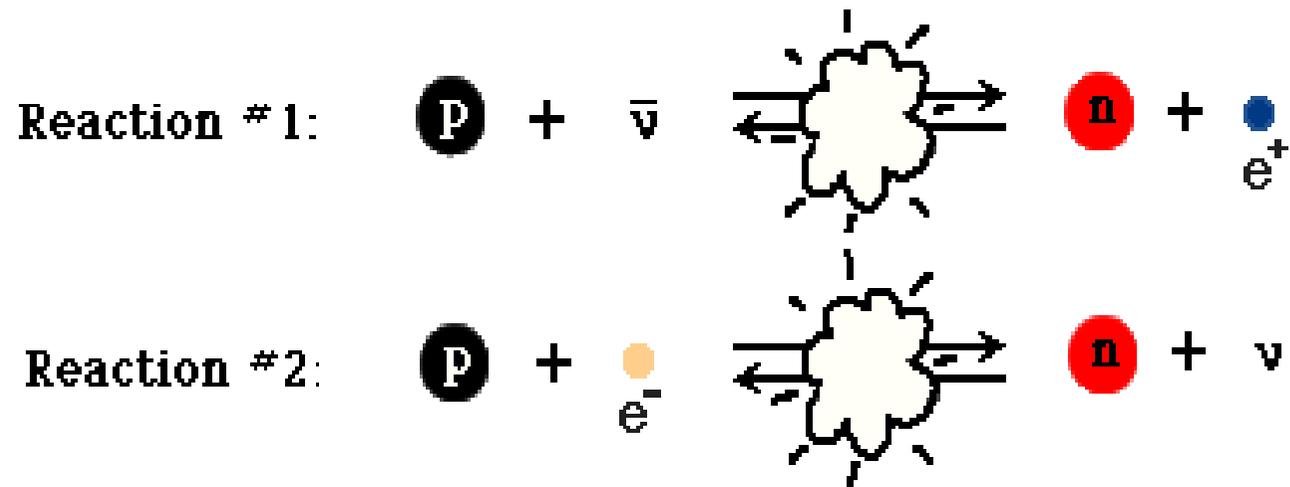


Figure 2.12 The differential cross-sections for O and C ions, in ratio to the Rutherford ones, scattering from various targets plotted against the distance of closest approach d , instead of the scattering angle θ , by using equation 2.20. The distance d has been divided by $(A_1^{1/3} + A_2^{1/3})$, where A_i is the mass number of nucleus i . The measured cross-sections then fall on a universal curve, showing that nuclear radii are approximately proportional to $A^{1/3}$. (a) $^{16}\text{O} + ^{40,48}\text{Ca}$ at 49 MeV, $^{16}\text{O} + ^{40,48}\text{Ca}$, ^{50}Ti , ^{52}Cr , ^{54}Fe , ^{62}Ni at 60 MeV and $^{18}\text{O} + ^{60}\text{Ni}$ at 60 MeV; (b) $^{12}\text{C} + ^{96}\text{Zr}$ at 38 MeV, $^{16}\text{O} + ^{96}\text{Zr}$ at 47, 49 MeV, $^{16}\text{O} + ^{88}\text{Sr}$, ^{93}Zr at 60 MeV and $^{18}\text{O} + ^{90}\text{Zr}$ at 60, 66 MeV. (After Christensen *et al.*, 1973)



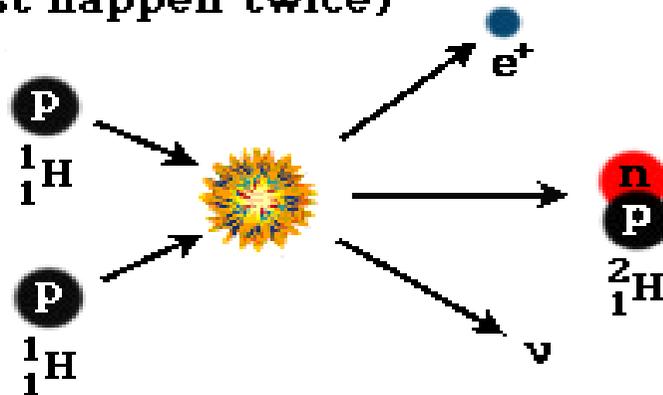


proton/neutron conversions

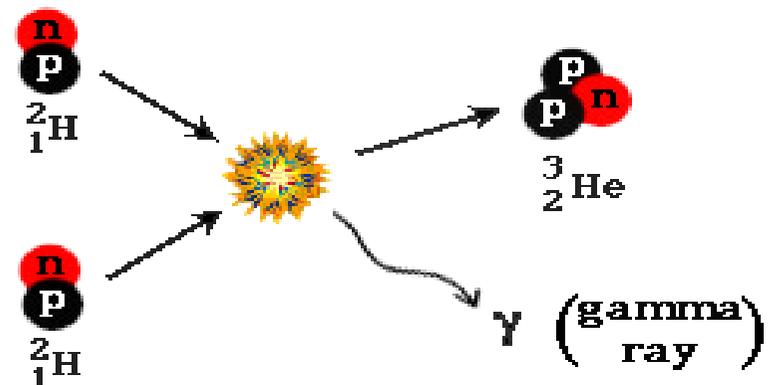


(The double arrows indicate these reactions go both ways.)

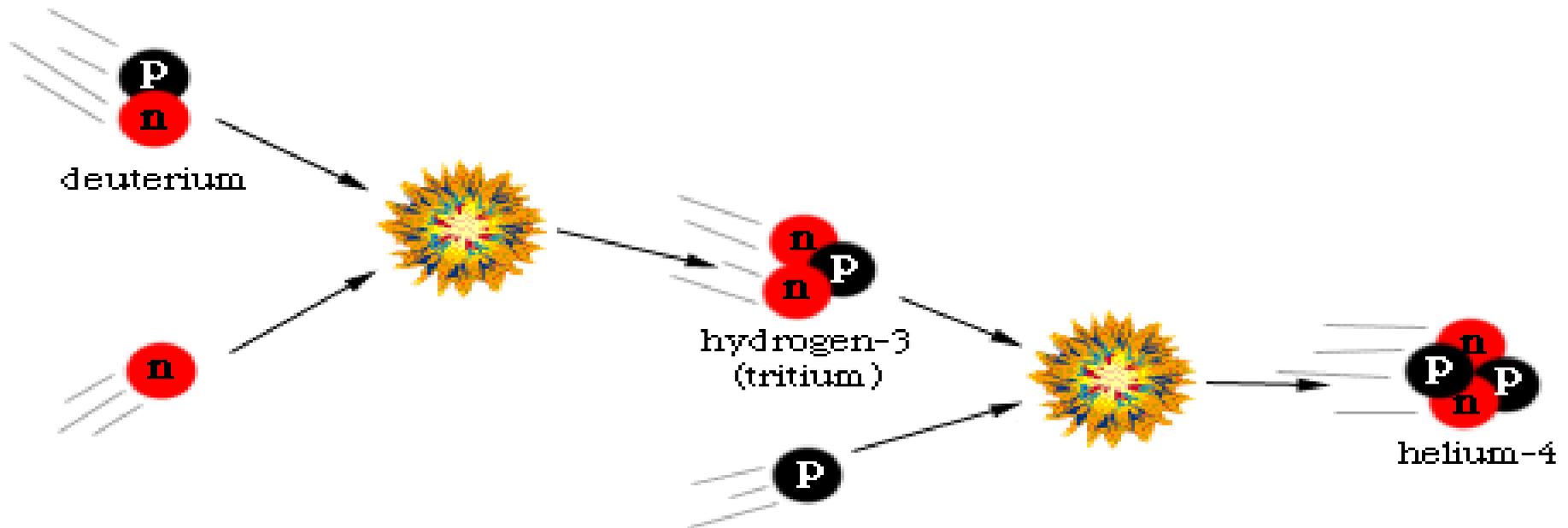
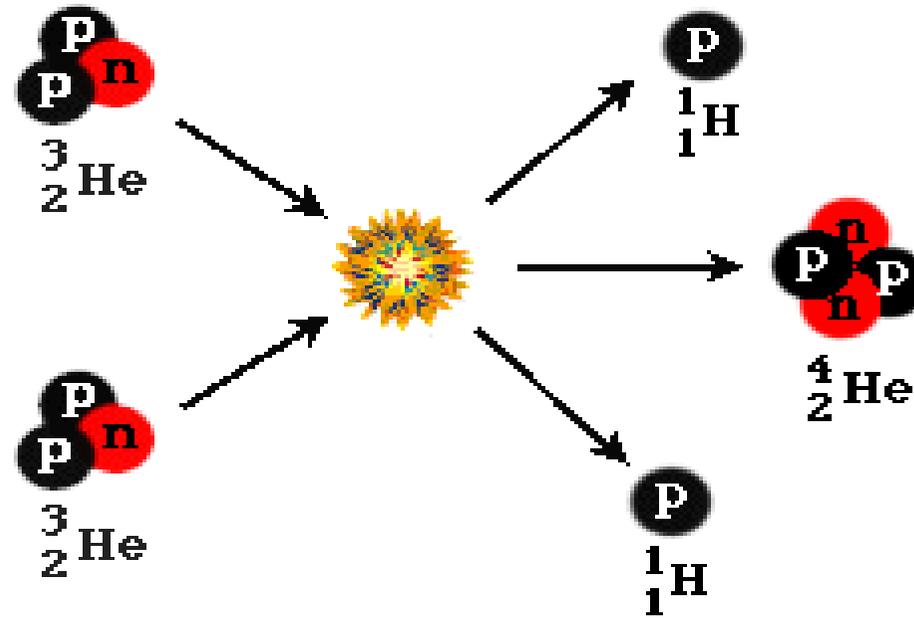
Step 1 (must happen twice)



Step 2 (must happen twice)



Step 3:

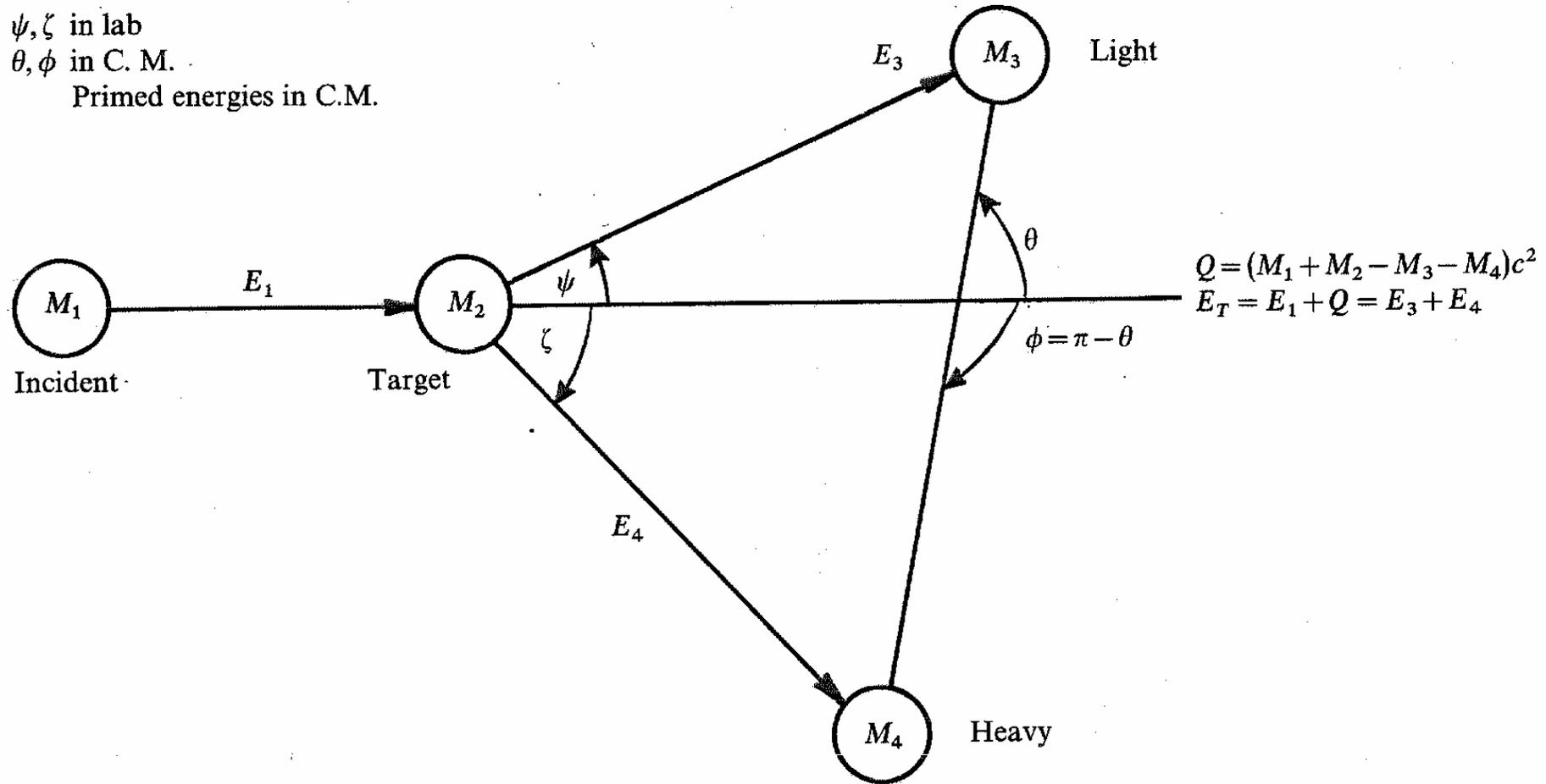


Kinematics of nuclear reactions and scattering (continued)

ψ, ζ in lab

θ, ϕ in C. M.

Primed energies in C.M.



Define:

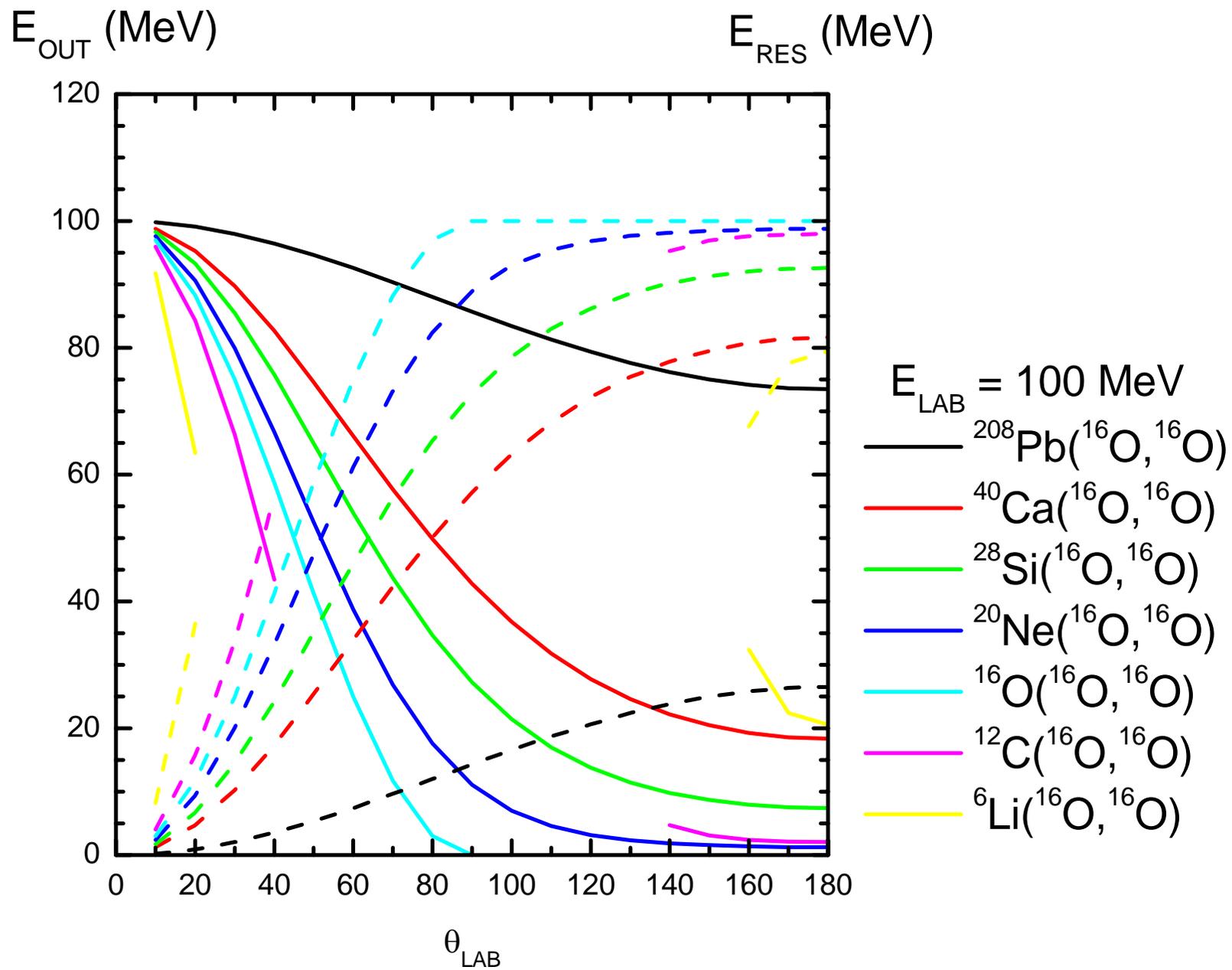
$$A = \frac{M_1 M_4 (E_1 / E_T)}{(M_1 + M_2)(M_3 + M_4)}, \quad C = \frac{M_2 M_3}{(M_1 + M_2)(M_3 + M_4)} \left(1 + \frac{M_1 Q}{M_2 E_T} \right) = \frac{E'_4}{E_T}$$

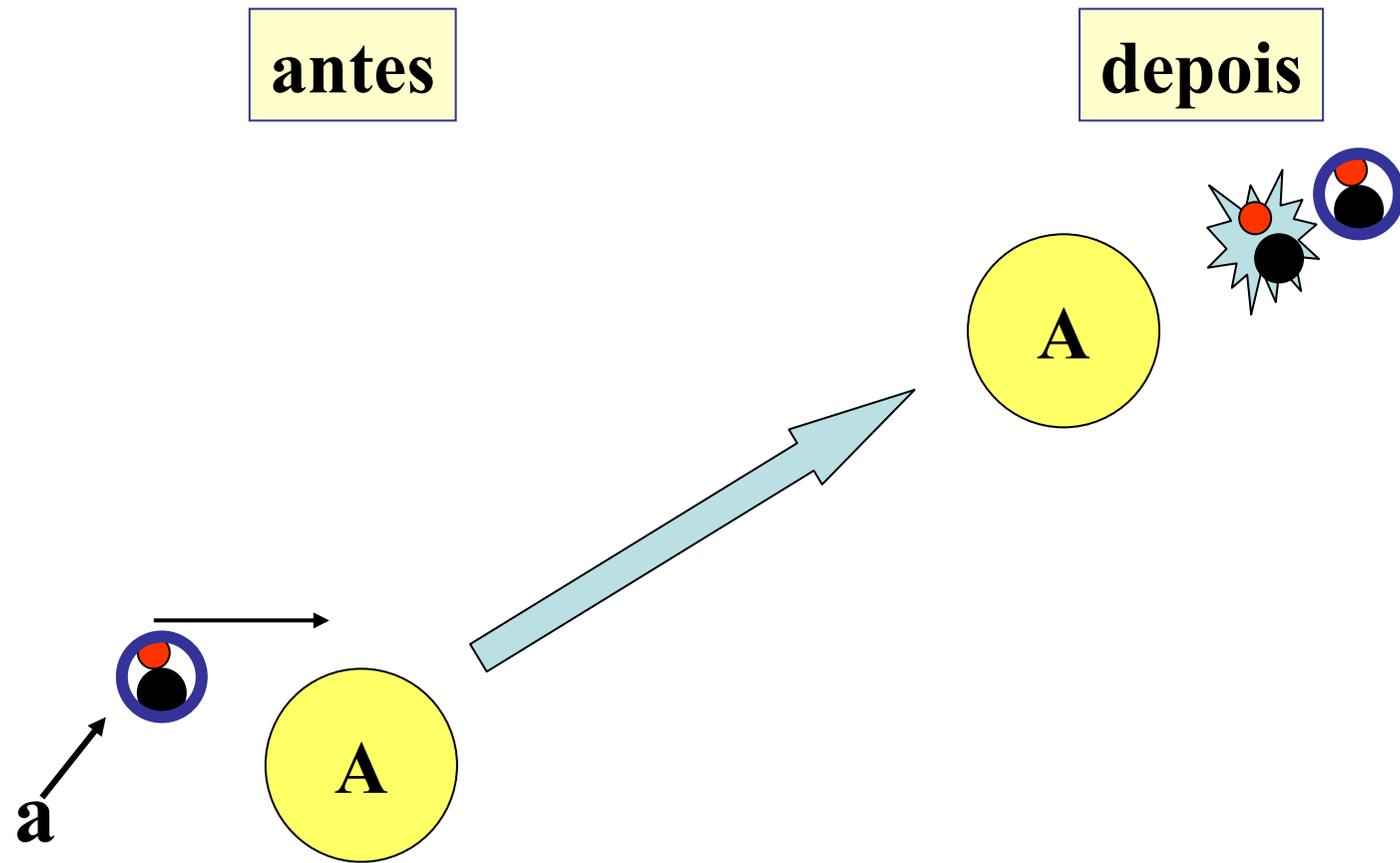
$$B = \frac{M_1 M_3 (E_1 / E_T)}{(M_1 + M_2)(M_3 + M_4)}, \quad D = \frac{M_2 M_4}{(M_1 + M_2)(M_3 + M_4)} \left(1 + \frac{M_1 Q}{M_2 E_T} \right) = \frac{E'_3}{E_T}$$

Note that $A + B + C + D = 1$ and $AC = BD$

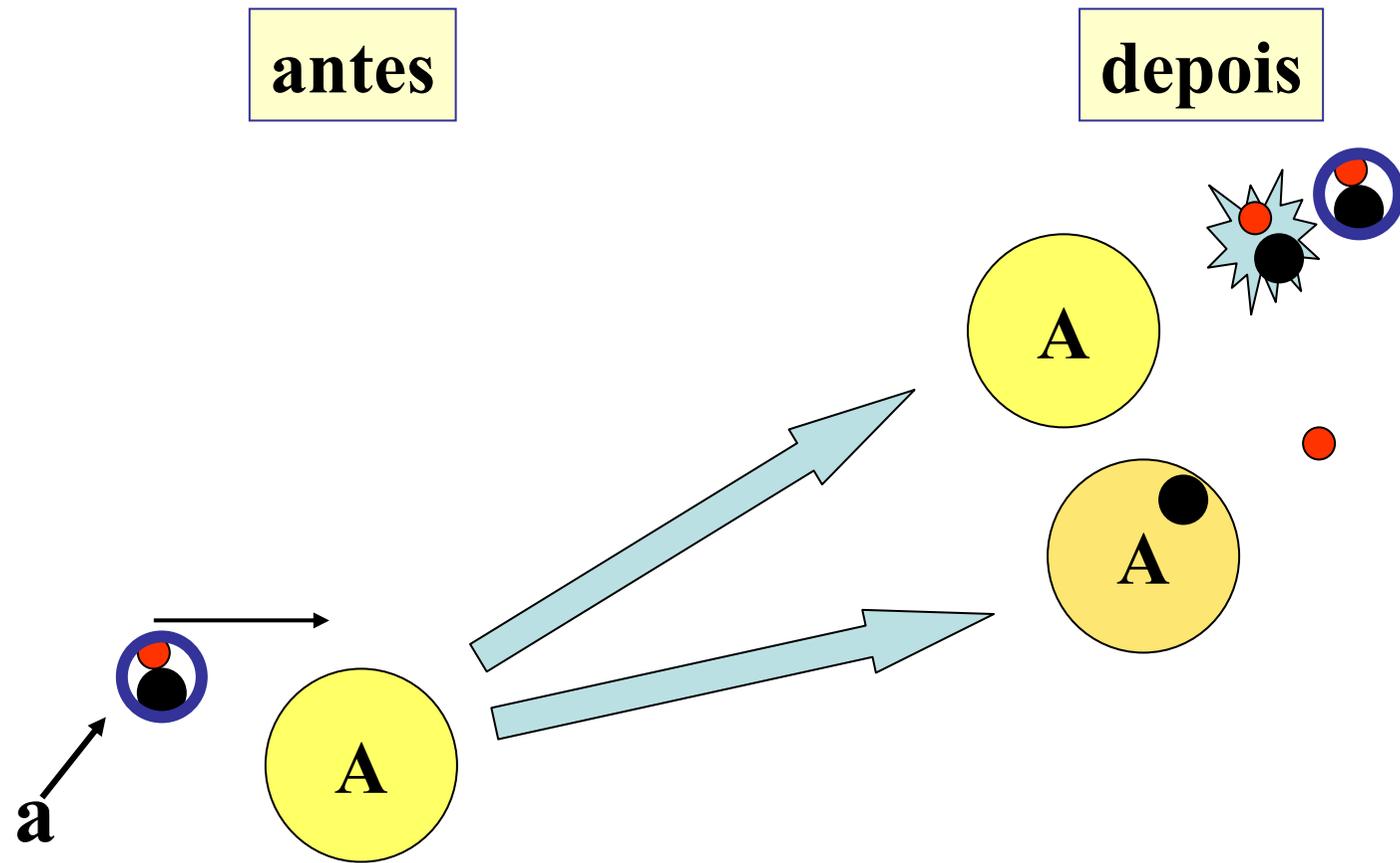
Note that $A + B + C + D = 1$ and $AC = BD$

Lab energy of light product :	$\frac{E_3}{E_T} = B + D + 2(AC)^{\frac{1}{2}} \cos \theta$ $= B[\cos \psi \pm (D/B - \sin^2 \psi)^{\frac{1}{2}}]^2$		Use only plus sign unless $B > D$, in which case $\psi_{\max} = \sin^{-1} (D/B)^{\frac{1}{2}}$
Lab energy of heavy product :	$\frac{E_4}{E_T} = A + C + 2(AC)^{\frac{1}{2}} \cos \phi$ $= A[\cos \zeta \pm (C/A - \sin^2 \zeta)^{\frac{1}{2}}]^2$		Use only plus sign unless $A > C$, in which case $\zeta_{\max} = \sin^{-1} (C/A)^{\frac{1}{2}}$
Lab angle of heavy product :	$\sin \zeta = \left(\frac{M_3 E_3}{M_4 E_4} \right)^{\frac{1}{2}} \sin \psi$	C.M. angle of light product :	$\sin \theta = \left(\frac{E_3/E_T}{D} \right) \sin \psi$
Intensity or solid-angle ratio for light product :	$\frac{\sigma(\theta)}{\sigma(\psi)} = \frac{I(\theta)}{I(\psi)} = \frac{\sin \psi d\psi}{\sin \theta d\theta} = \frac{\sin^2 \psi}{\sin^2 \theta} \cos(\theta - \psi) = \frac{(AC)^{\frac{1}{2}} (D/B - \sin^2 \psi)^{\frac{1}{2}}}{E_3/E_T}$		
Intensity or solid-angle ratio for heavy product :	$\frac{\sigma(\phi)}{\sigma(\zeta)} = \frac{I(\phi)}{I(\zeta)} = \frac{\sin \zeta d\zeta}{\sin \phi d\phi} = \frac{\sin^2 \zeta}{\sin^2 \phi} \cos(\phi - \zeta) = \frac{(AC)^{\frac{1}{2}} (C/A - \sin^2 \zeta)^{\frac{1}{2}}}{E_4/E_T}$		
Intensity or solid-angle ratio for associated particles in the lab system :	$\frac{\sigma(\zeta)}{\sigma(\psi)} = \frac{I(\zeta)}{I(\psi)} = \frac{\sin \psi d\psi}{\sin \zeta d\zeta} = \frac{\sin^2 \psi \cos(\theta - \psi)}{\sin^2 \zeta \cos(\phi - \zeta)}$		

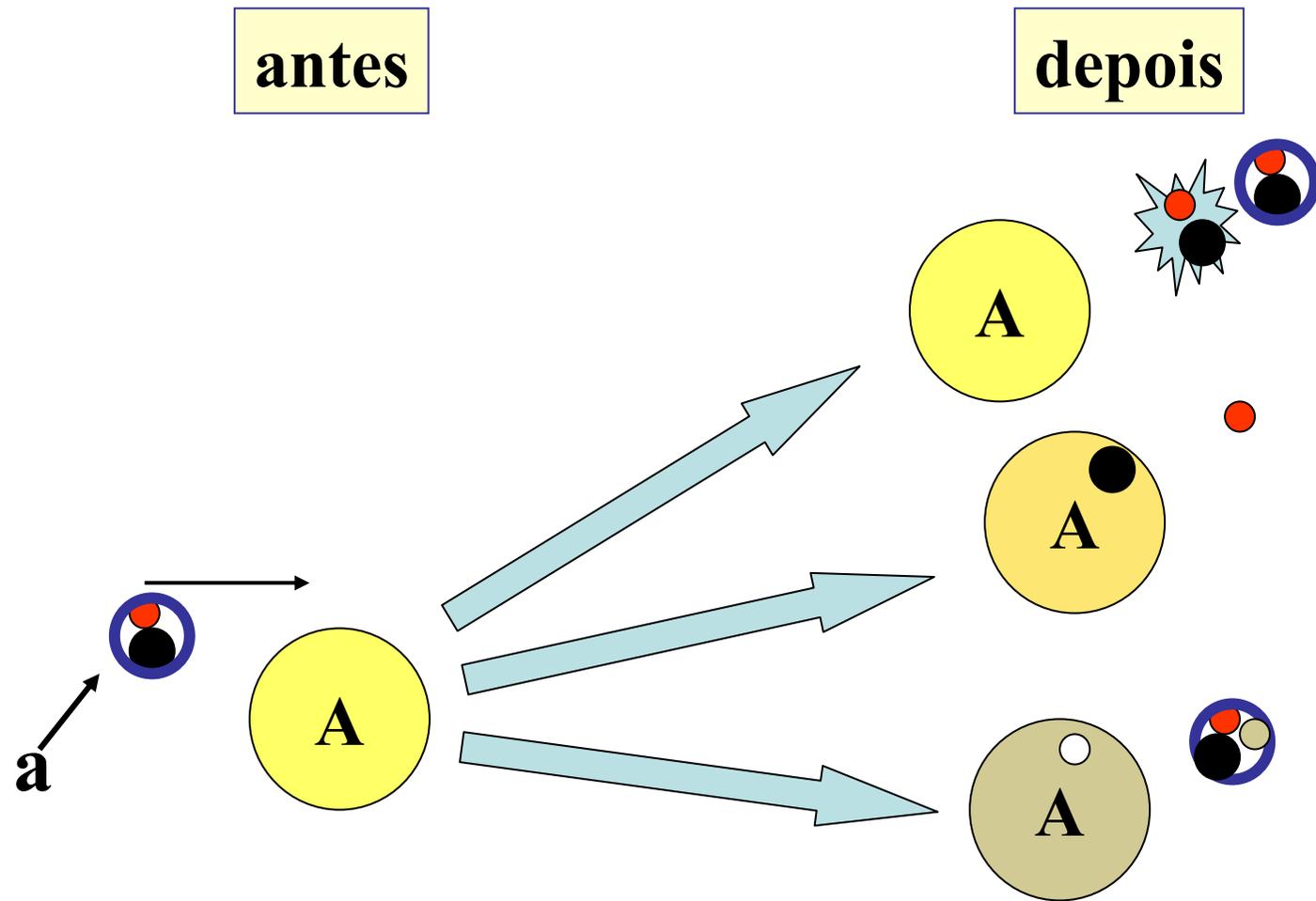




Reações diretas (rápidas)



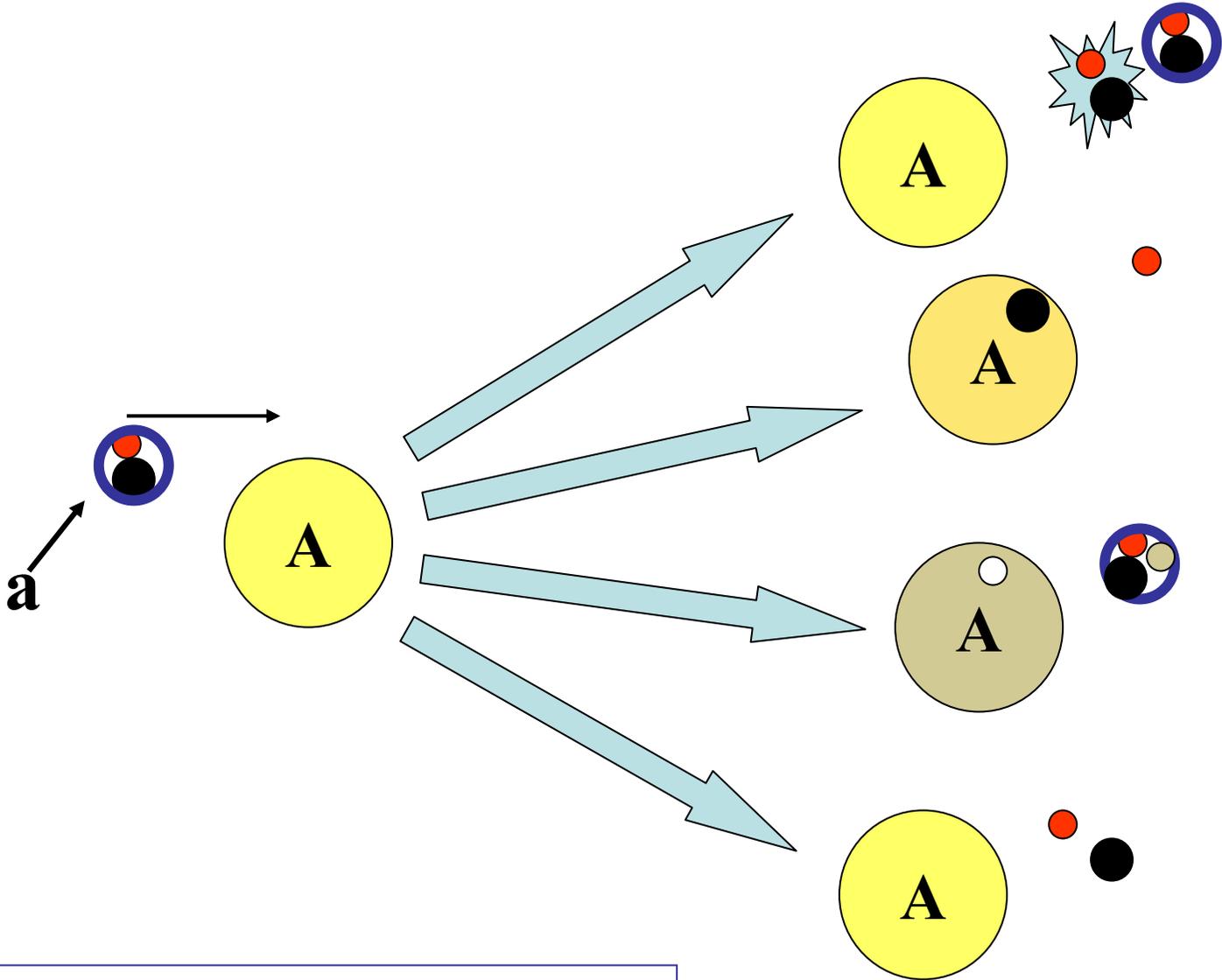
Reações diretas (rapidas)



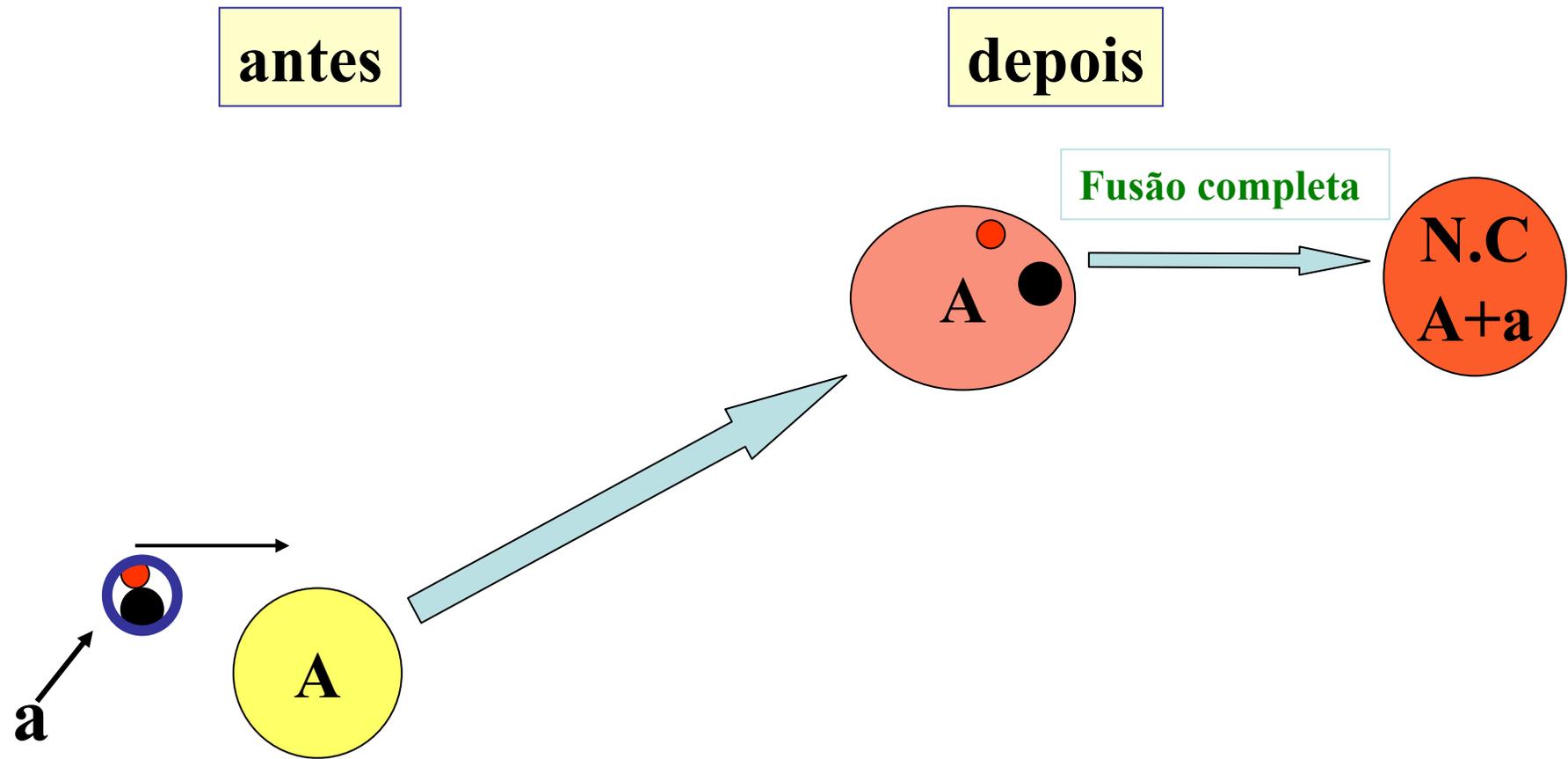
Reações diretas (rapidas)

antes

depois

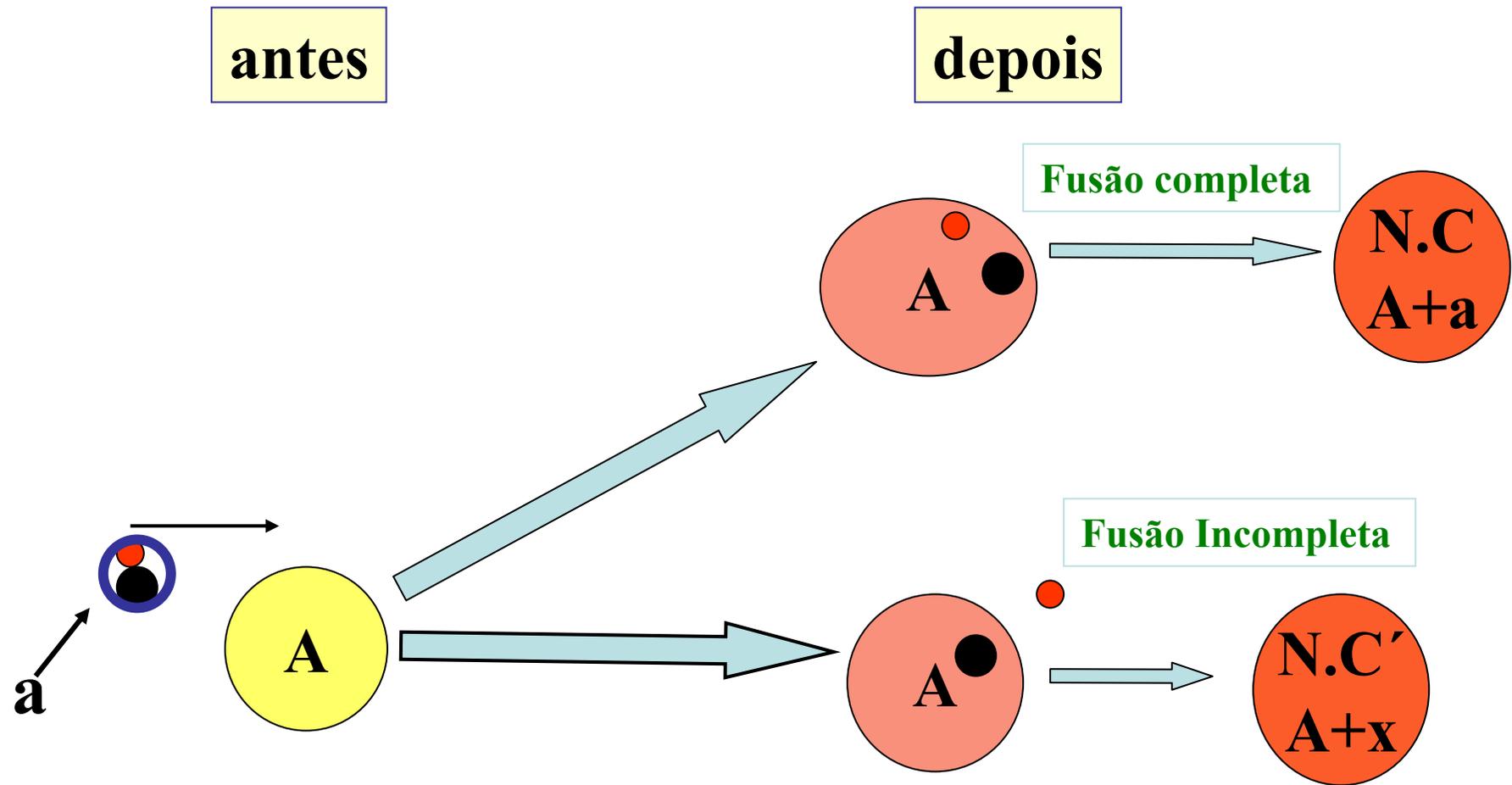


Reações diretas (rapidas)



Processos estatísticos (lentos)

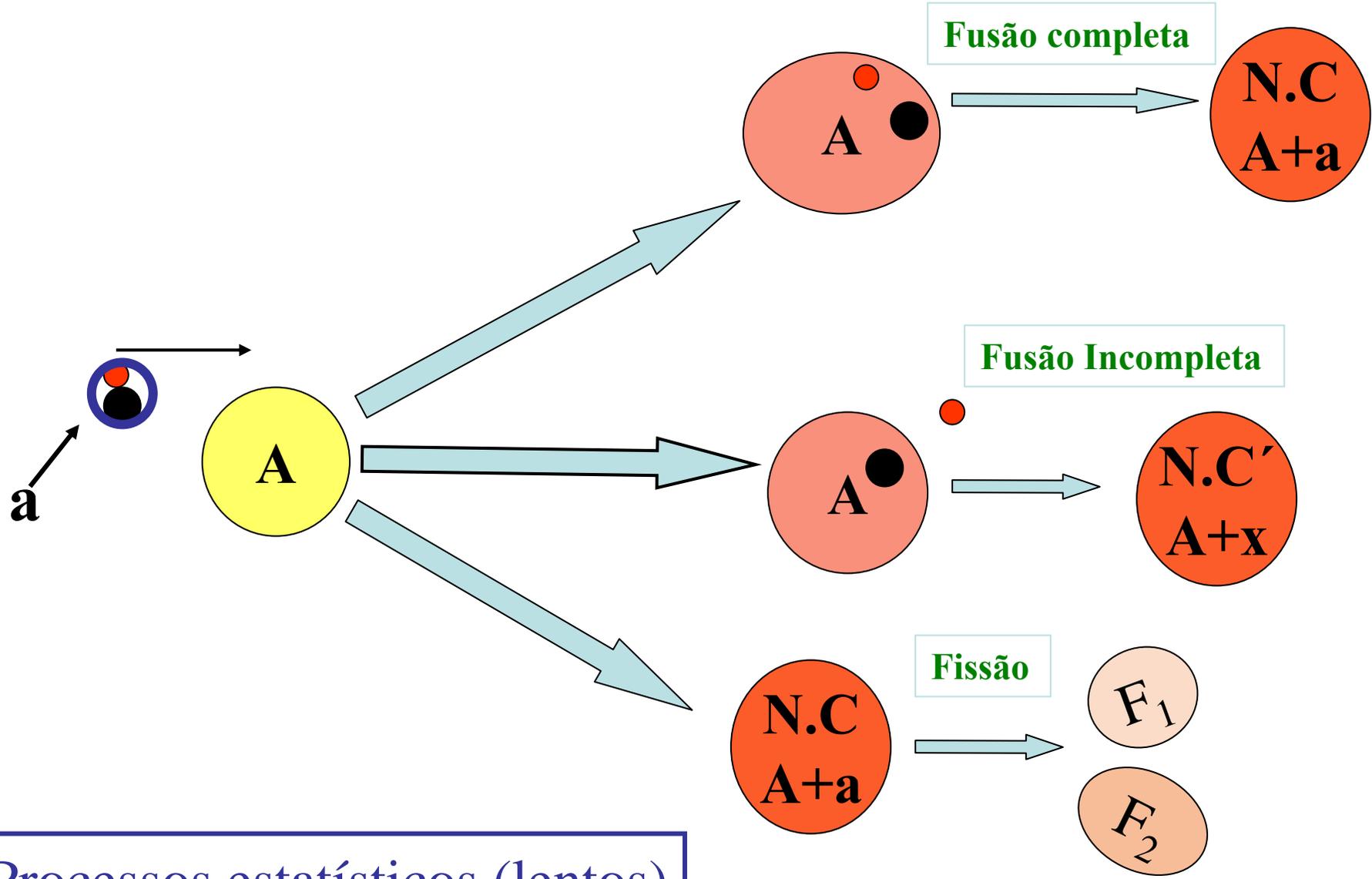
via **Núcleo Composto (N.C.)**



Processos estatísticos (lentos)

antes

depois



Processos estatísticos (lentos)

ESPALHAMENTO ELASTICO

REAÇÕES DIRETAS

ESPALHAMENTO INELASTICO

TRANSFERENCIA DE NUCLEONS

“STRIPPING”

“PICK-UP”

“KNOCK-OUT”

QUEBRA NUCLEAR (“BREAK-UP”)

PRÉ-EQUILIBRIO

NUCLEO COMPOSTO

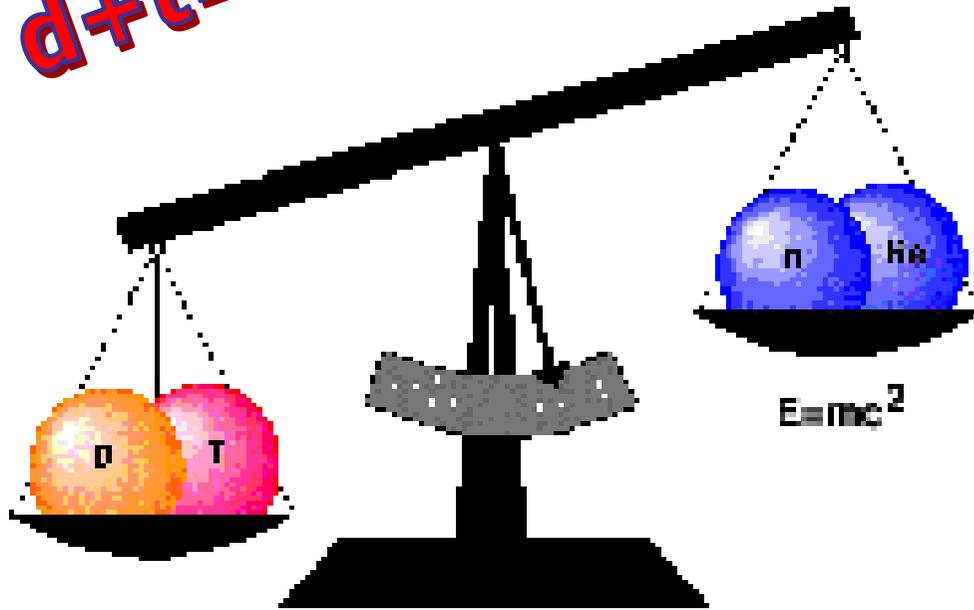
FUSÃO

FUSÃO COMPLETA

FUSÃO INCOMPLETA

FISSÃO

$p+p+n+n \equiv \text{ALPHA}$
 $d+t \equiv n + \text{ALPHA}$



$n = 1.00866 \text{ u.m.a.}$

$p = 1.0079 \text{ u.m.a.}$

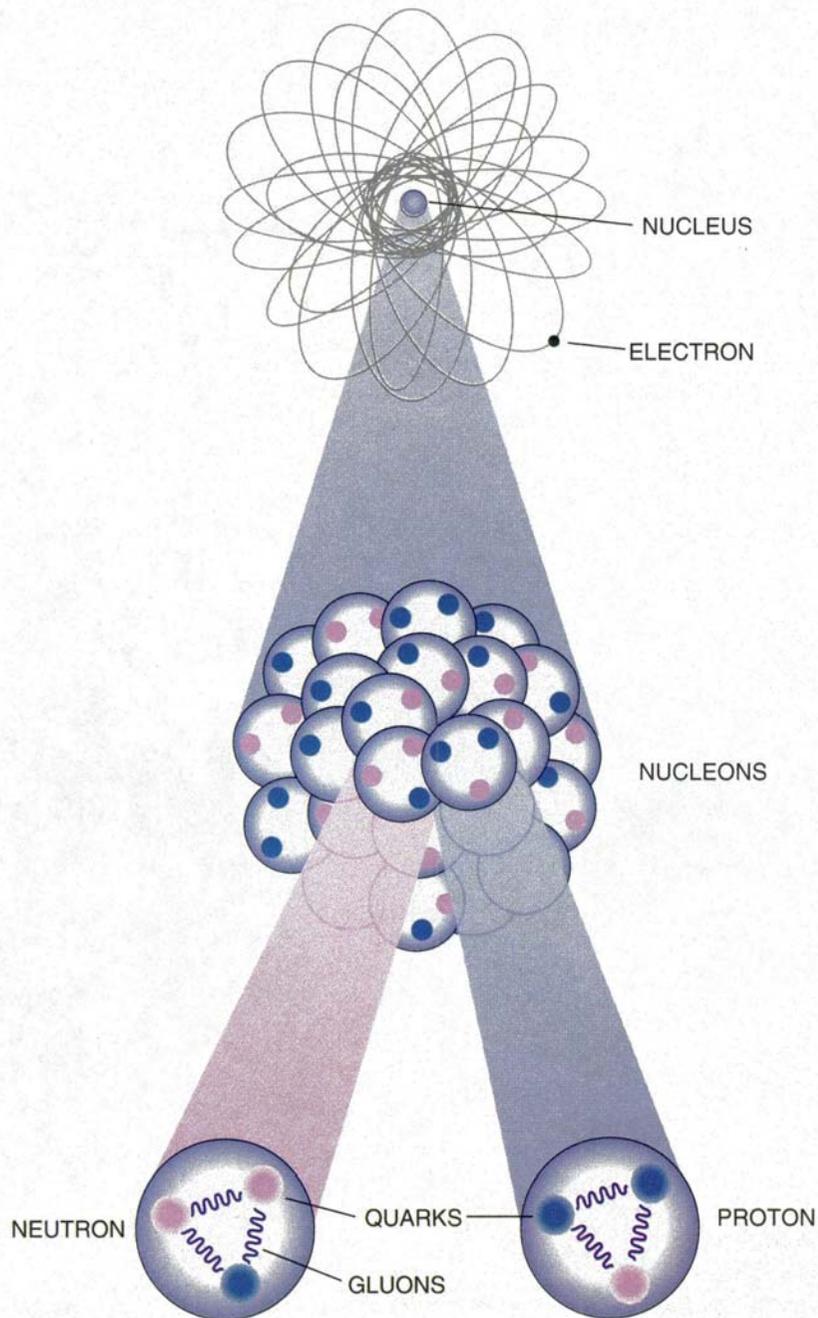
$d = 2.01410 \text{ u.m.a.}$

$t = 3.01860 \text{ u.m.a.}$

${}^4\text{He} = 4.00260 \text{ u.m.a.}$

${}^6\text{Li} = 6.01512 \text{ u.m.a.}$

${}^{12}\text{C} = 0.00000 \text{ u.m.a.}$



$$M_{(\text{proton})} = M_p = 938.27 \text{ MeV}$$

$$M_{(\text{neutron})} = M_n = 939.56 \text{ MeV}$$

$$M_{(\text{eletron})} = M_e = 0.511 \text{ MeV}$$

$$M_{(\text{átomo de Hidrogênio})} = M_H = 938.58 \text{ MeV}$$

$$M_p + M_e = 938.27 + 0.511 = 938.78 \text{ MeV}$$

$$M_H = 938.58 \text{ MeV}$$

$$M_{(16O)} = 14\,899.17 \text{ MeV}$$

$$8M_p + 8M_n = 15\,022.64 \text{ MeV}$$

$$\Delta M = 123.47 \text{ MeV}$$

$$p = (u + u + d) \quad M_{(u)} \approx 3 \text{ MeV}$$

$$M_{(d)} \approx 6 \text{ MeV}$$

$$2 M_{(u)} + M_{(d)} \approx 12 \text{ MeV}$$

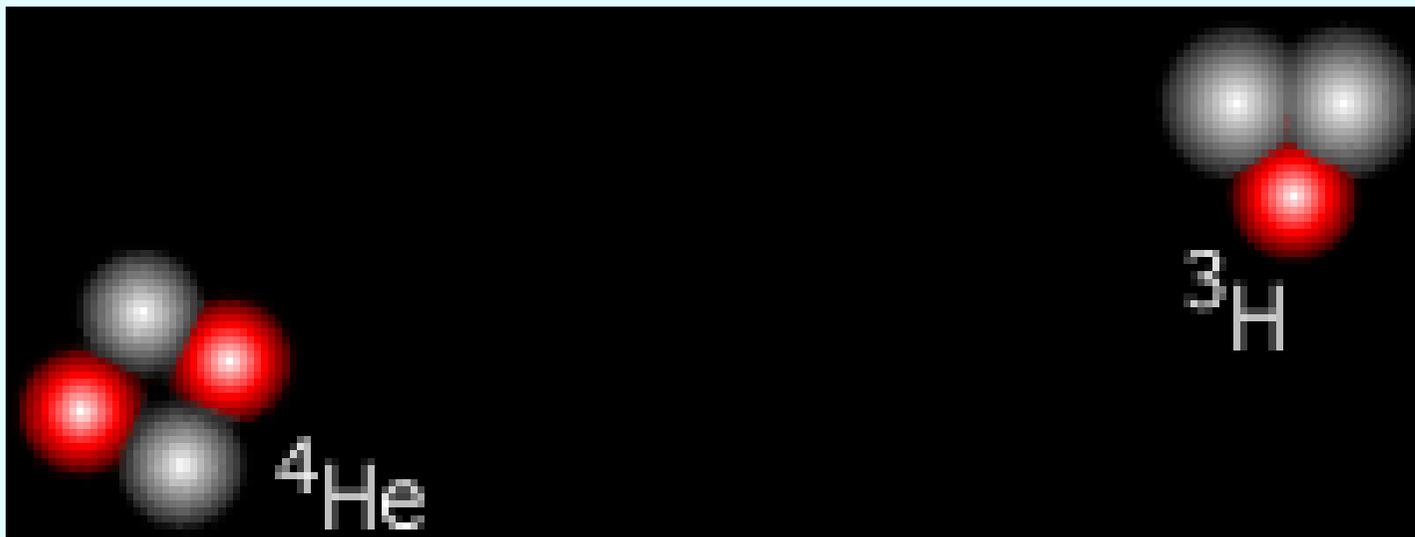
$$M_p = 938.27 \text{ MeV}$$

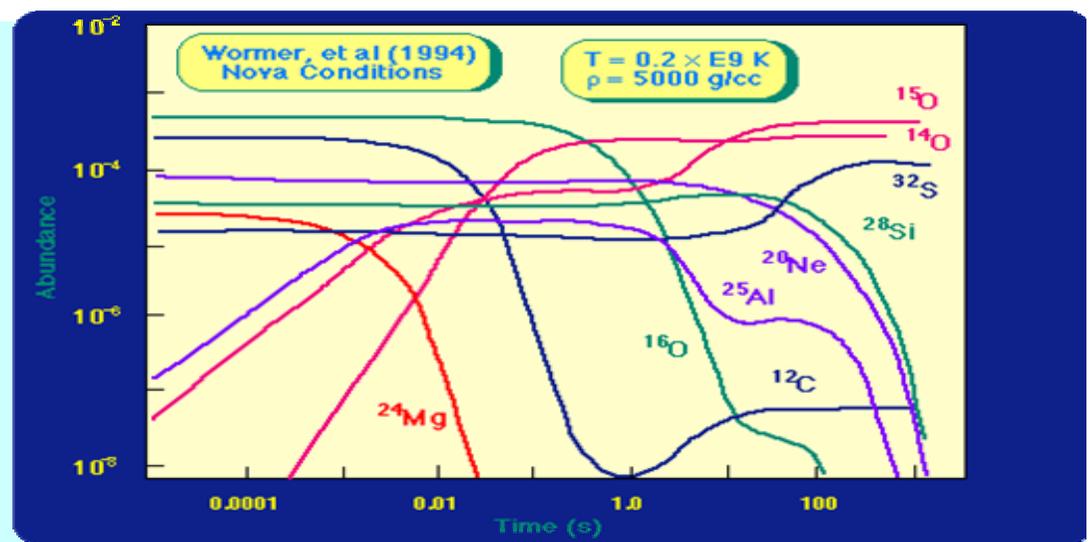
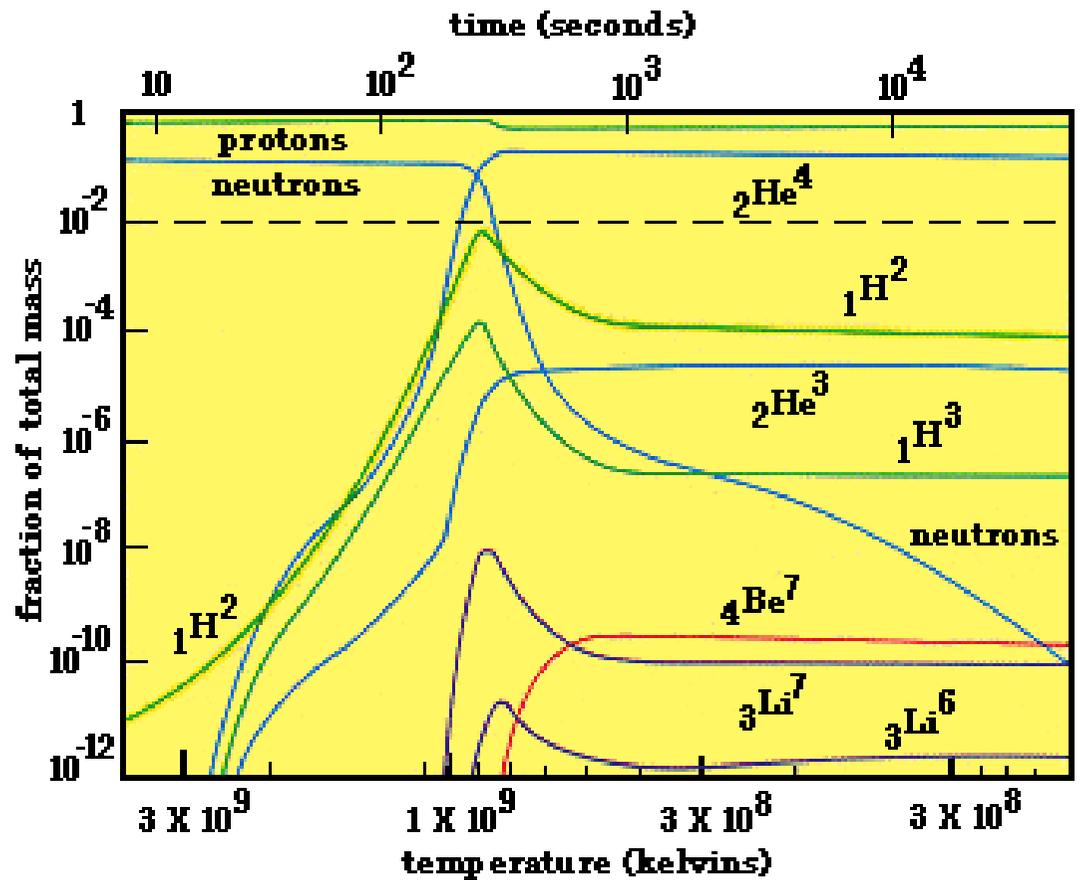
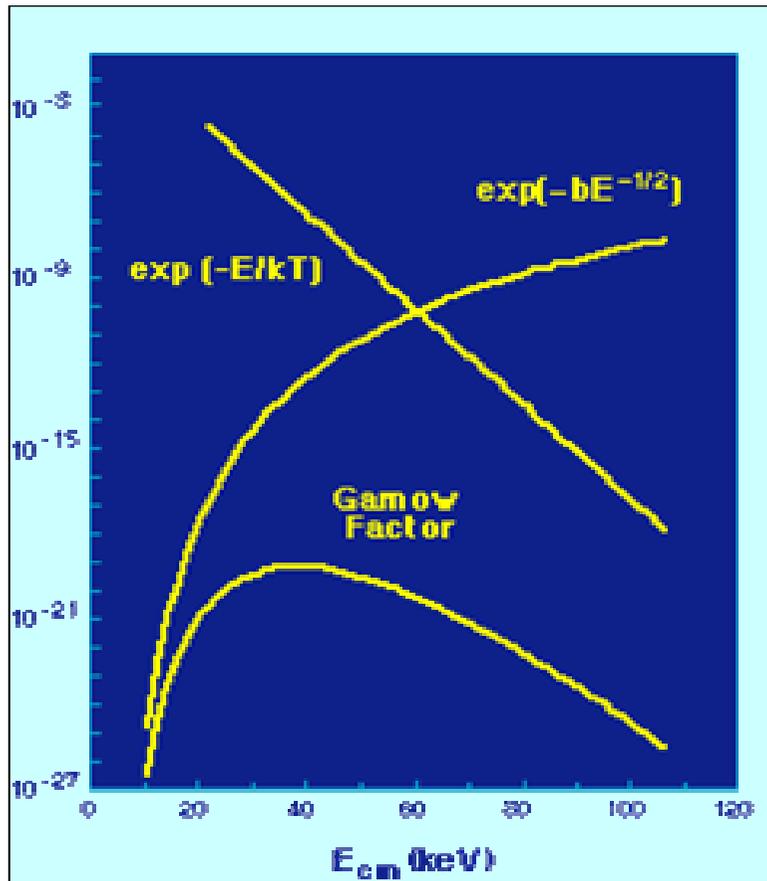
$$\Delta M \approx 926 \text{ MeV}$$

ESPALHAMENTO



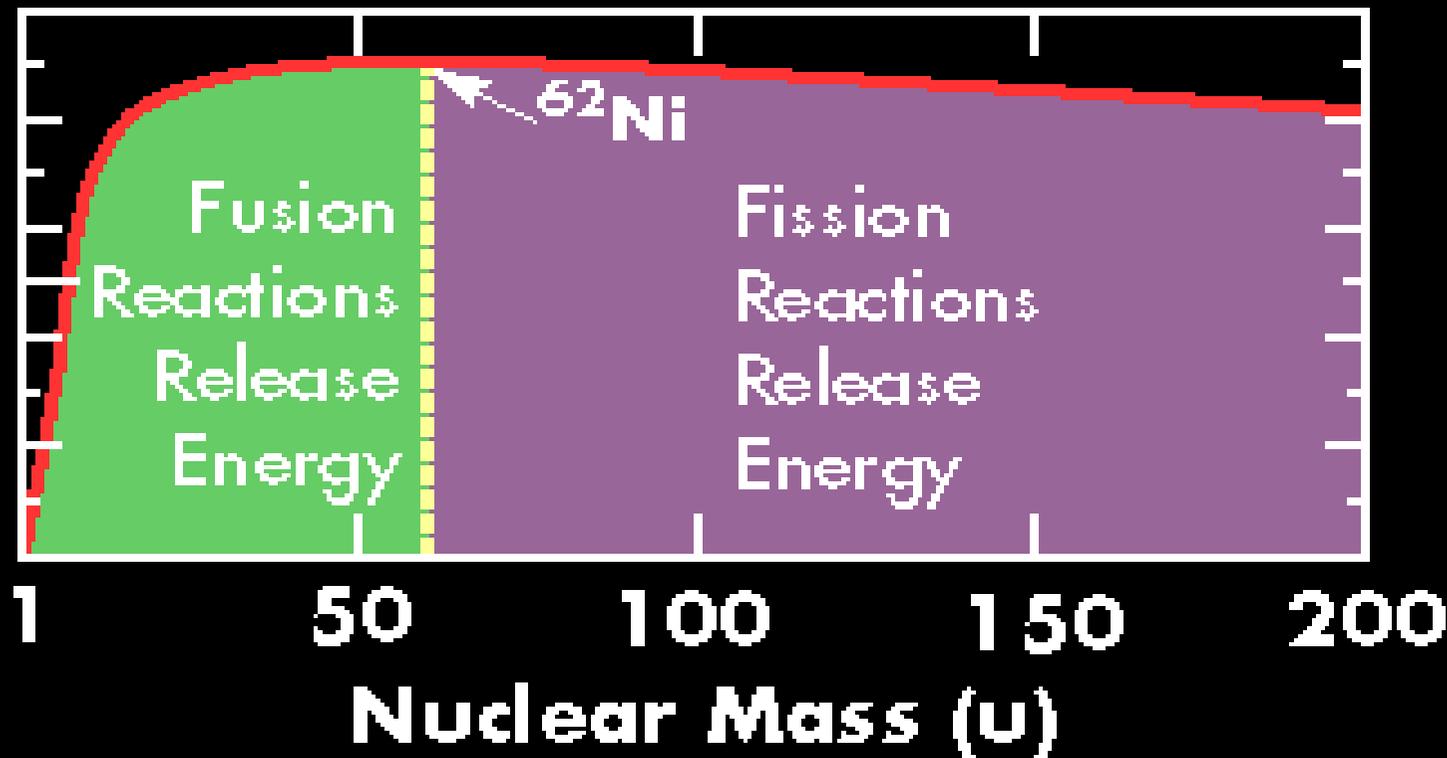
TUNELAMENTO



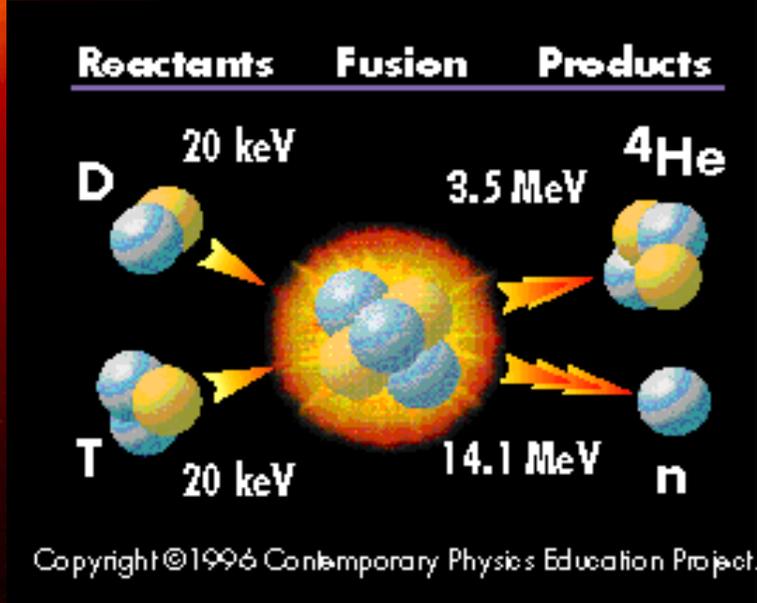
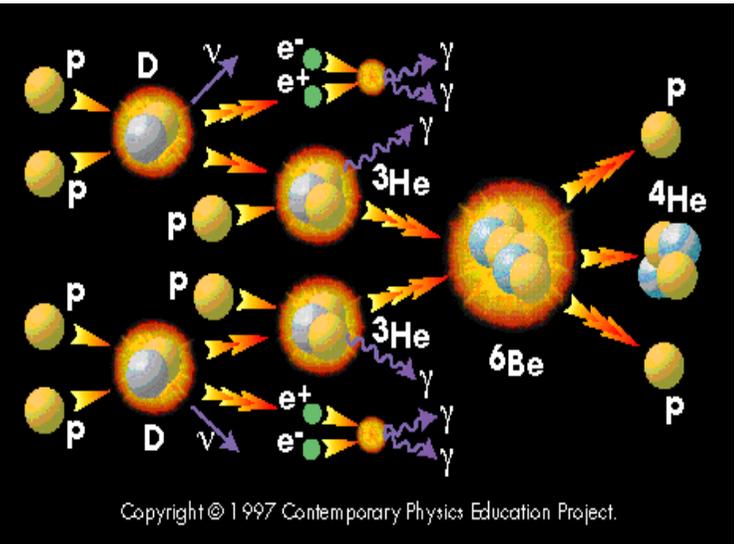
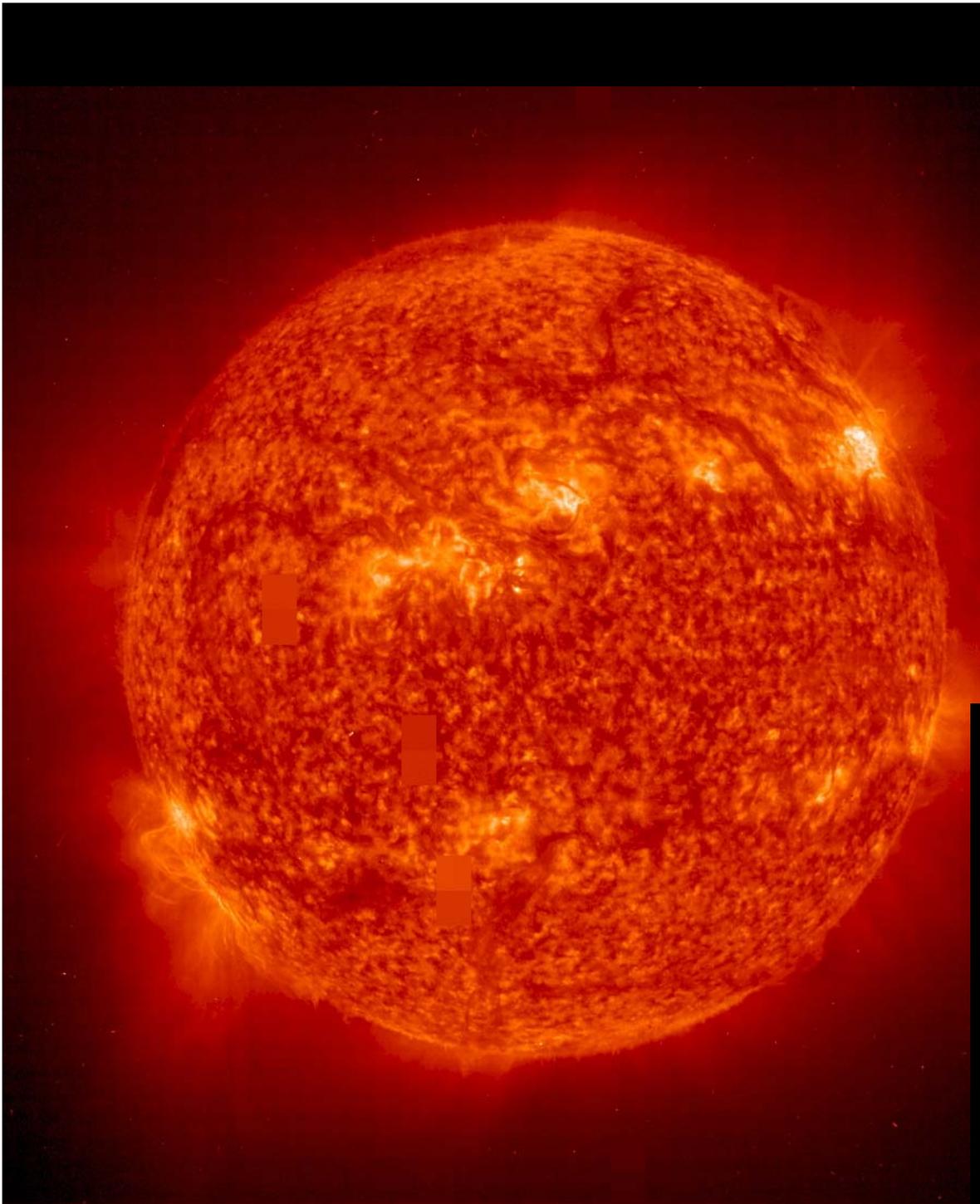


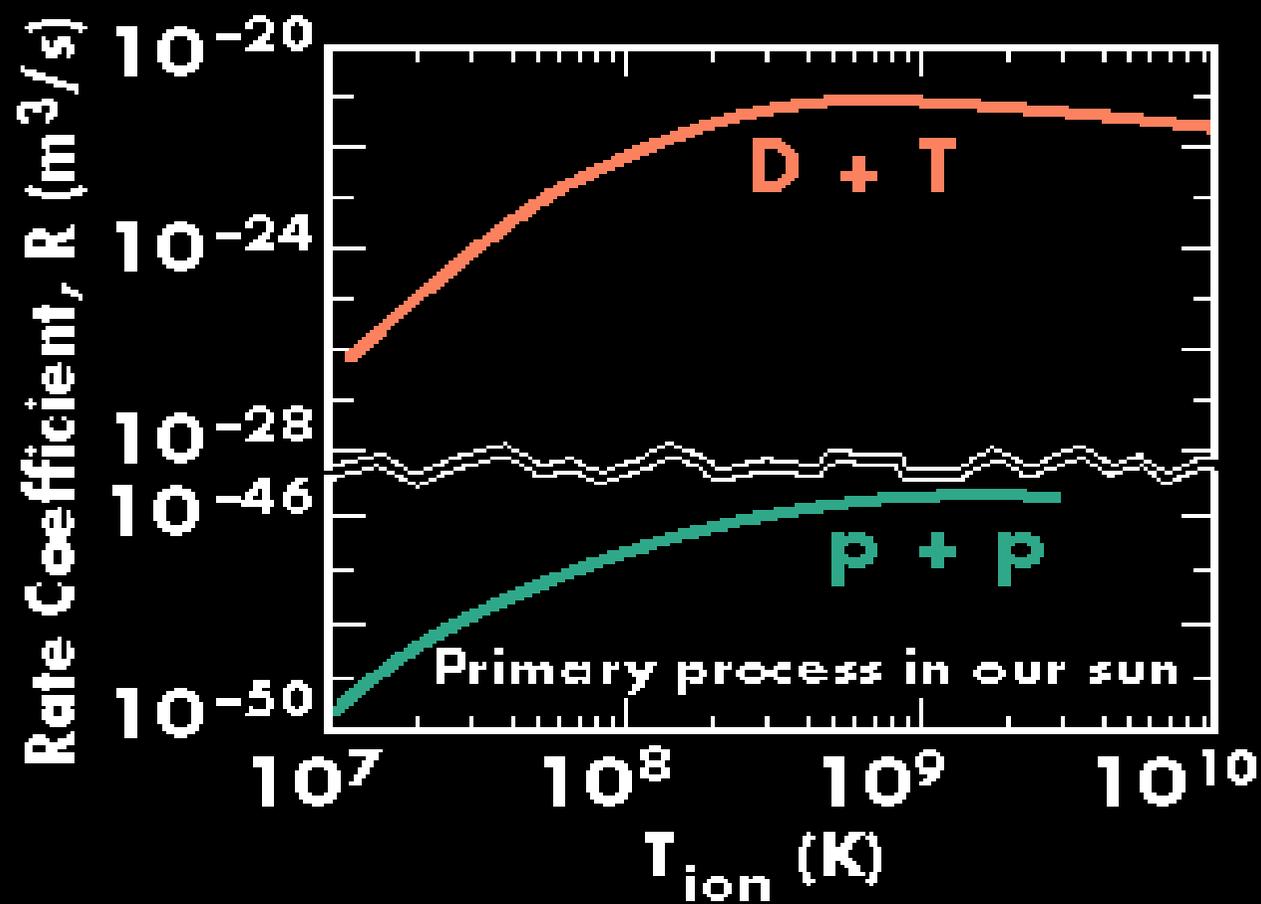
Binding Energy
Per Nucleon (MeV)

10
5
0

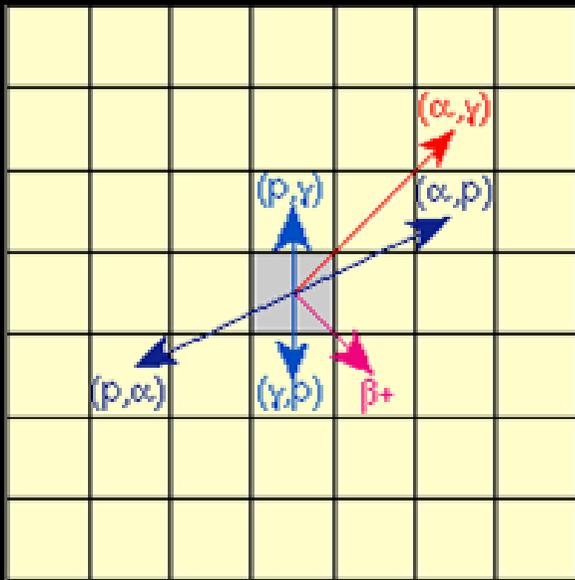


Copyright © 1996 Contemporary Physics Education Project.

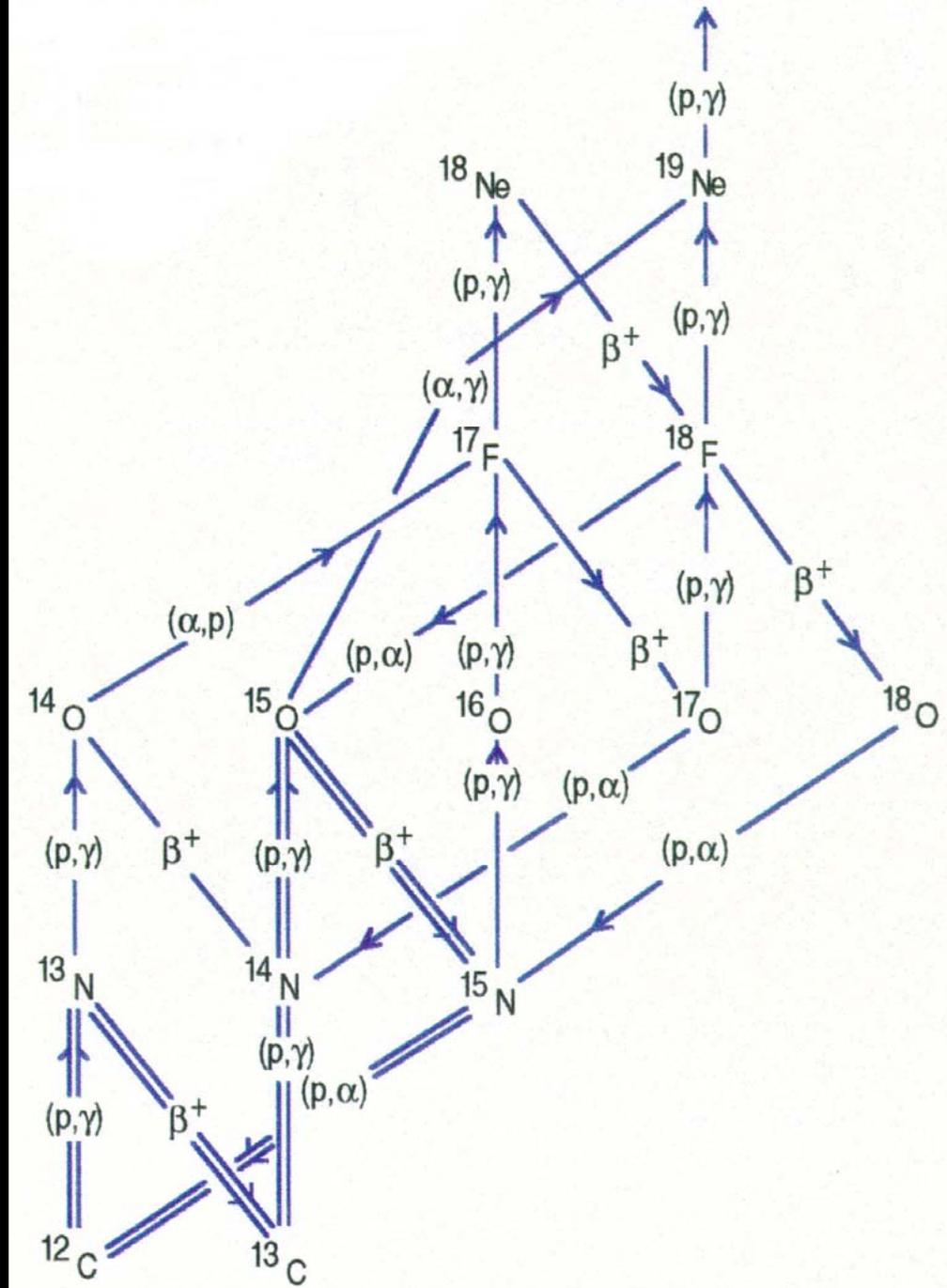
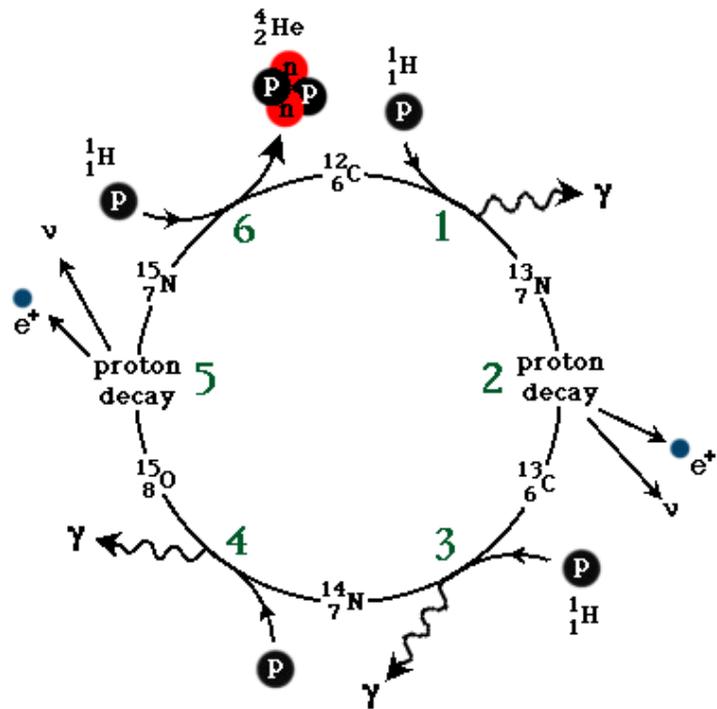


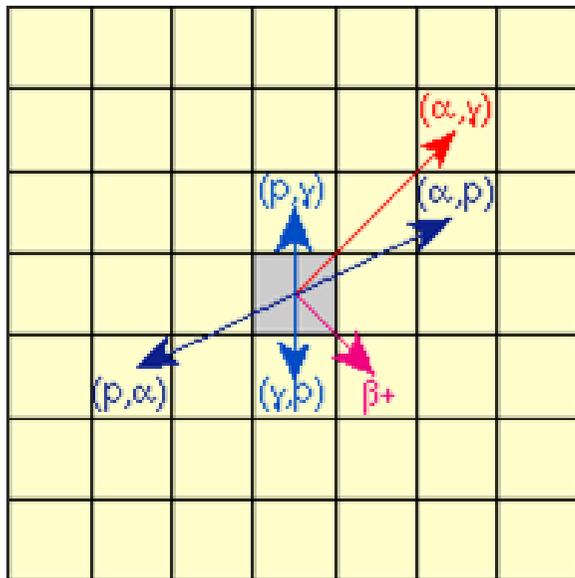


Copyright © 1996 Contemporary Physics Education Project.

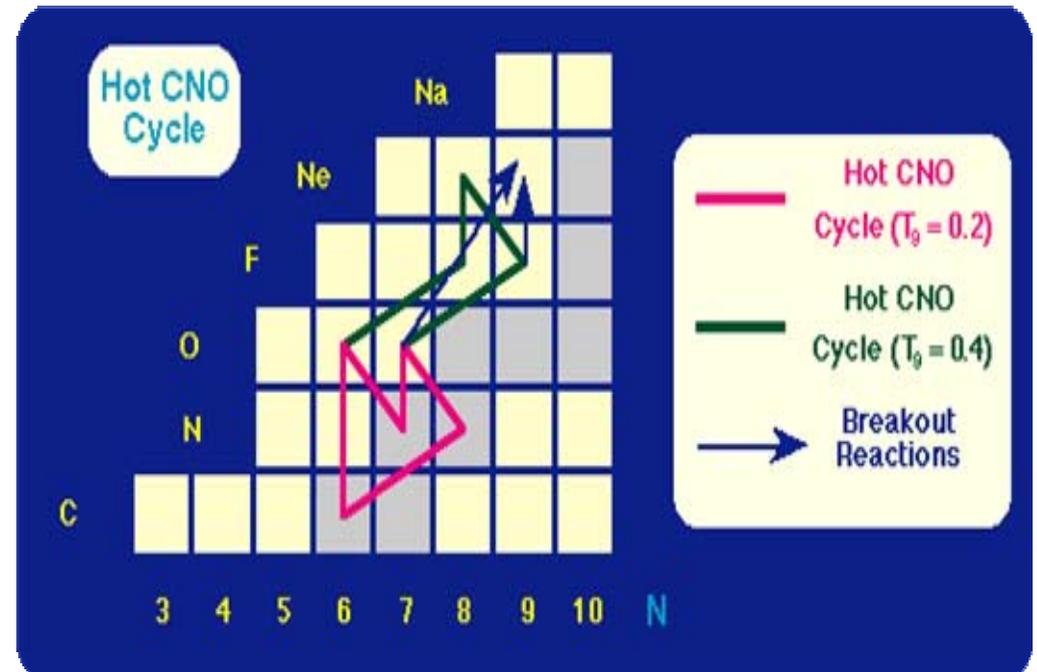
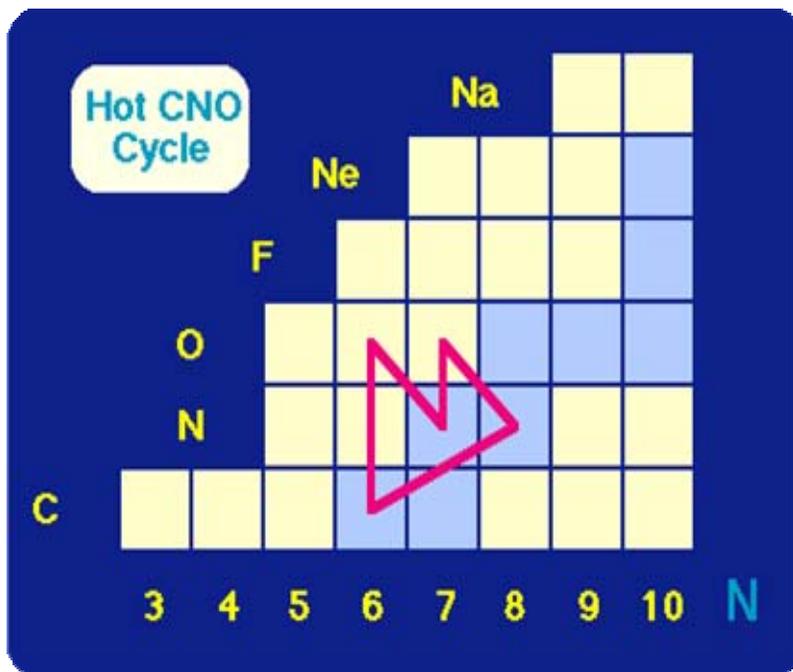
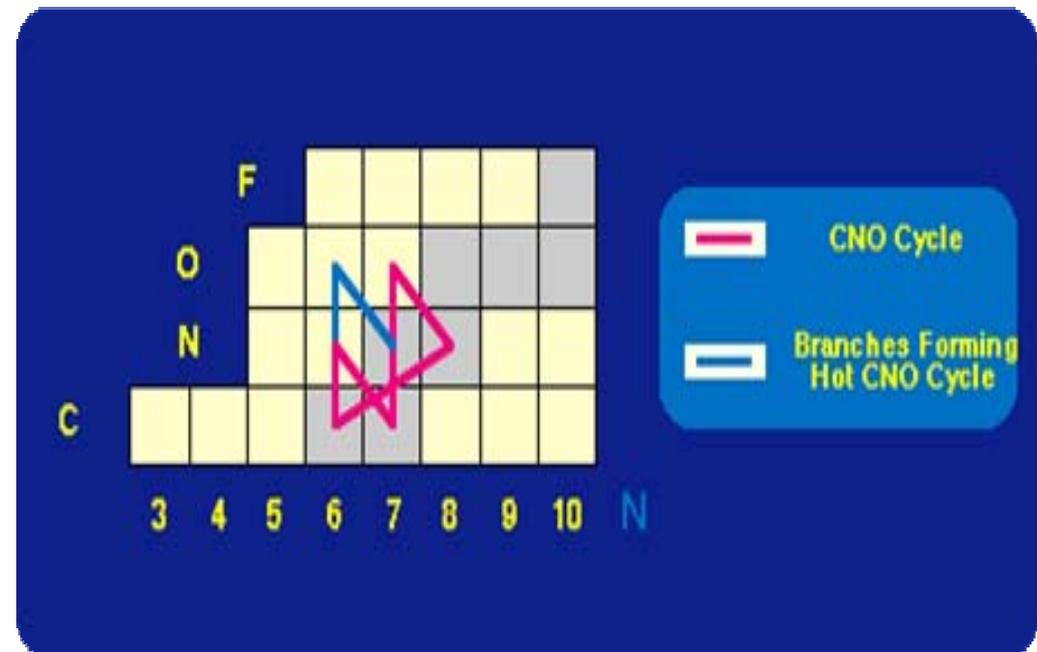


Carbon-Nitrogen-Oxygen Cycle

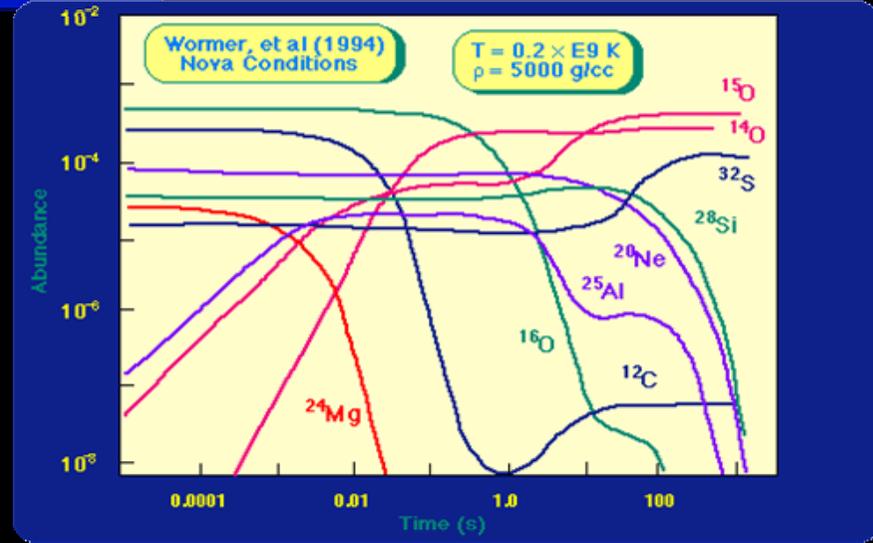
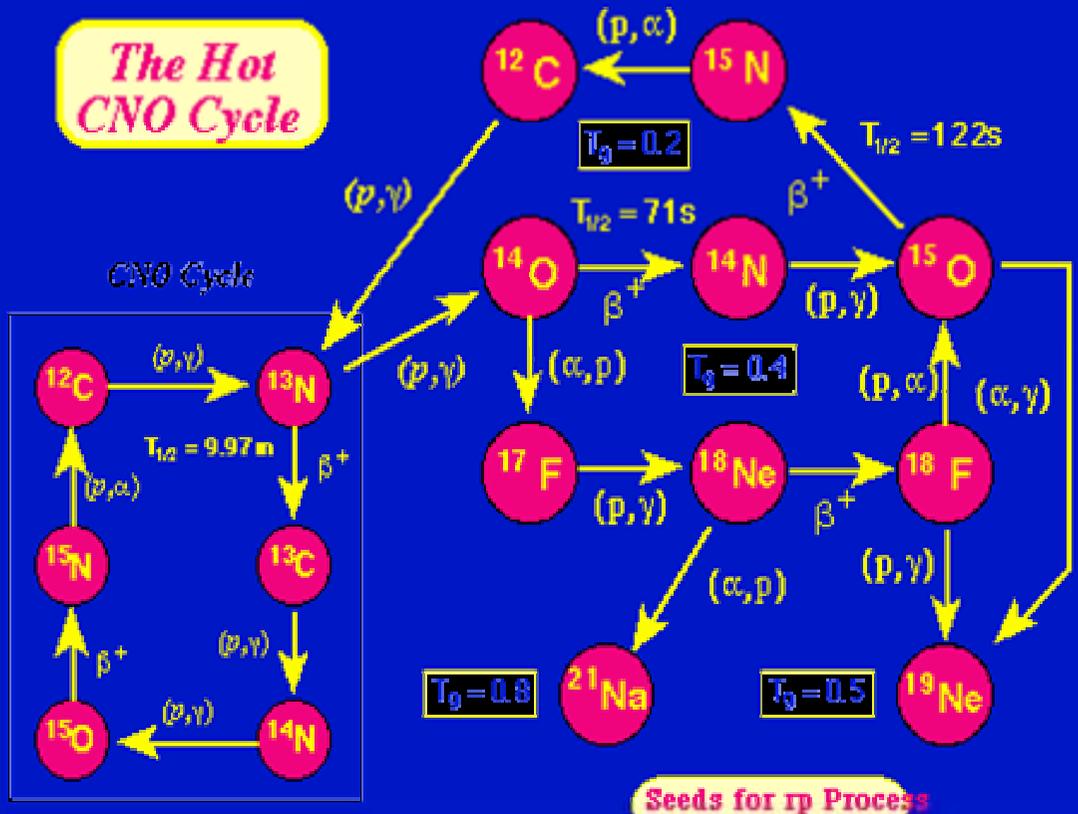




Reaction Vectors



The Hot CNO Cycle

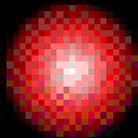




n



p



p



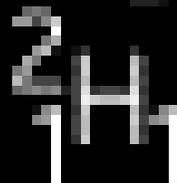
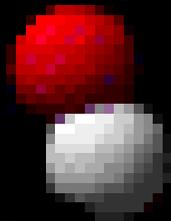
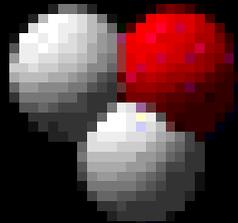
${}^2\text{H}$

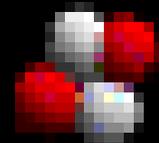
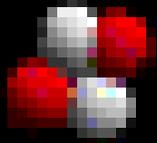


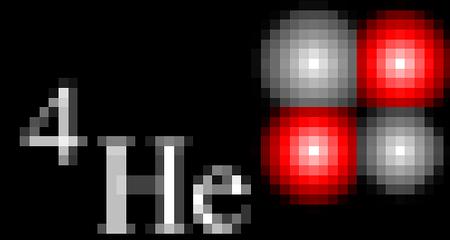
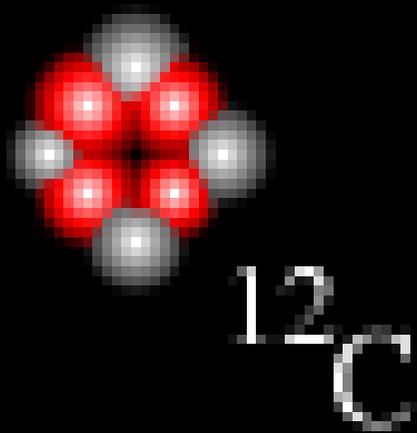
${}^3_1\text{H}_2$

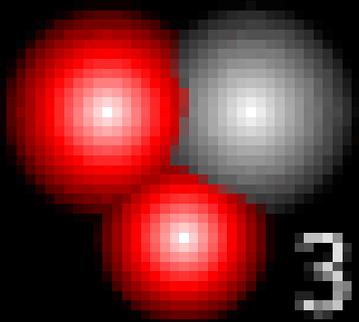


${}^2_1\text{H}_1$

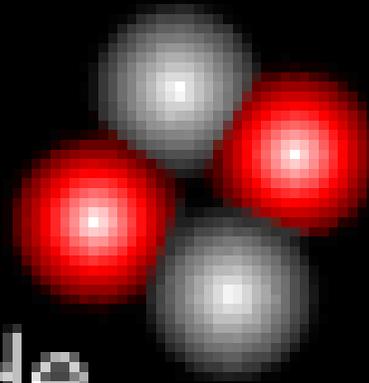








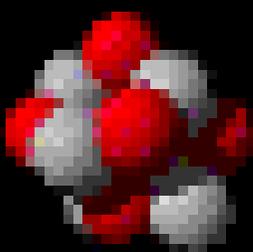
^3He



^4He



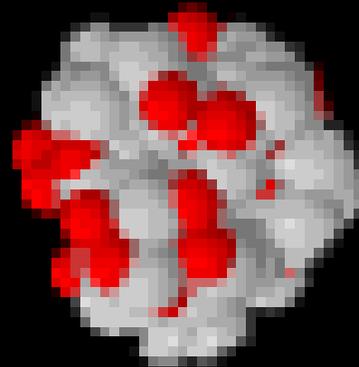
p



^{12}C



n



^{235}U

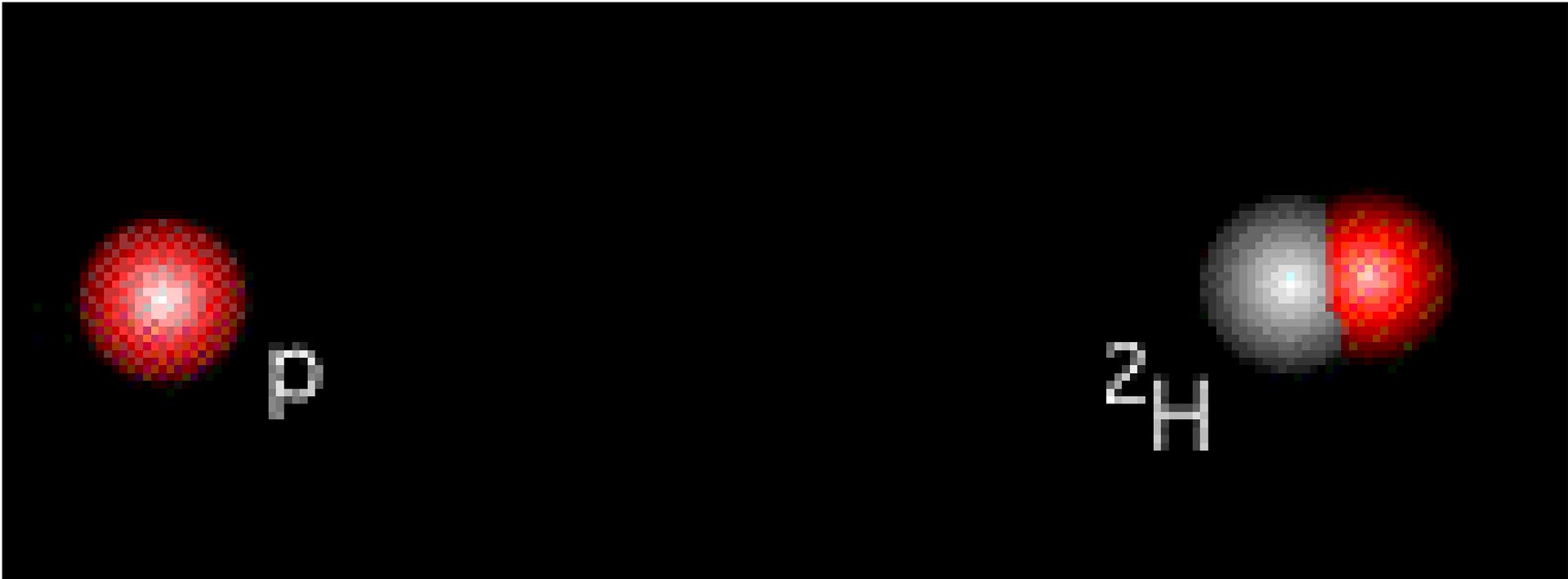
92

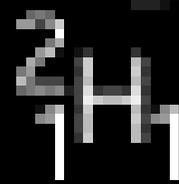
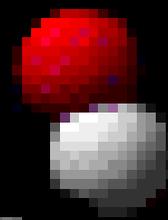
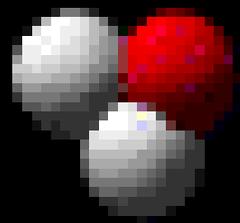
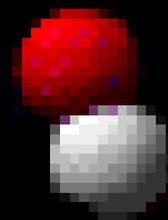
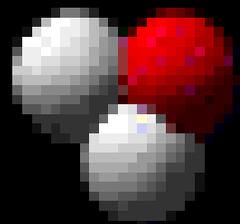
143

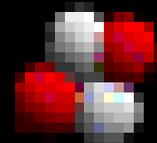
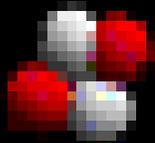
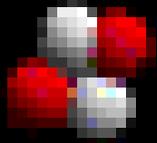
ESTRUTURA NUCLEAR

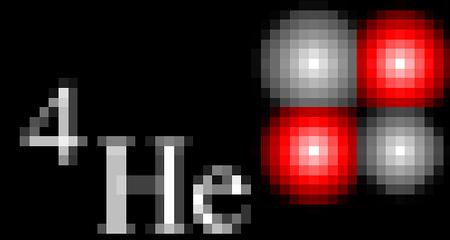
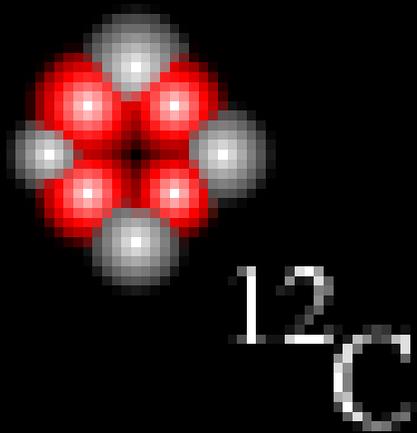
MECANISMO DE REAÇÕES

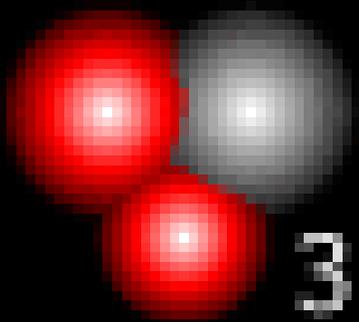
DINÂMICA DE REAÇÕES



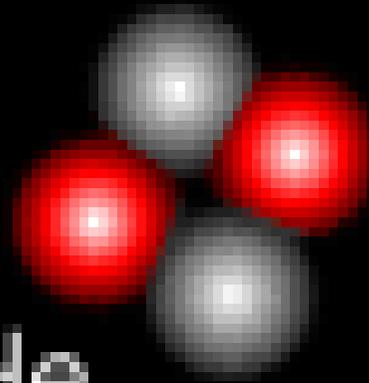








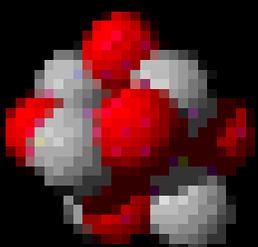
^3He



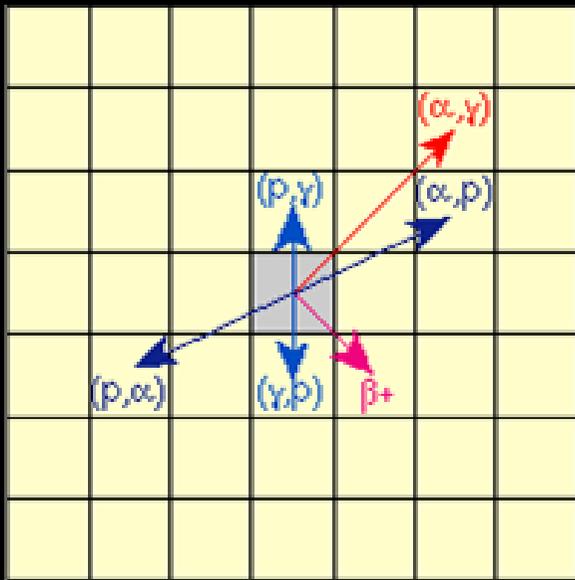
^4He



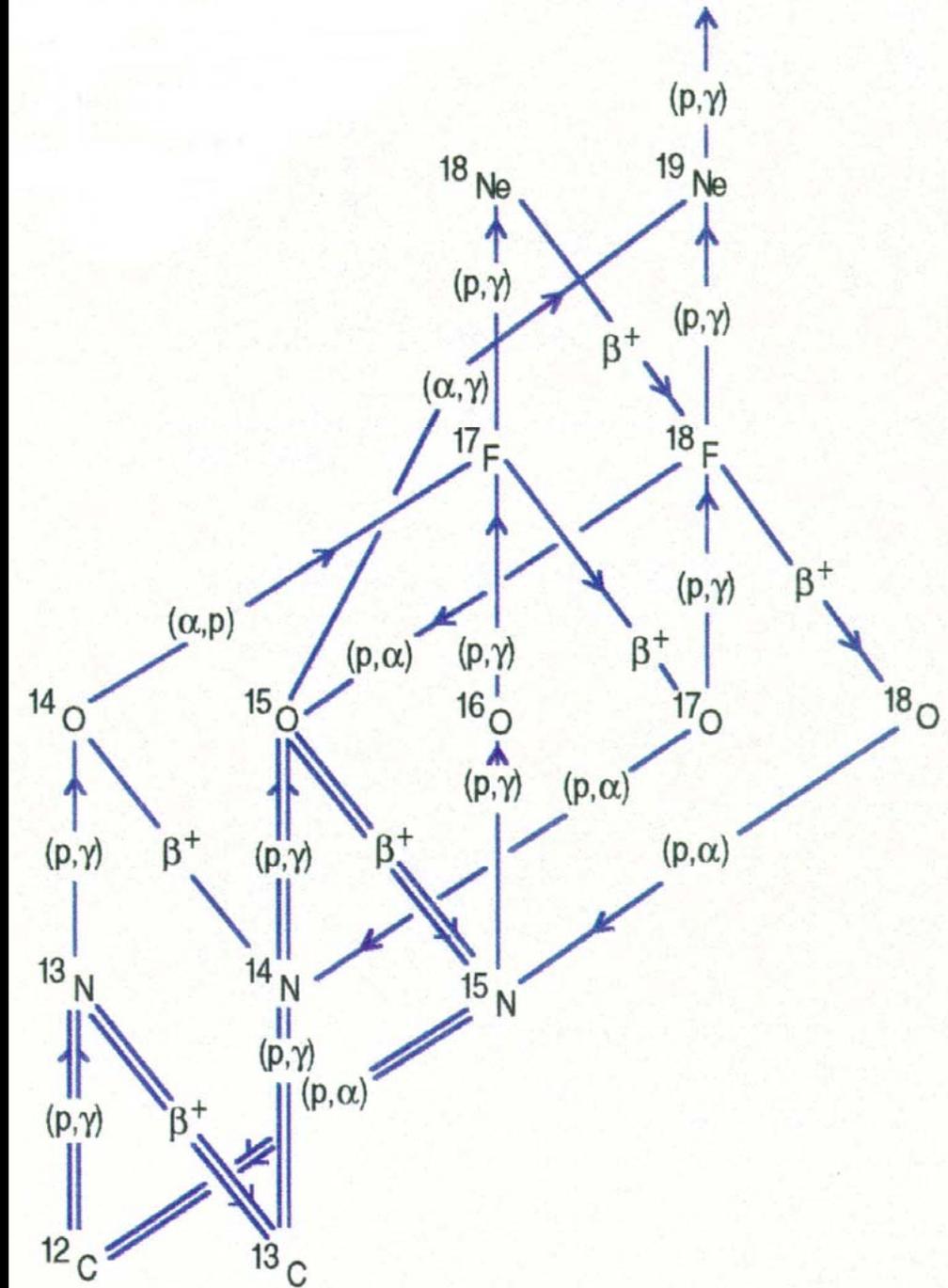
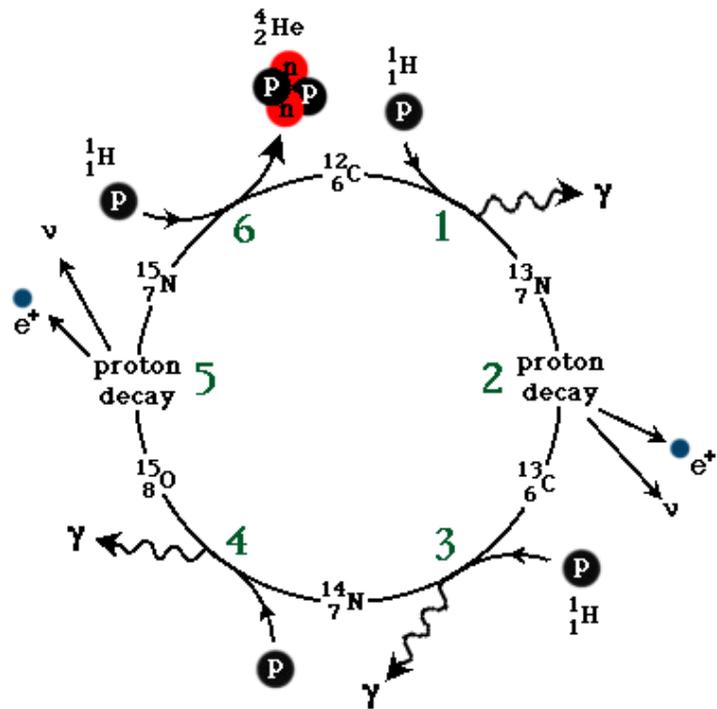
p



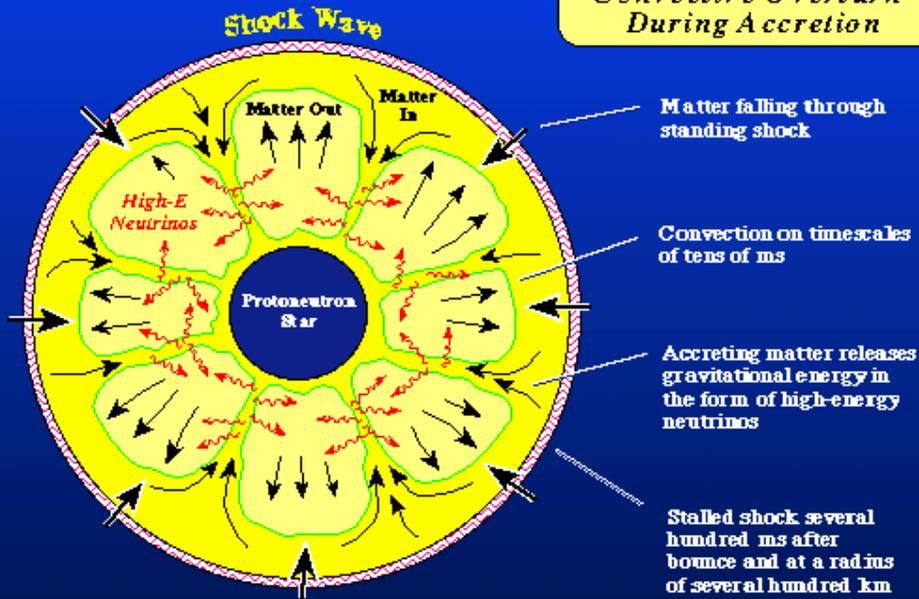
^{12}C



Carbon-Nitrogen-Oxygen Cycle



Convective Overturn During Accretion



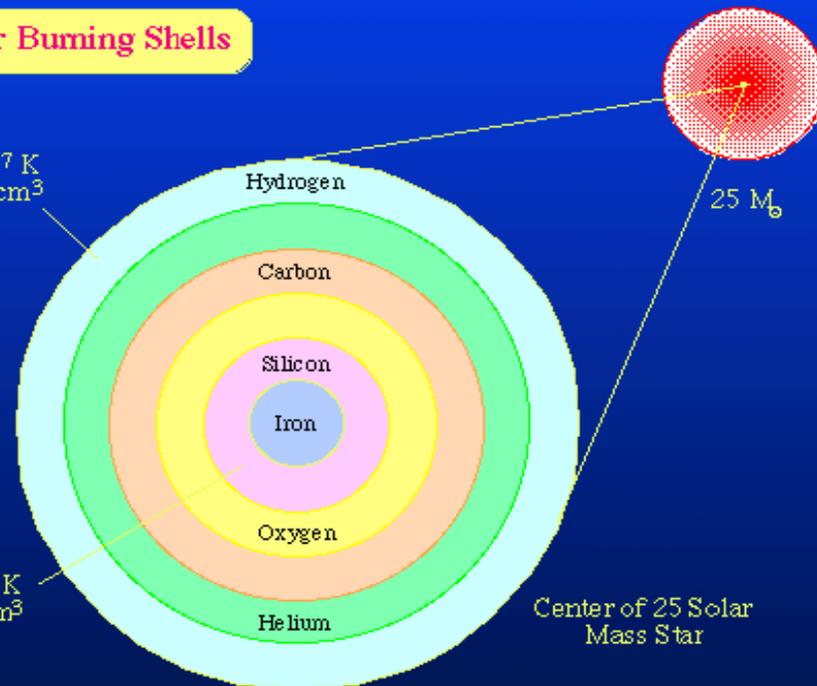
Stellar Burning Shells

$$T = 2 \times 10^7 \text{ K}$$

$$\rho = 10^2 \text{ g/cm}^3$$

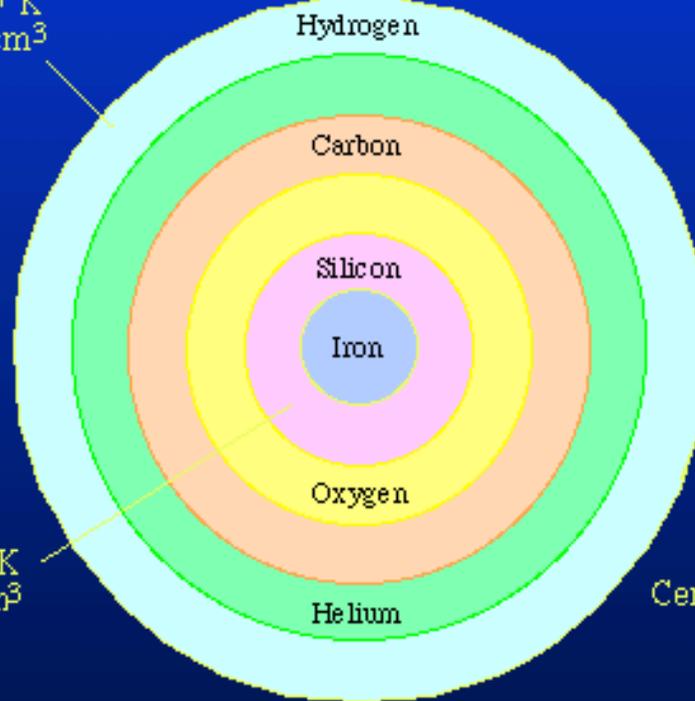
$$T = 4 \times 10^9 \text{ K}$$

$$\rho = 10^7 \text{ g/cm}^3$$



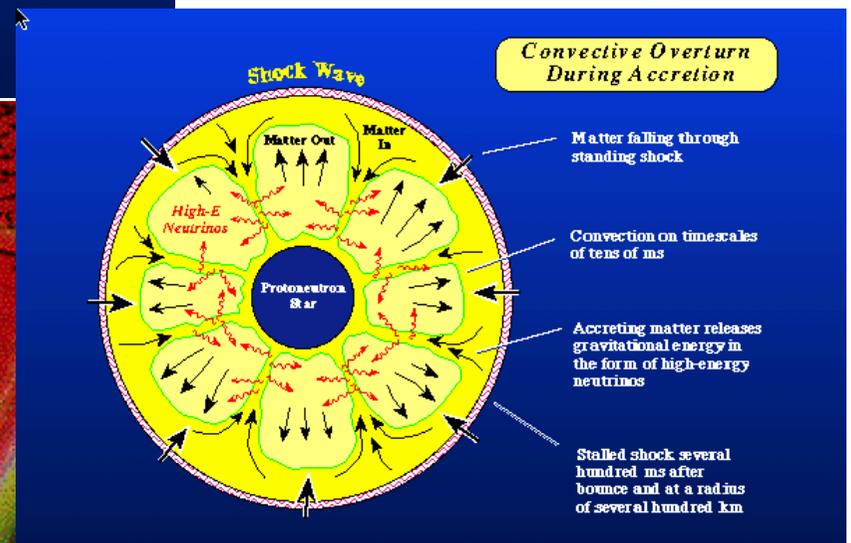
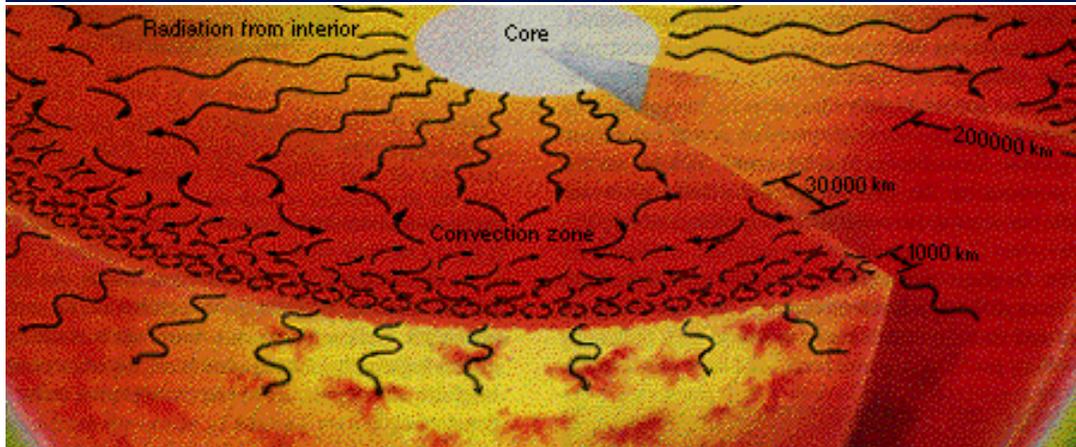
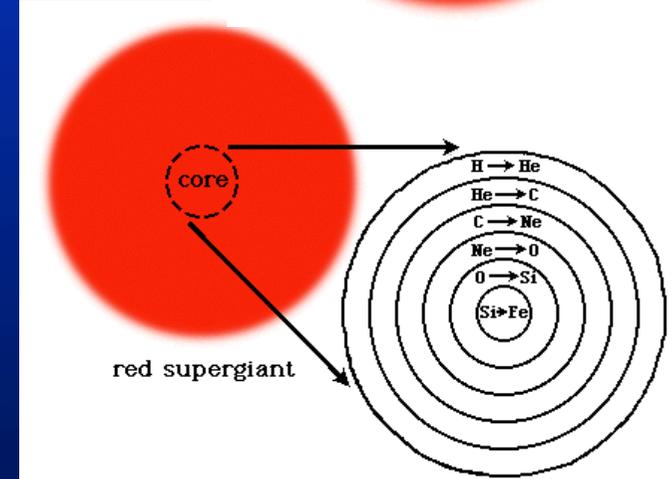
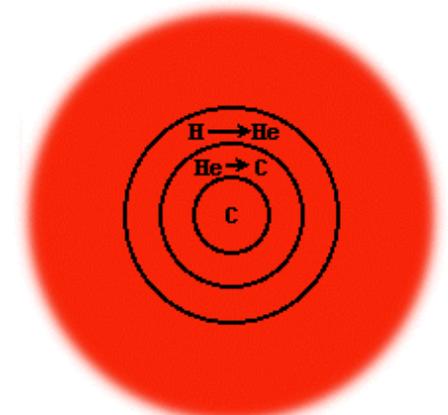
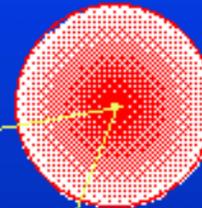
Stellar Burning Shells

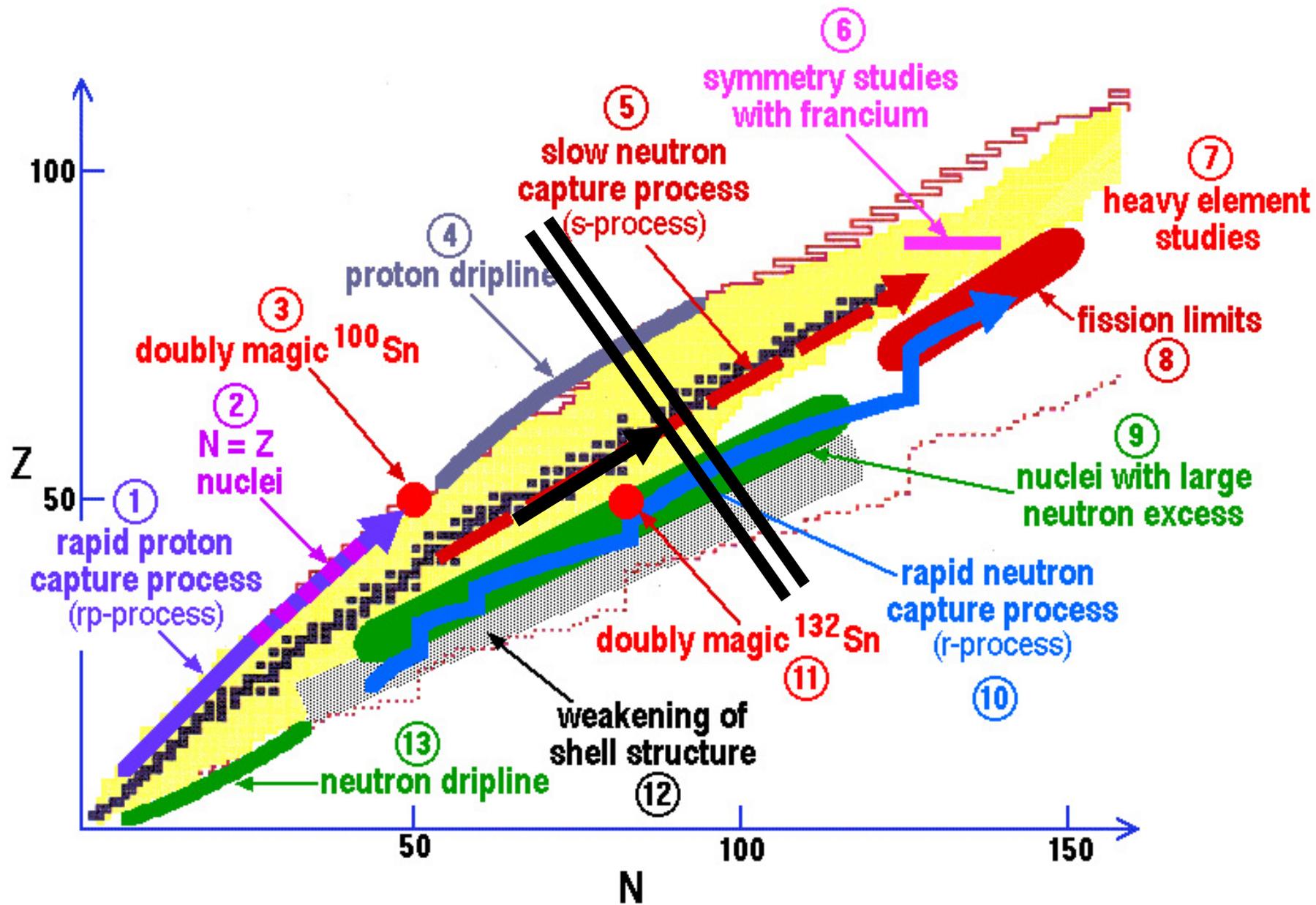
$T = 2 \times 10^7 \text{ K}$
 $\rho = 10^2 \text{ g/cm}^3$

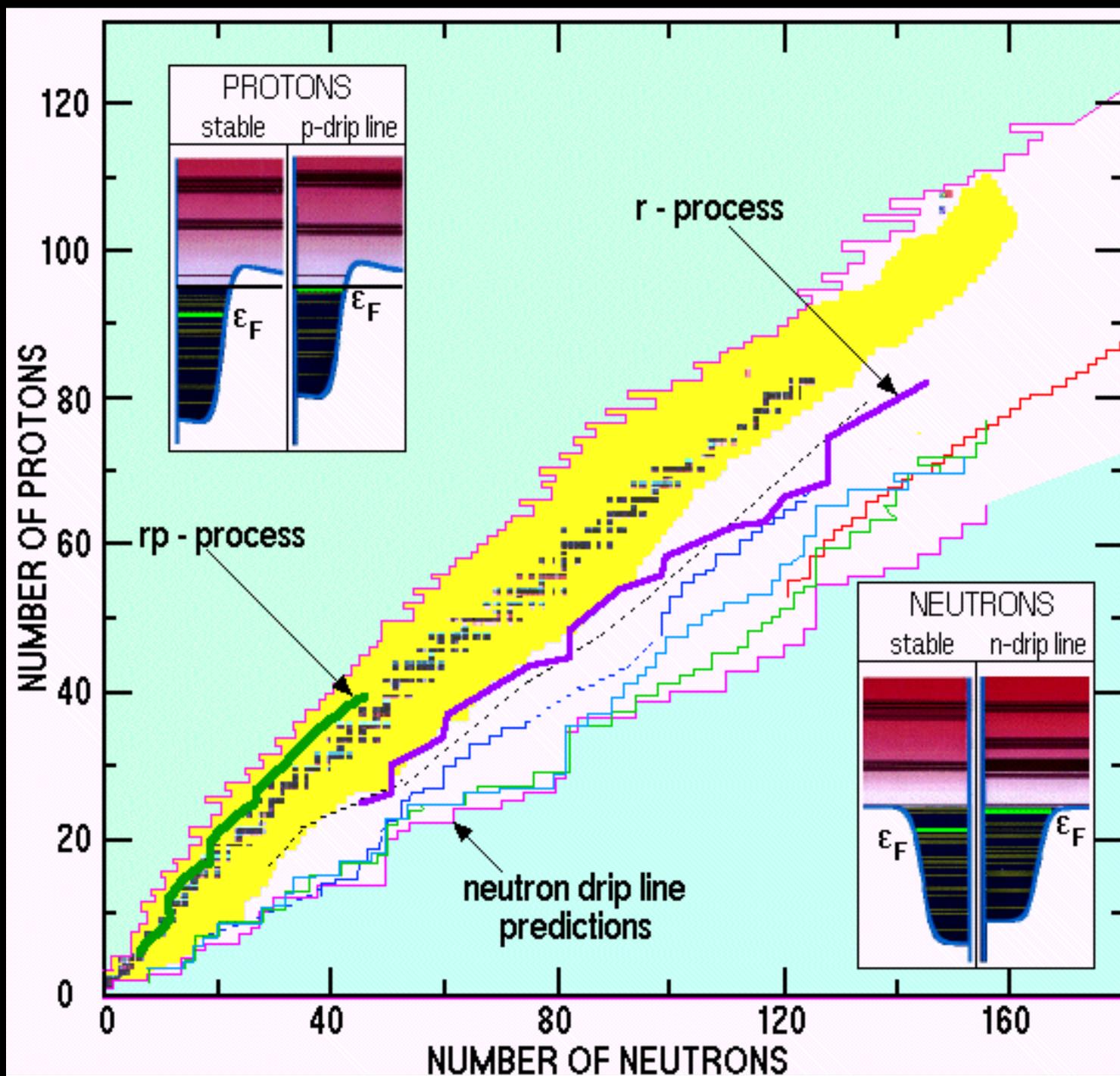


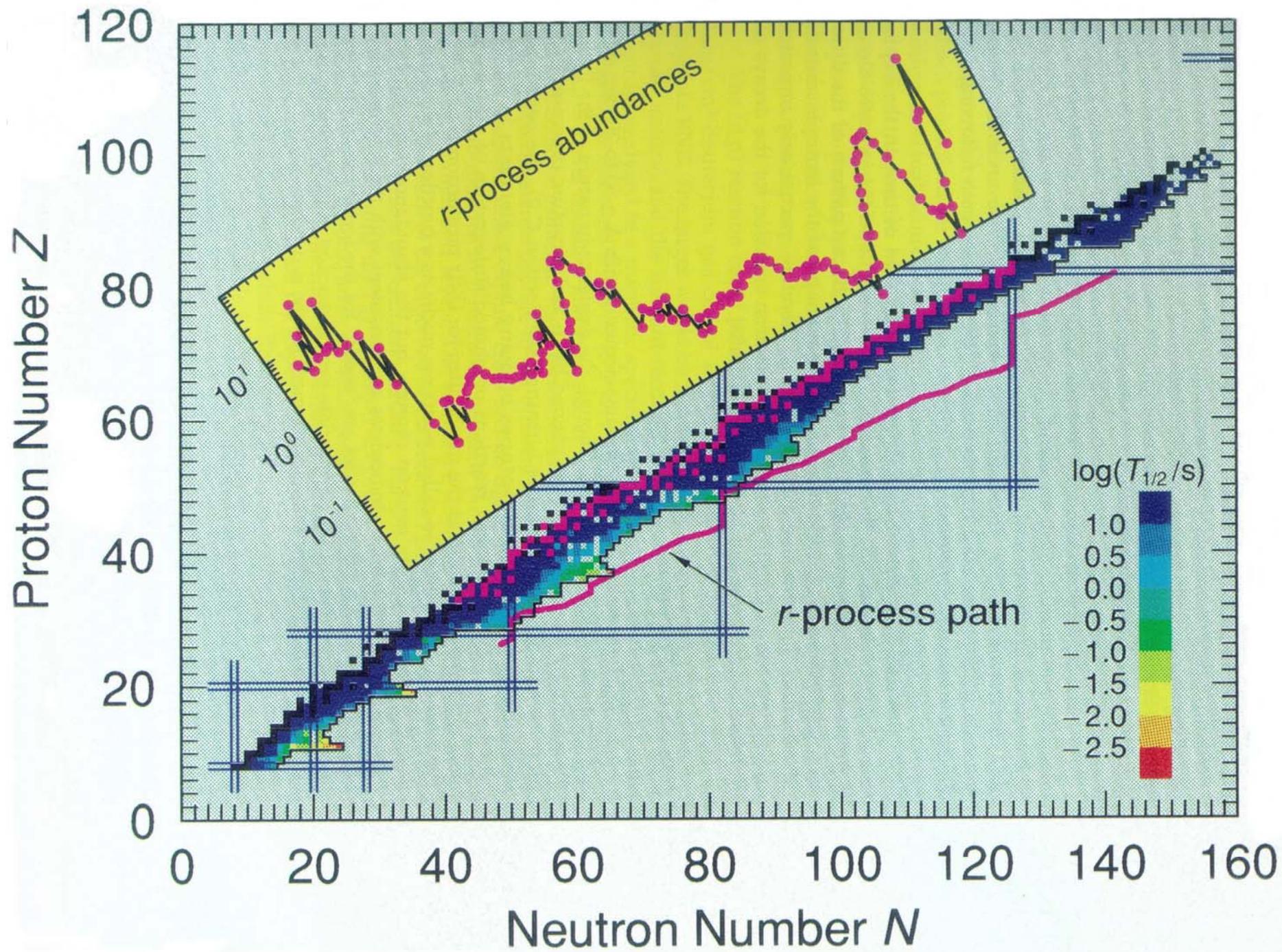
$T = 4 \times 10^9 \text{ K}$
 $\rho = 10^7 \text{ g/cm}^3$

Center of 25 Solar Mass Star

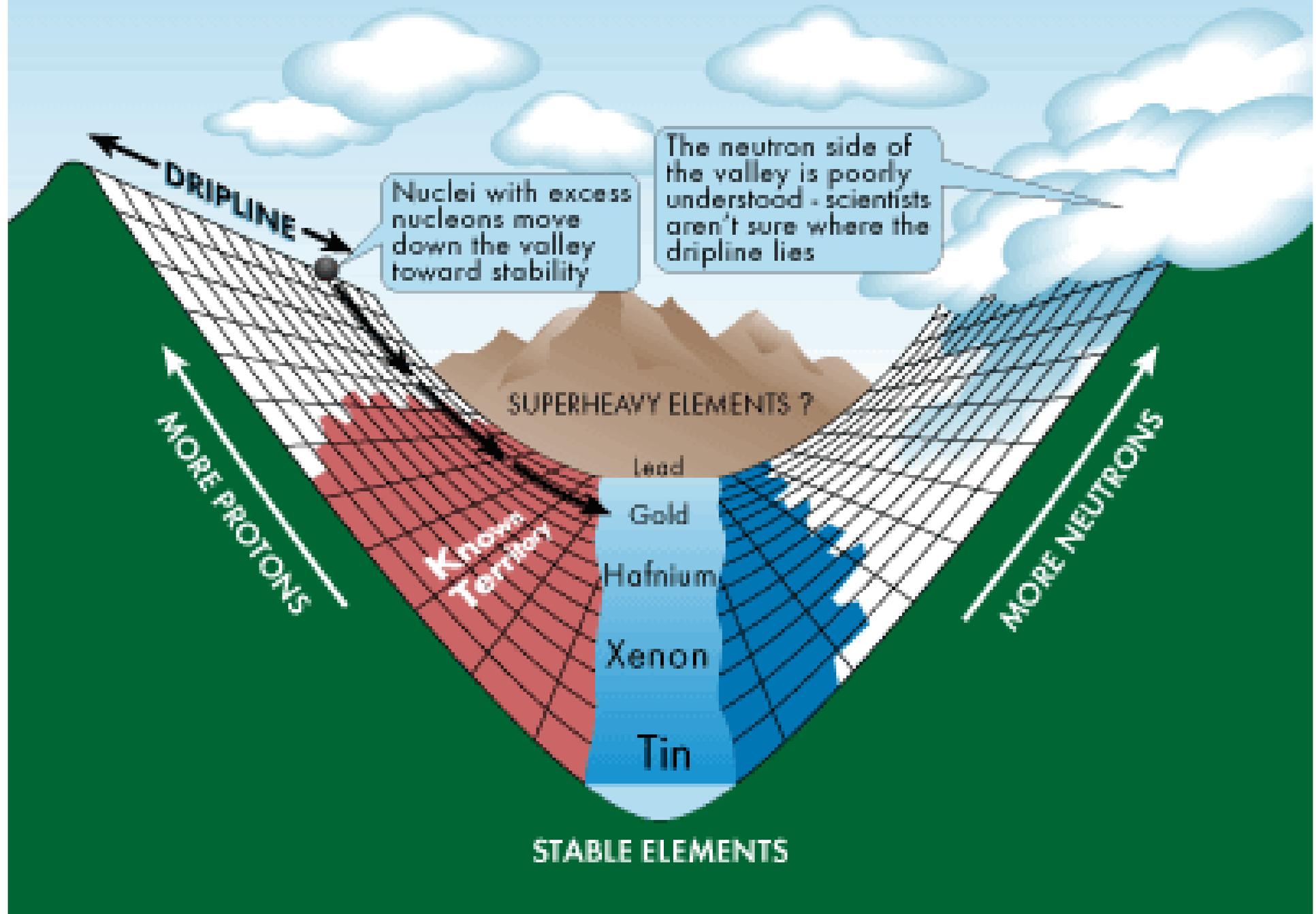








THE VALLEY OF STABILITY



BARREIRA COULOMBIANA

$$V_C = \frac{Z_A Z_a e^2}{R} = \frac{C}{R}$$

Para a colisão razante:

$$\begin{aligned} R_c &= r_0 A_A^{1/3} + r_0 A_a^{1/3} \\ &= r_0 (A_A^{1/3} + A_a^{1/3}) \end{aligned}$$

$$V_{\text{Coul}} = \frac{Z_A Z_a e^2}{r_0 (A_A^{1/3} + A_a^{1/3})}$$

$$\begin{aligned} E_{\text{cm}} > V_{\text{coul}} &\implies T_1(E_{\text{cm}}) = 1 \\ E_{\text{cm}} < V_{\text{coul}} &\implies T_1(E_{\text{cm}}) = 0 \end{aligned}$$

$$e^2 = 1.44 \text{ MeV}\cdot\text{fm}$$

$$r_0 \sim 1.25 \text{ fm}$$

$$V_{\text{coul}} (\text{MeV})$$

REAÇÃO DIRETA

$^{12}\text{C}(^{16}\text{O}, \alpha)^{24}\text{Mg}$

E_{cm}
 $^{16}\text{O}+^{12}\text{C}$

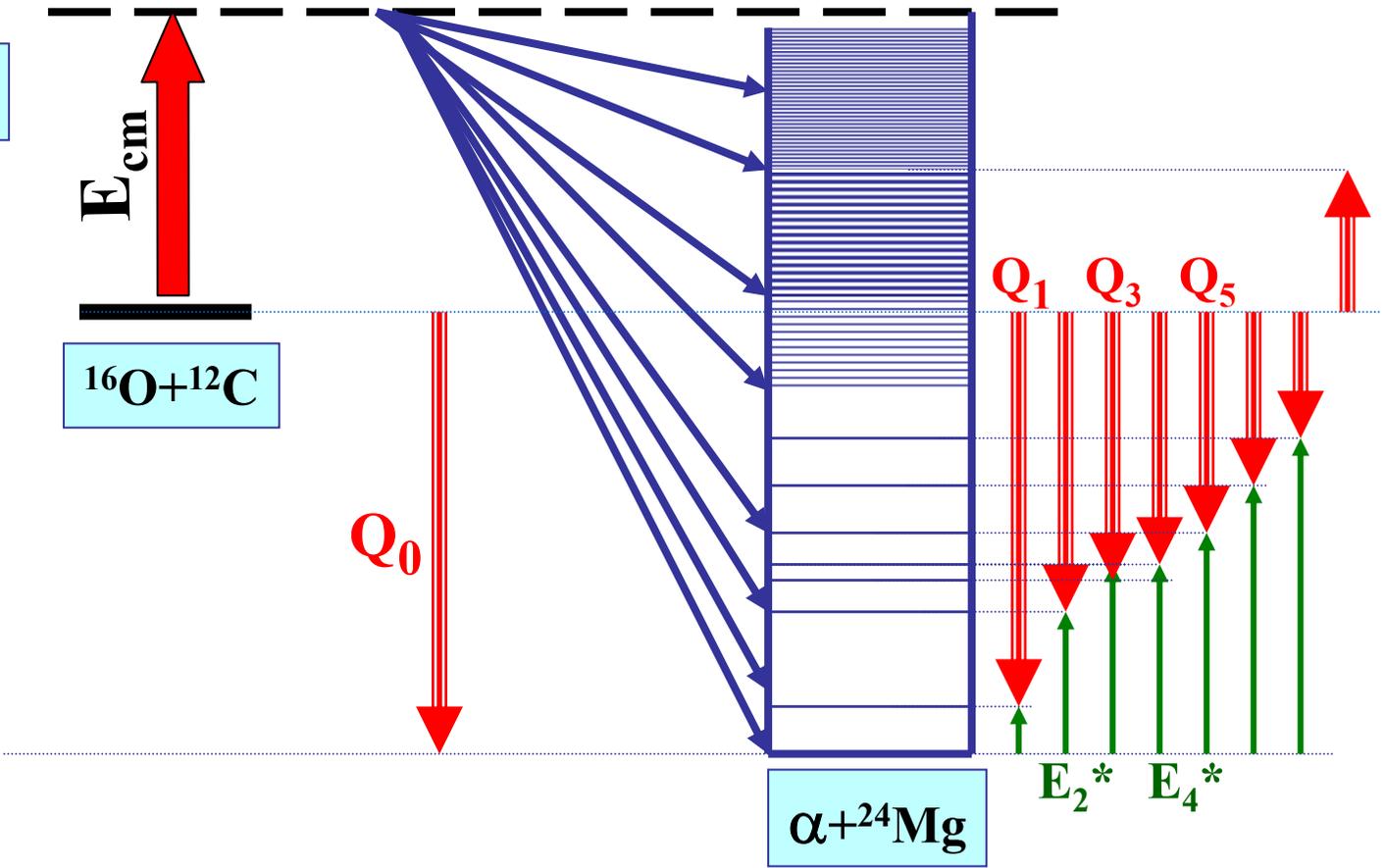
$\alpha+^{24}\text{Mg}$

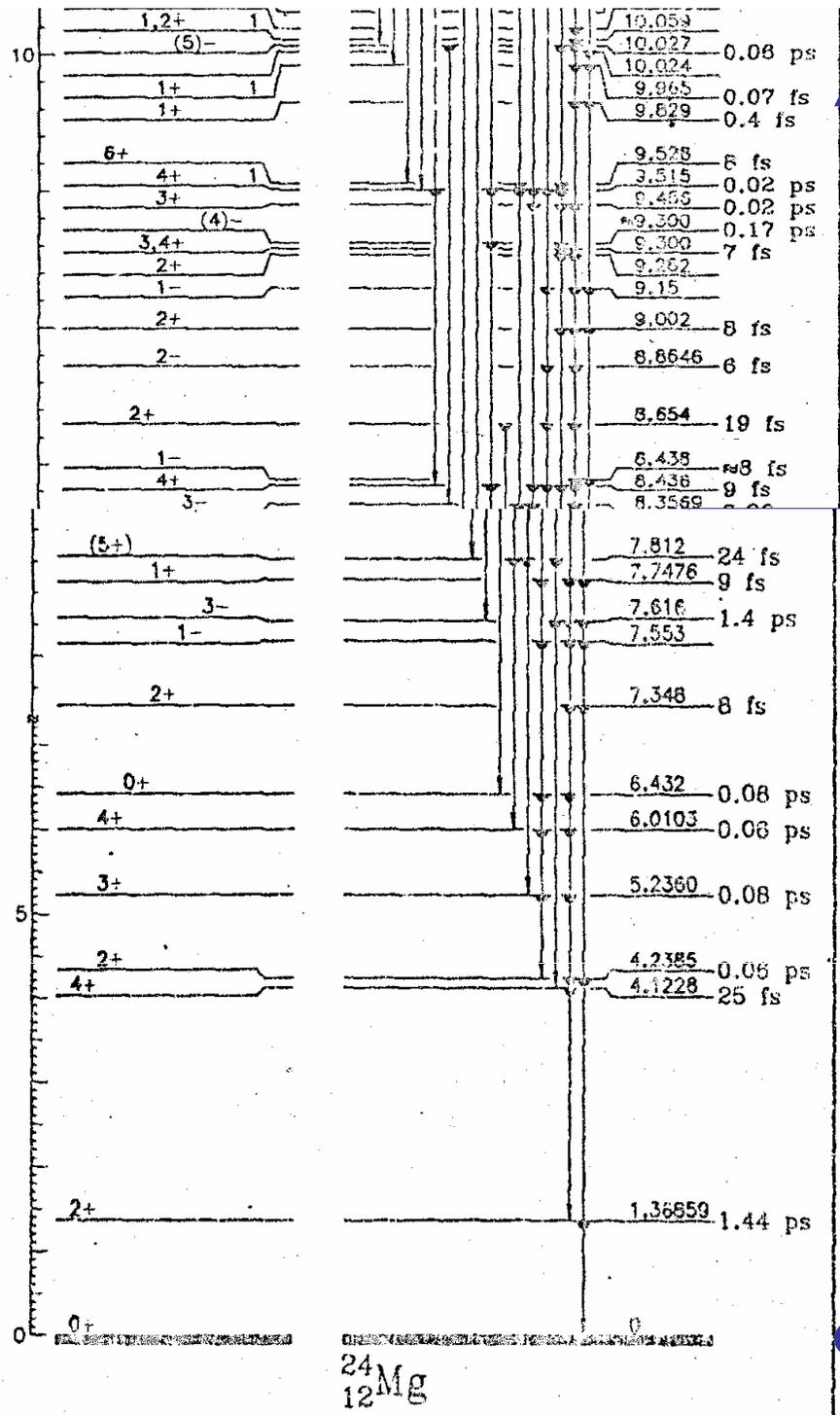
Q_0

Q_1 Q_3 Q_5

E_2^* E_4^*

- $\longrightarrow E_{\text{cm}}(\alpha)$
- $\longrightarrow E_i^*[^{24}\text{Mg}(i)]$
- $\longrightarrow Q_i[^{24}\text{Mg}(i)]$

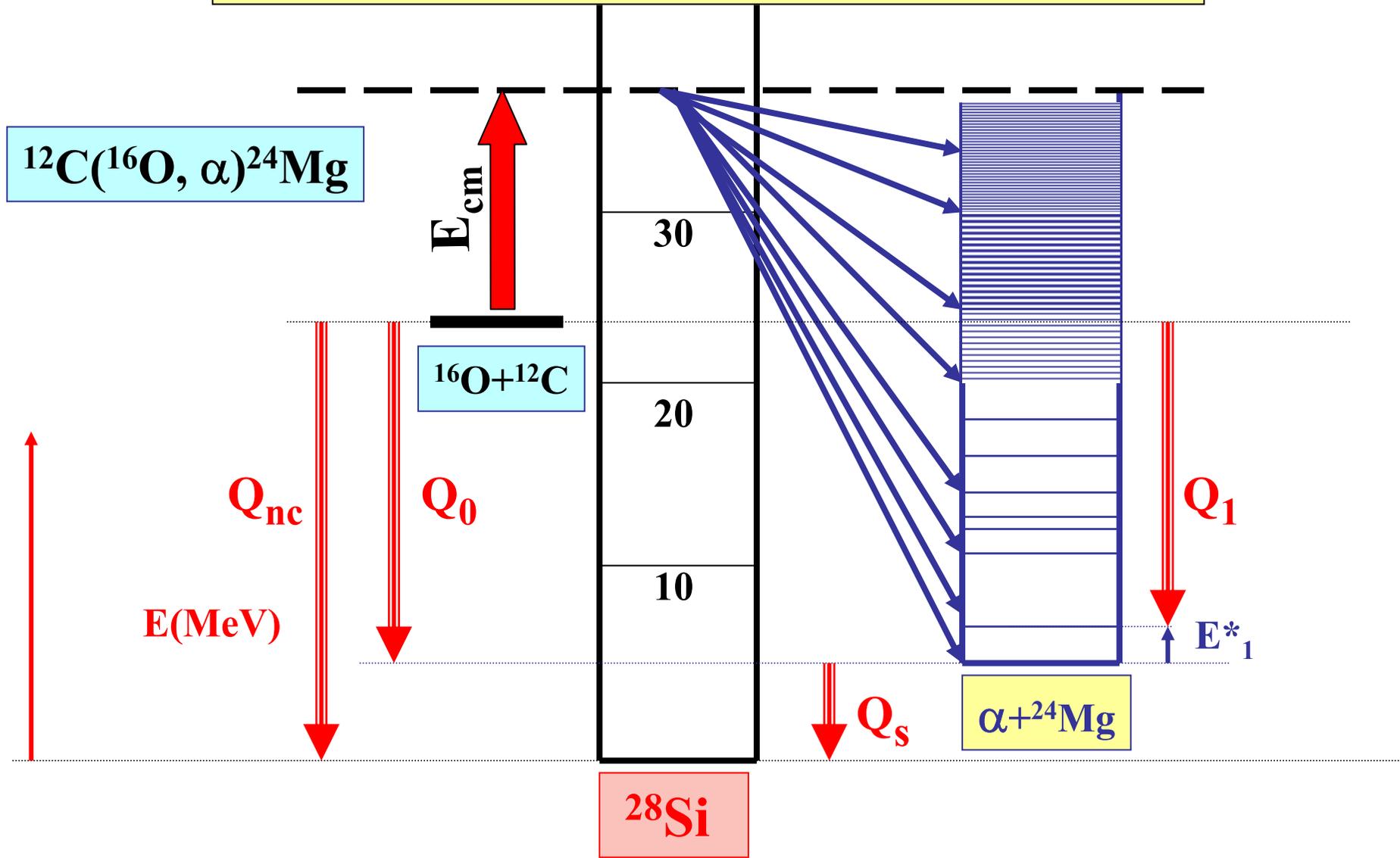




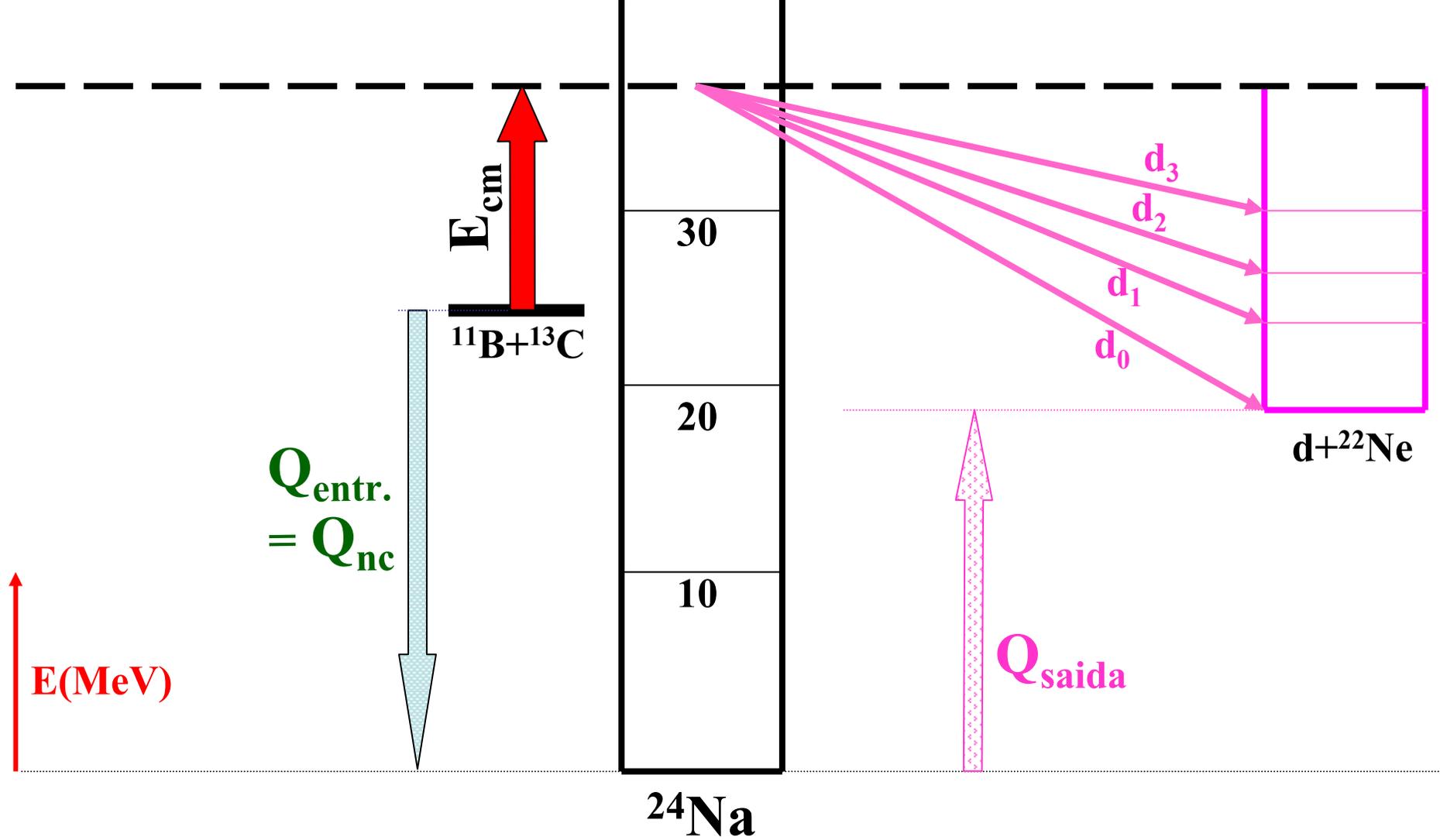
energia de excitação (E^*)

^{24}Mg
 ^{12}Mg

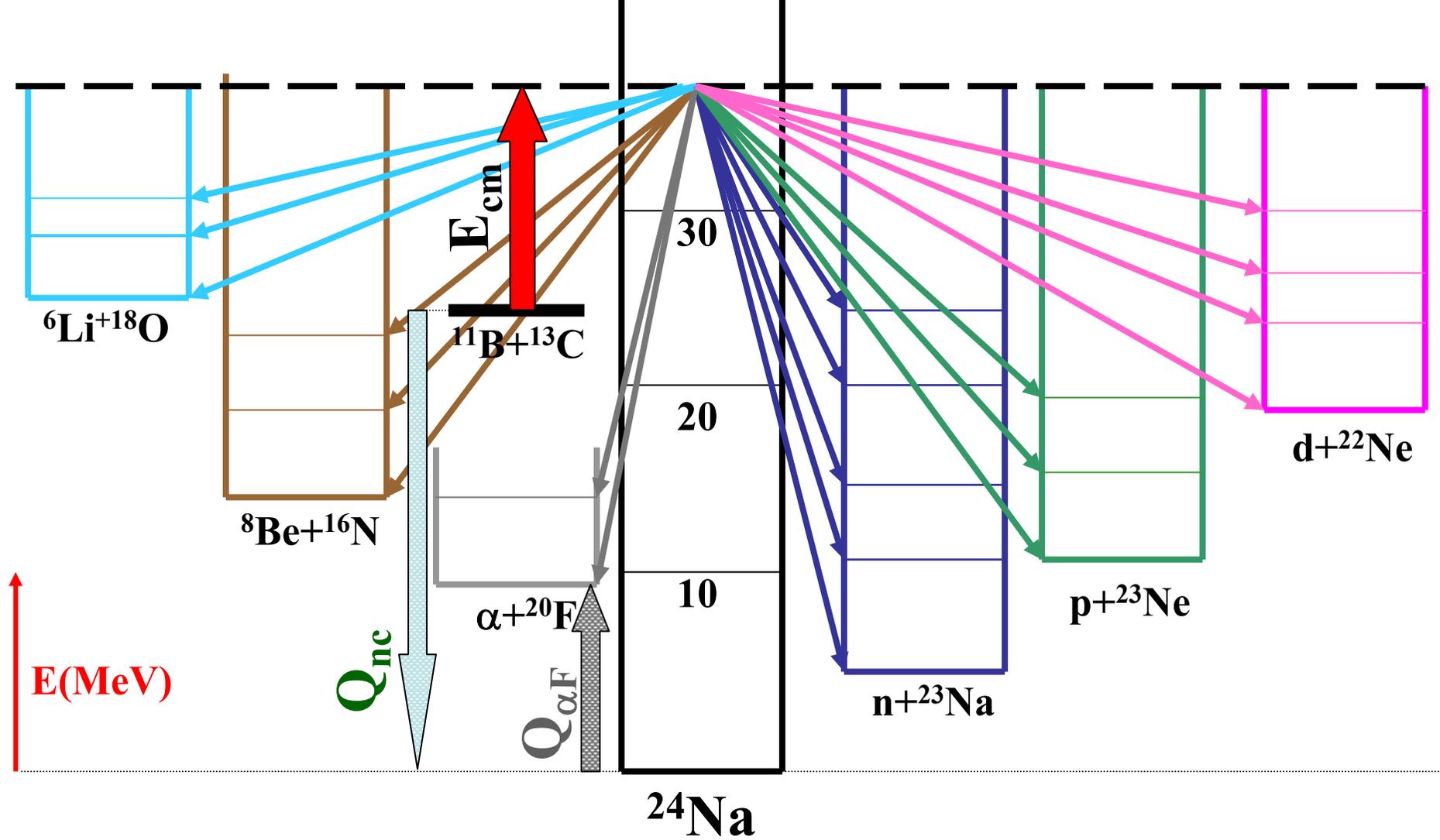
REAÇÃO VIA NÚCLEO COMPOSTO

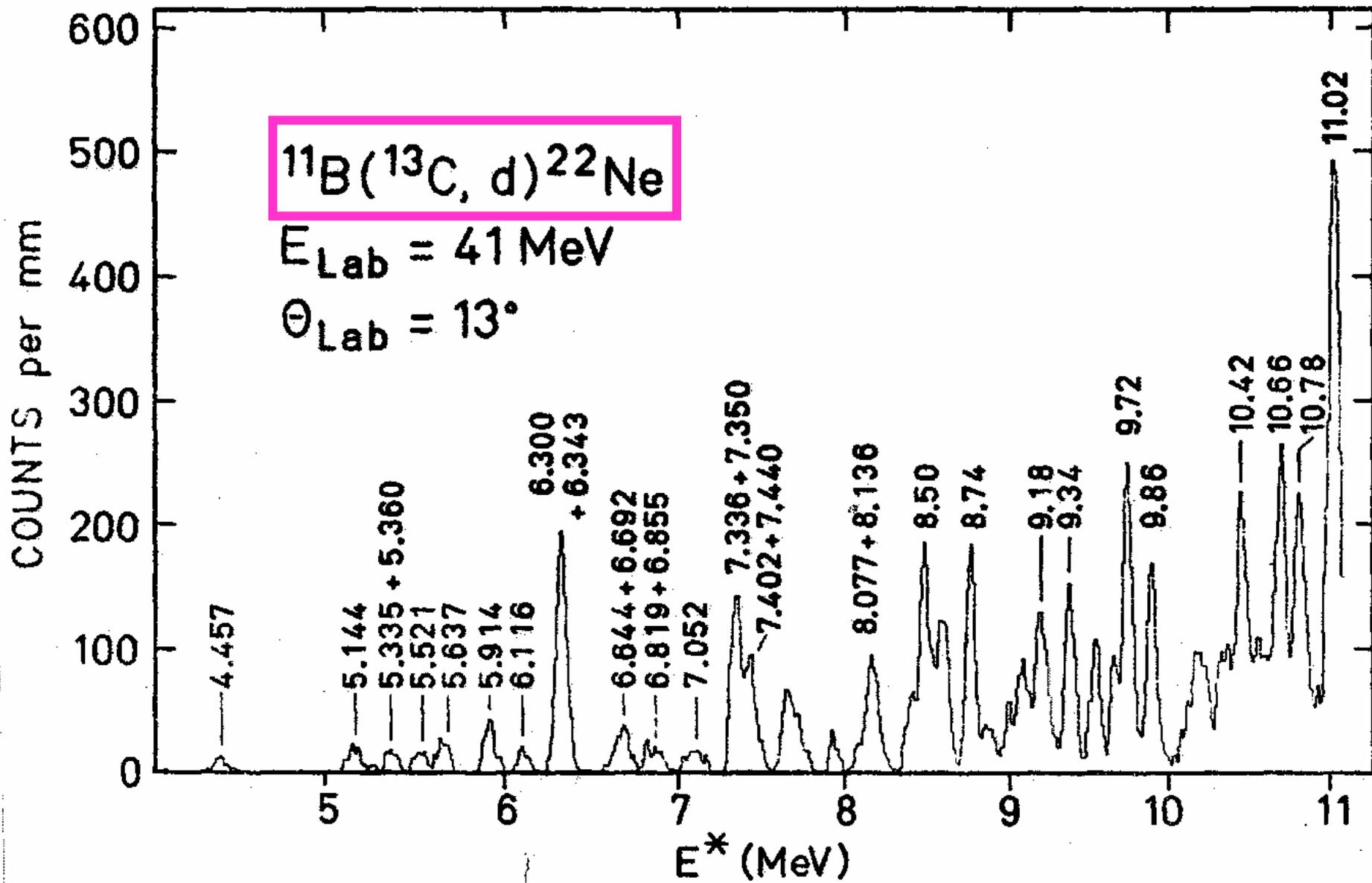


REAÇÃO VIA NÚCLEO COMPOSTO

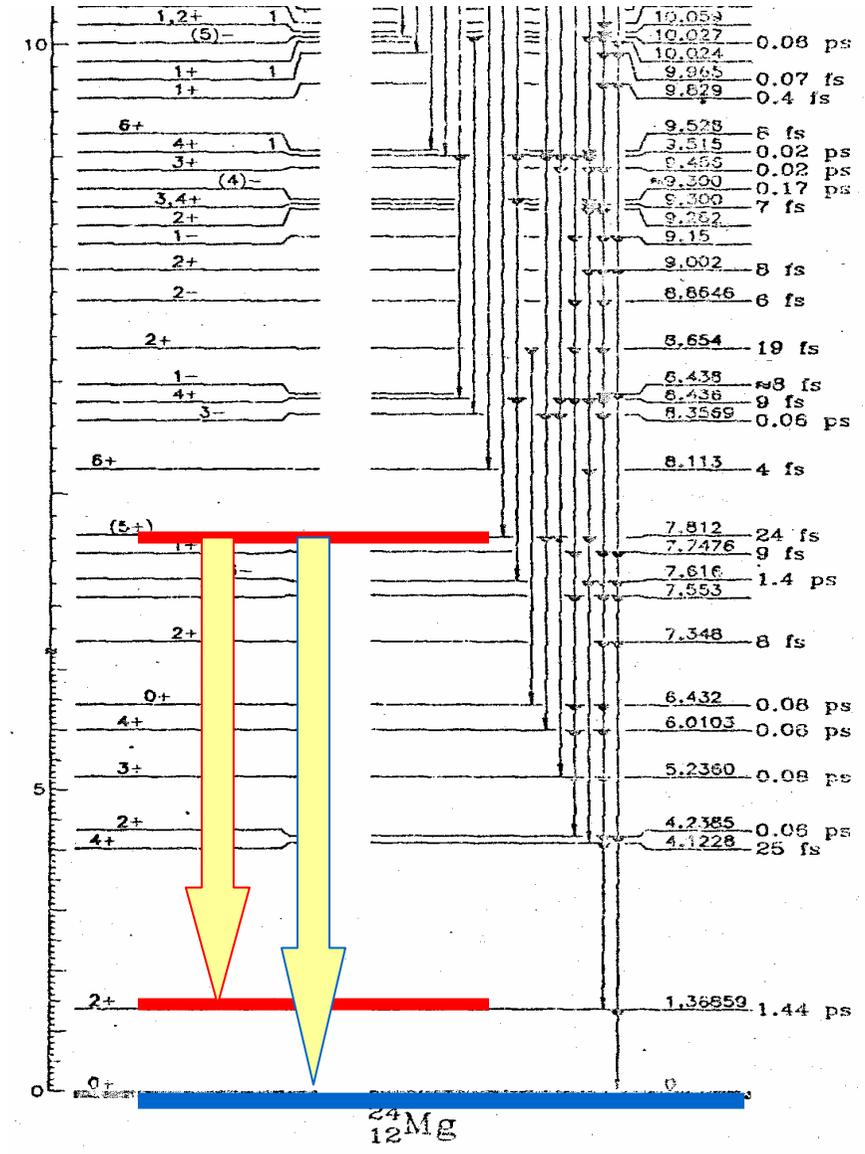


REAÇÕES VIA NÚCLEO COMPOSTO

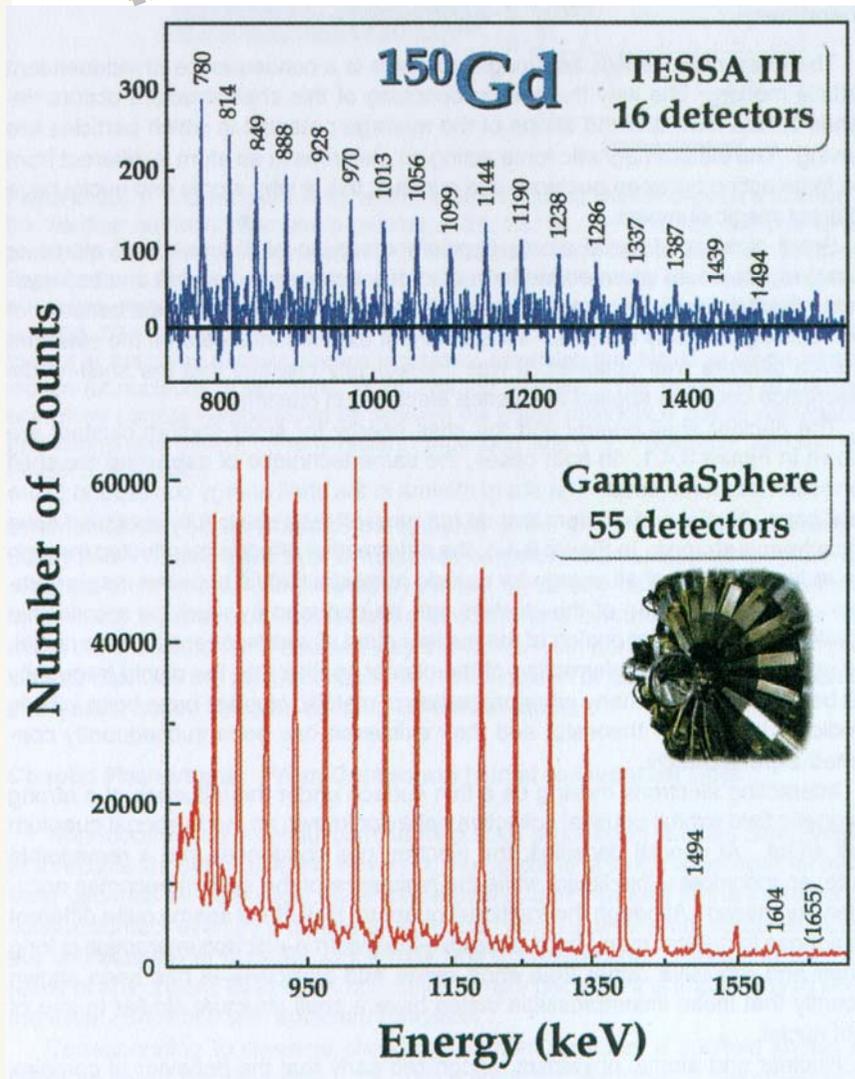
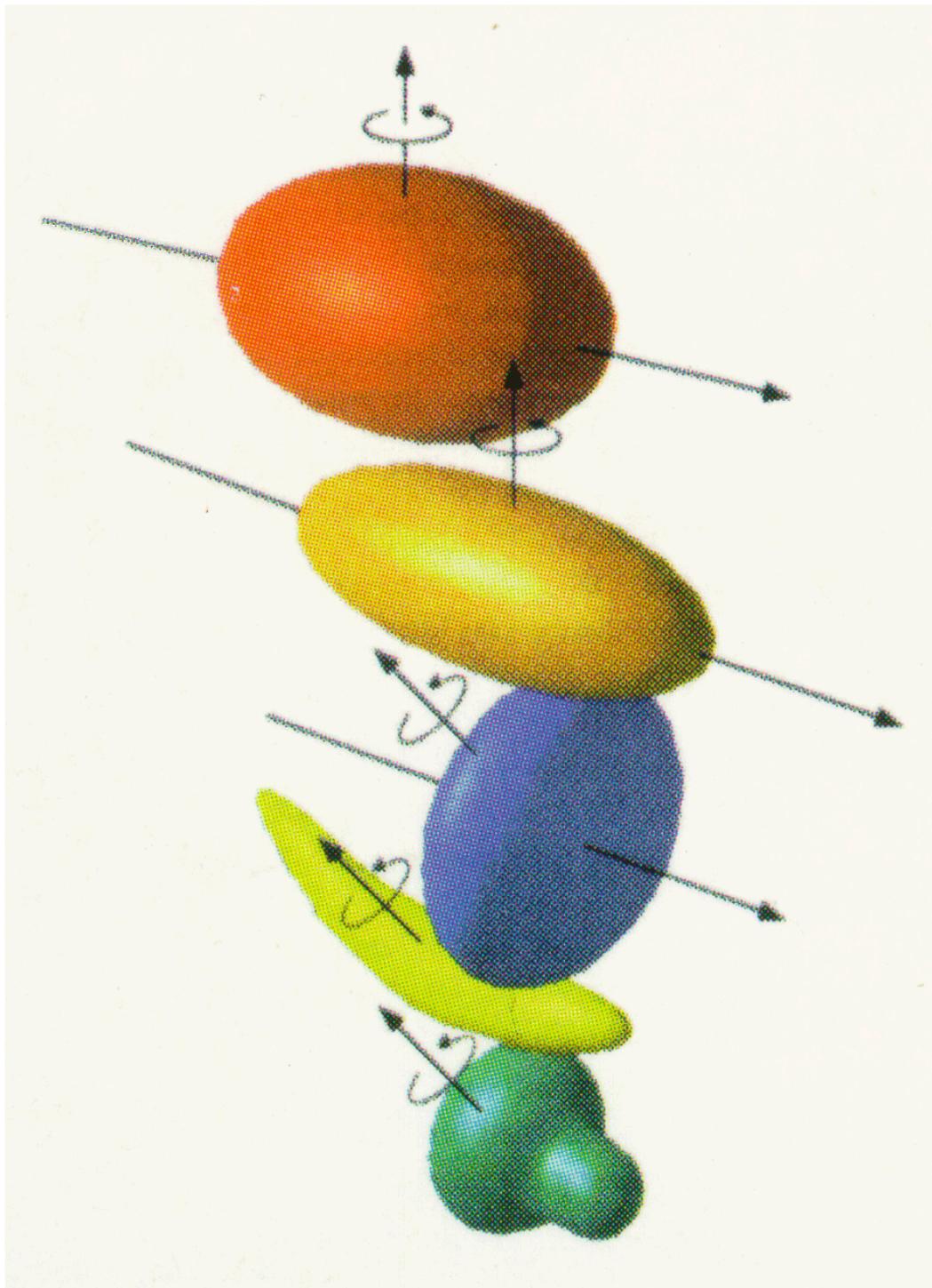


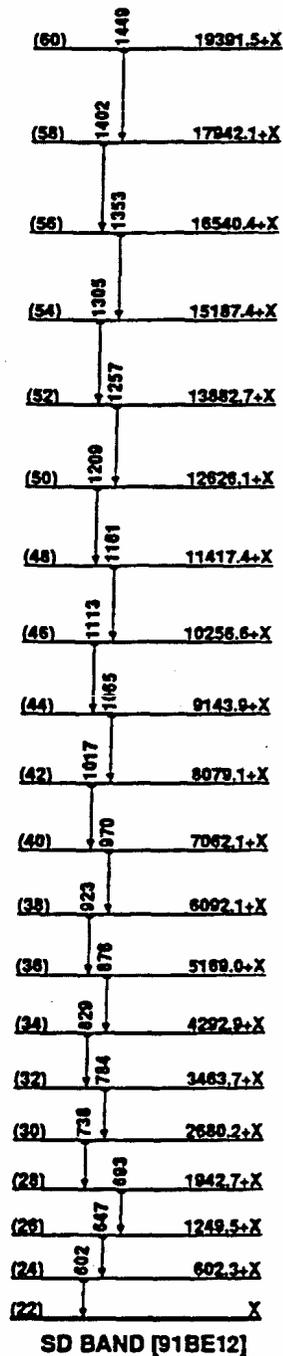


DECAIMENTO GAMMA (γ)



INCLUR DIAGRAMA DE NIVEIS





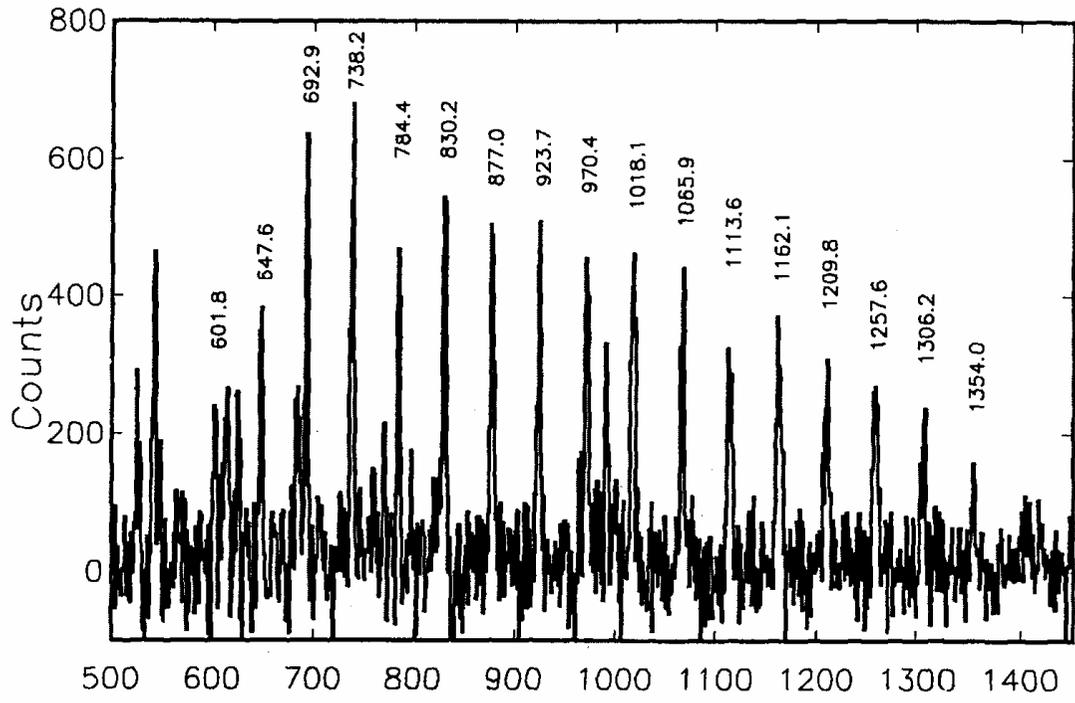
¹⁵²₆₆Dy Rotational Band

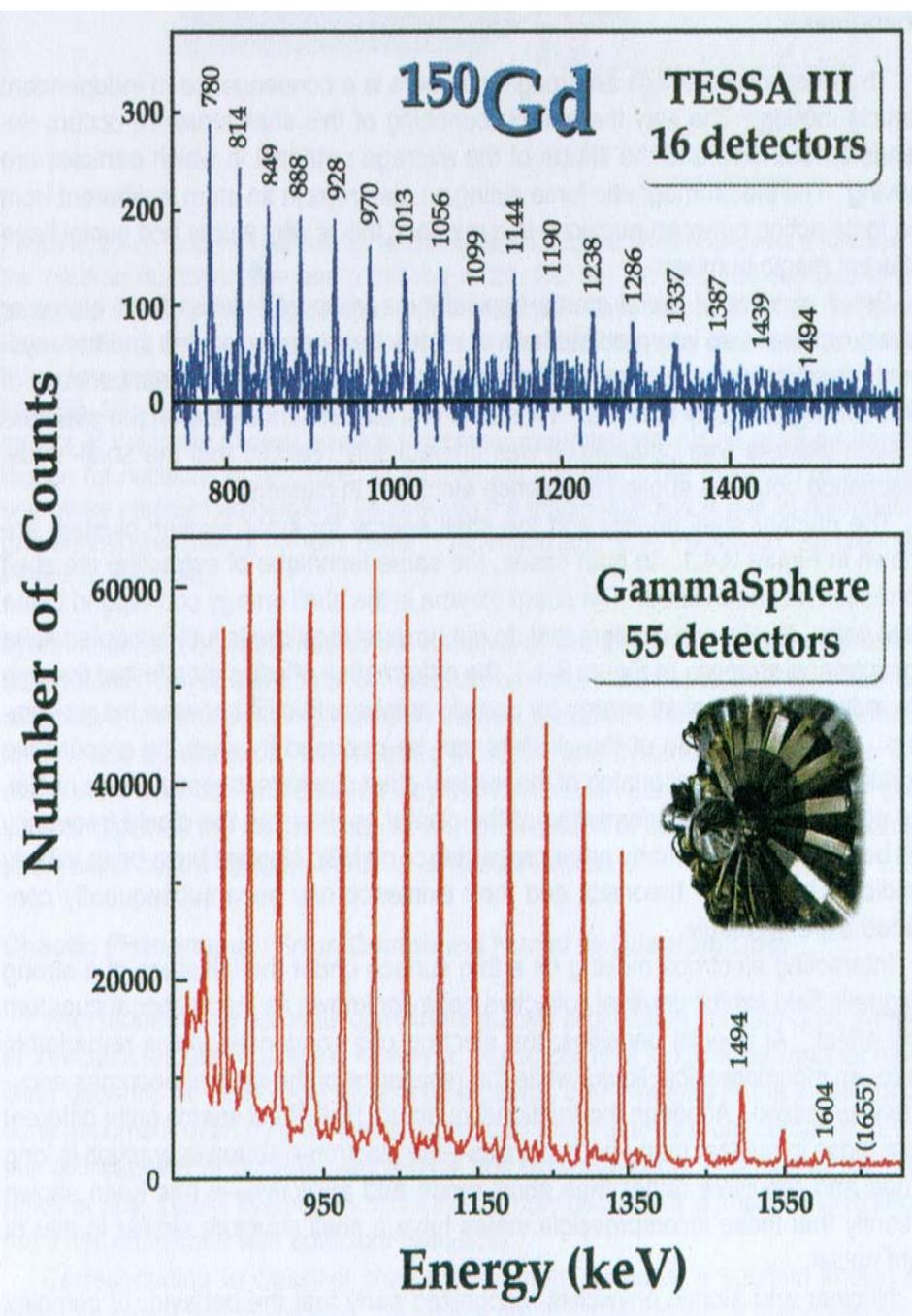
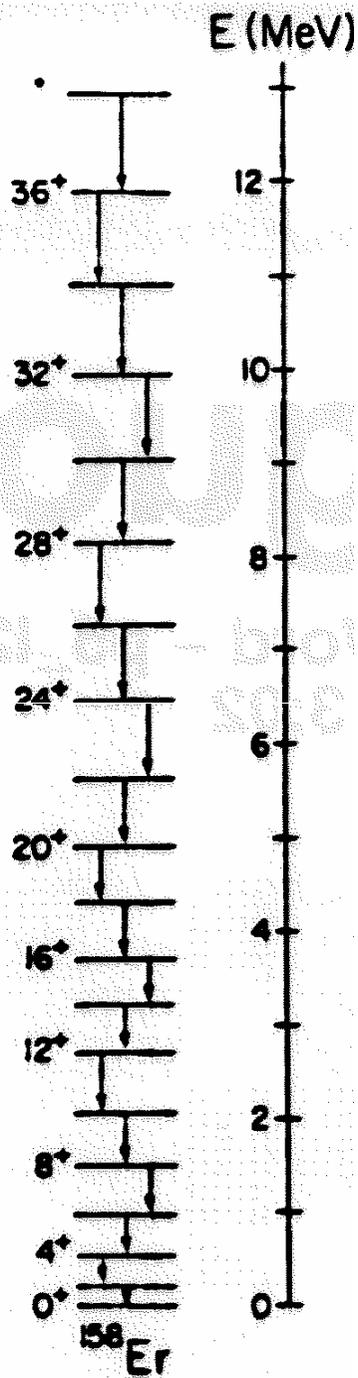
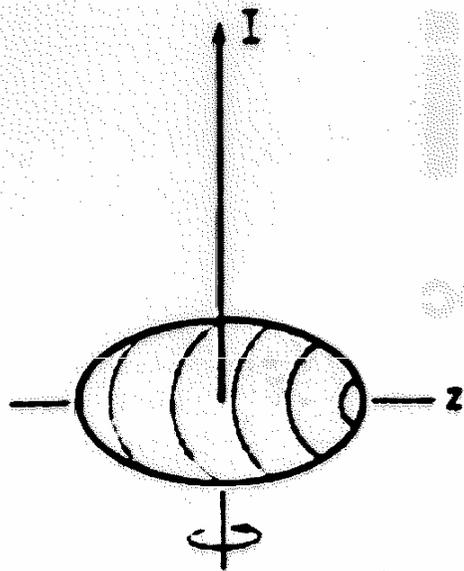
$$E_{\text{rot}}^{\gamma} = \frac{\hbar^2}{2\mathfrak{I}} [L(L+1)]$$

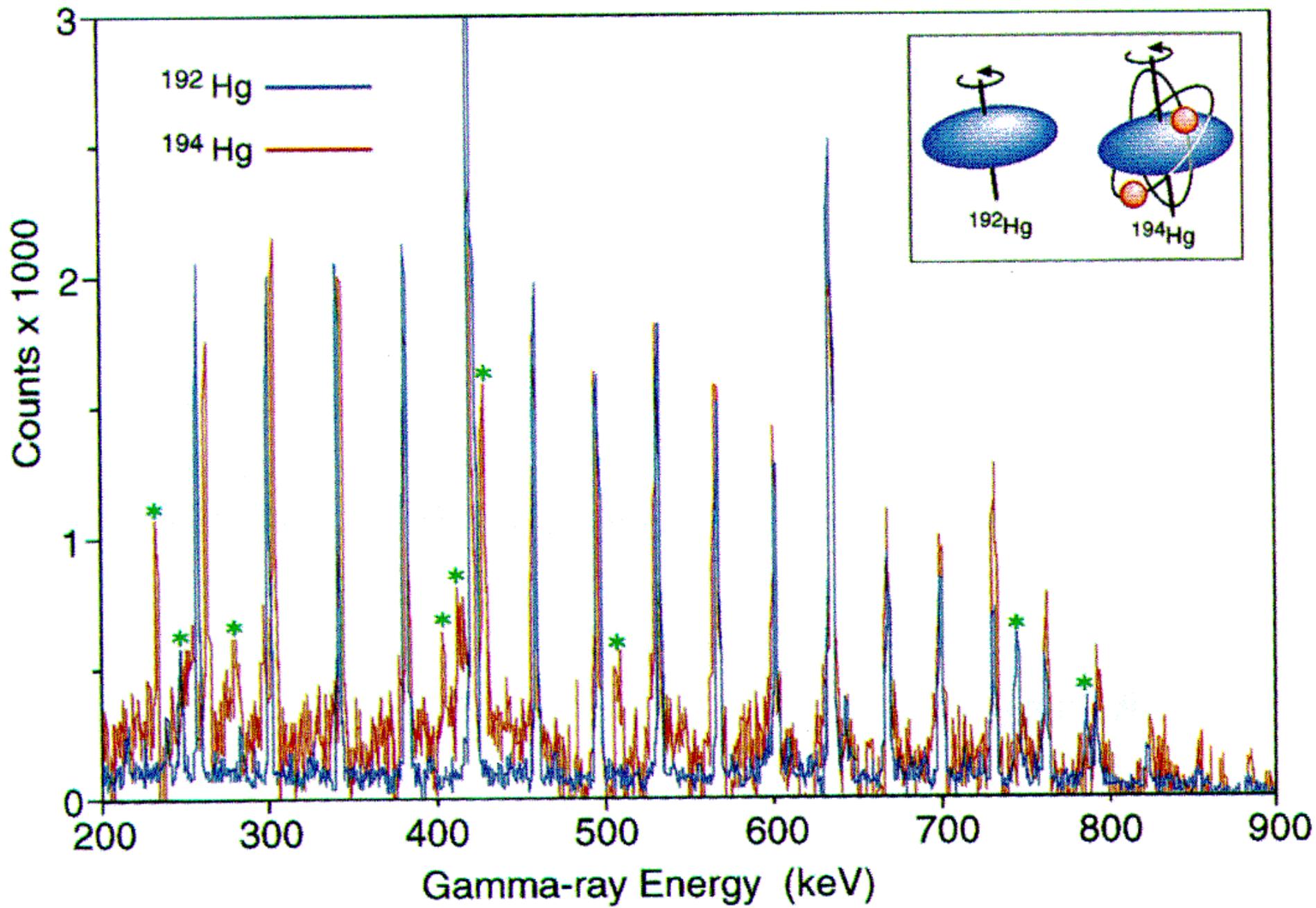
$$E_{(L+2) \rightarrow L}^{\gamma} = \frac{\hbar^2}{\mathfrak{I}} (2L+3) = E_{(L+2)}^{\gamma}$$

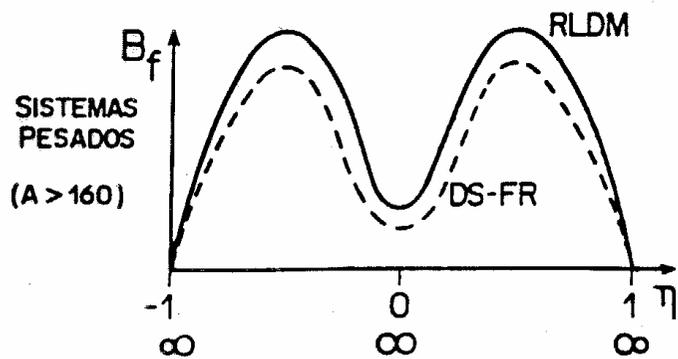
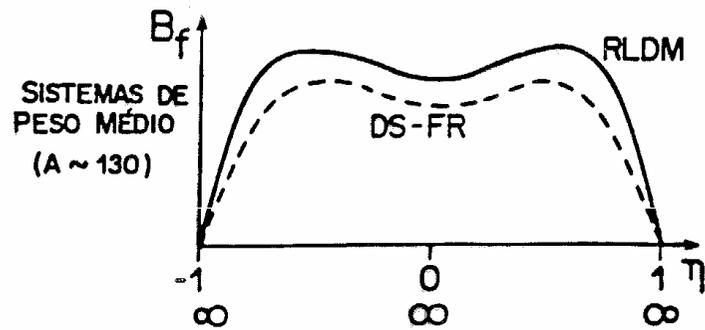
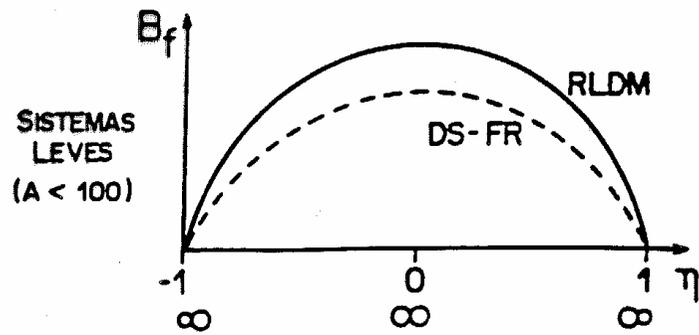
$$\Delta E^{\gamma} = E_{(L+2)}^{\gamma} - E_L^{\gamma} = 4 \frac{\hbar^2}{\mathfrak{I}}$$

152Dy SD Band

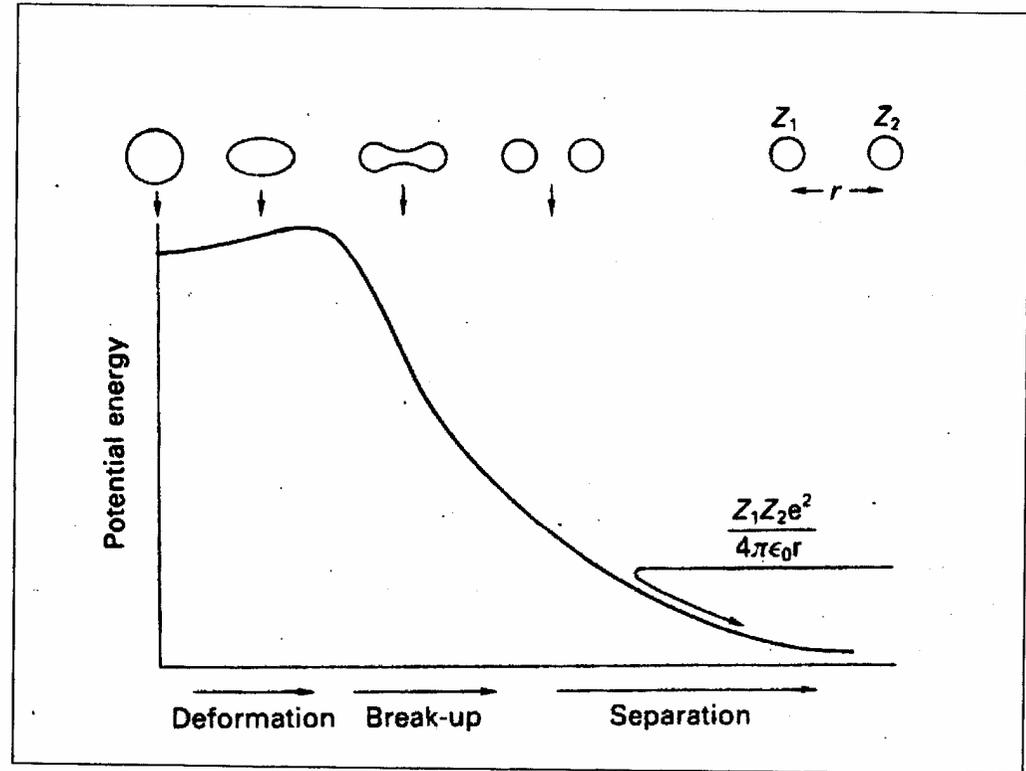


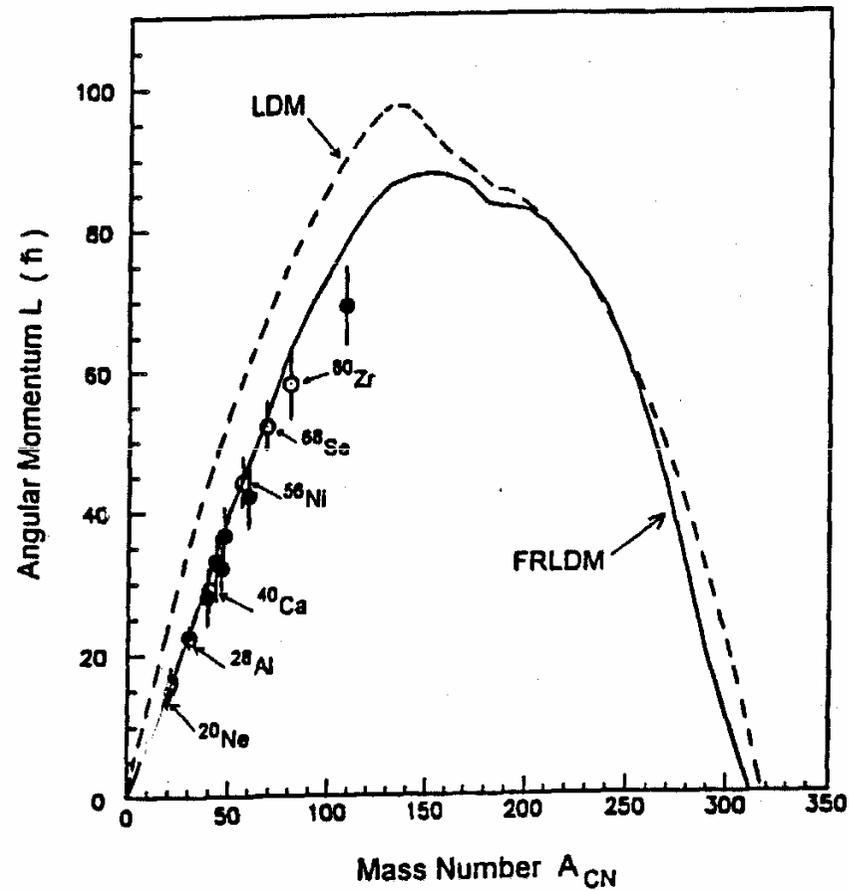
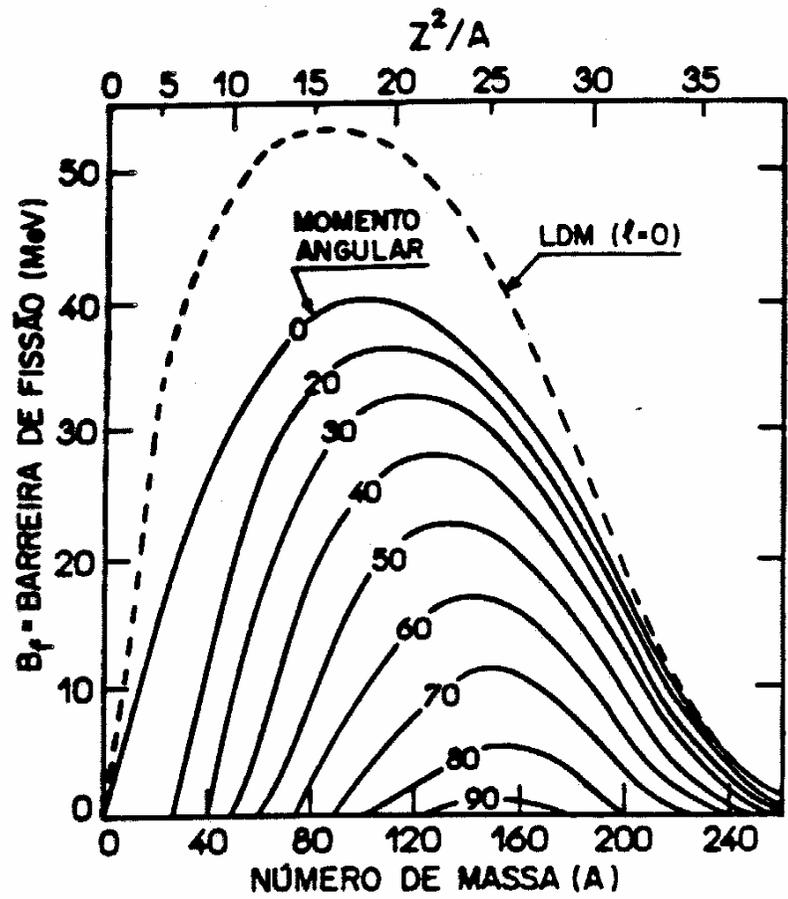




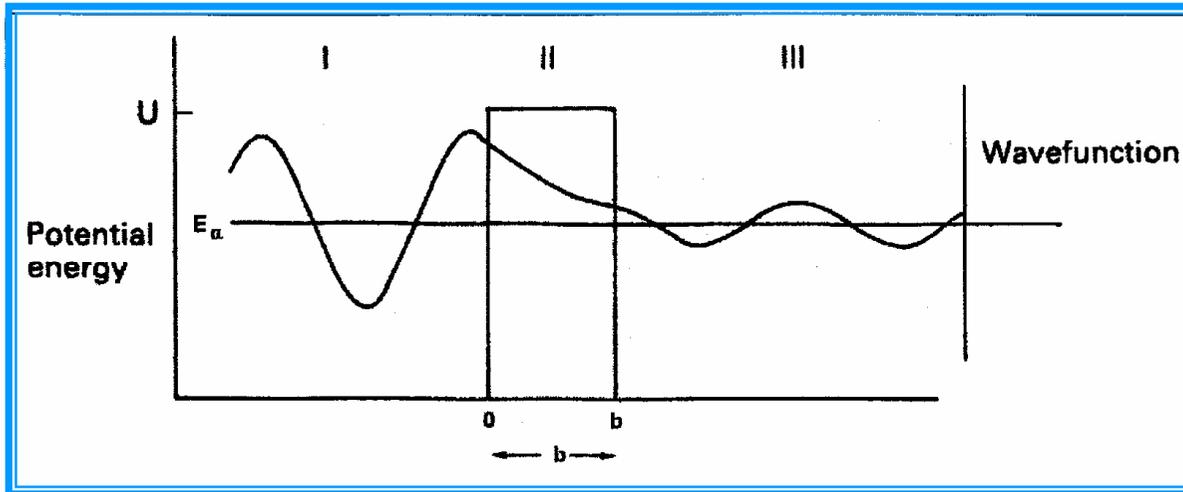


$$\eta = \frac{A_1 - A_2}{A_1 + A_2}$$

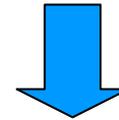




Barrier Potential, $E < V_0$



$$T = |F|^2 / |A|^2$$



$$\psi_1 = A e^{ik_1 x} + B e^{-ik_1 x}$$

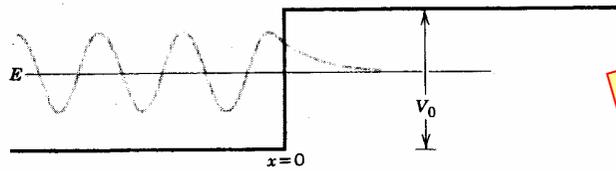
$$\psi_2 = C e^{k_2 x} + D e^{-k_2 x}$$

$$\psi_3 = F e^{ik_3 x} + \cancel{G} e^{-ik_3 x}$$

$$k^I = k^3 = \sqrt{2mE} / \hbar \quad \text{and} \quad k^S = \sqrt{2m(E - V_0)} / \hbar$$

$$T = \frac{1}{1 + \frac{1}{4} \frac{V_0^2}{E(E - V_0)} \sin^2 k_2 a}$$

$$k_2 = \sqrt{2m(V_0 - E)} / \hbar$$



ref. Krane

Figure 2.3 The wave function of a particle of energy E encountering a step of height V_0 , for the case $E < V_0$. The wave function decreases exponentially in the classically forbidden region, where the classical kinetic energy would be negative. At $x = 0$, ψ and $d\psi/dx$ are continuous.

the classically forbidden region. All (classical) particles are reflected at the boundary; the quantum mechanical wave packet, on the other hand, can penetrate a short distance into the forbidden region. The (classical) particle is never directly observed in that region; since $E < V_0$, the kinetic energy would be negative in region 2. The solution is illustrated in Figure 2.3

Barrier Potential, $E > V_0$

The potential is

$$V(x) = \begin{cases} 0 & x < 0 \\ V_0 & 0 \leq x \leq a \\ 0 & x > a \end{cases} \quad (2.35)$$

In the three regions 1, 2, and 3, the solutions are

$$\begin{aligned} \psi_1 &= A e^{ik_1x} + B e^{-ik_1x} \\ \psi_2 &= C e^{ik_2x} + D e^{-ik_2x} \\ \psi_3 &= F e^{ik_3x} + G e^{-ik_3x} \end{aligned} \quad (2.36)$$

where $k_1 = k_3 = \sqrt{2mE/\hbar^2}$ and $k_2 = \sqrt{2m(E - V_0)/\hbar^2}$.

Using the continuity conditions at $x = 0$ and at $x = a$, and assuming again that particles are incident from $x = -\infty$ (so that G can be set to zero), after considerable algebraic manipulation we can find the transmission coefficient $T = |F|^2/|A|^2$:

$$T = \frac{1}{1 + \frac{1}{4} \frac{V_0^2}{E(E - V_0)} \sin^2 k_2 a} \quad (2.37)$$

The solution is illustrated in Figure 2.4.

Barrier Potential, $E < V_0$

For this case, the ψ_1 and ψ_3 solutions are as above, but ψ_2 becomes

$$\psi_2 = C e^{k_2x} + D e^{-k_2x} \quad (2.38)$$

where now $k_2 = \sqrt{2m(V_0 - E)/\hbar^2}$. Because region 2 extends only from $x = 0$

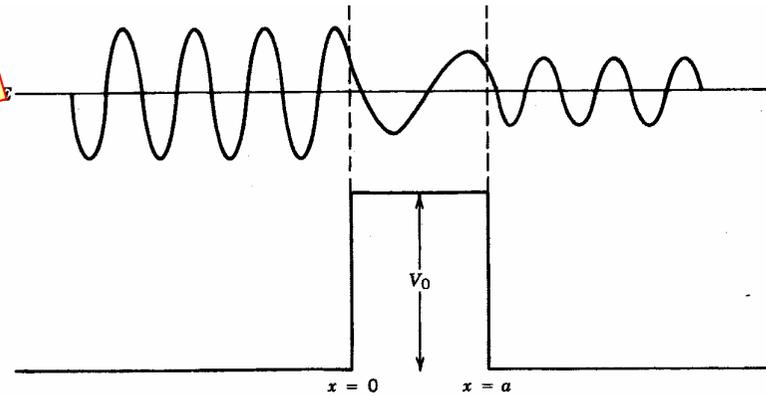


Figure 2.4 The wave function of a particle of energy $E > V_0$ encountering a barrier potential. The particle is incident from the left. The wave undergoes reflections at both boundaries, and the transmitted wave emerges with smaller amplitude.

to $x = a$, the question of an exponential solution going to infinity does not arise, so we cannot set C or D to zero.

Again, applying the boundary conditions at $x = 0$ and $x = a$ permits the solution for the transmission coefficient:

$$T = \frac{1}{1 + \frac{1}{4} \frac{V_0^2}{E(V_0 - E)} \sinh^2 k_2 a} \quad (2.39)$$

Classically, we would expect $T = 0$ —the particle is not permitted to enter the forbidden region where it would have negative kinetic energy. The quantum wave can penetrate the barrier and give a nonzero probability to find the particle beyond the barrier. The solution is illustrated in Figure 2.5.

This phenomenon of *barrier penetration* or quantum mechanical *tunneling* has important applications in nuclear physics, especially in the theory of α decay, which we discuss in Chapter 8.

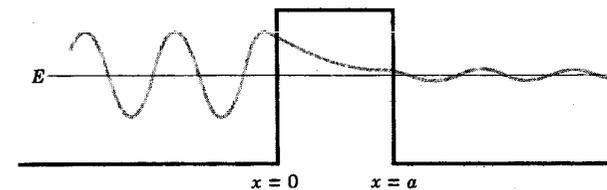
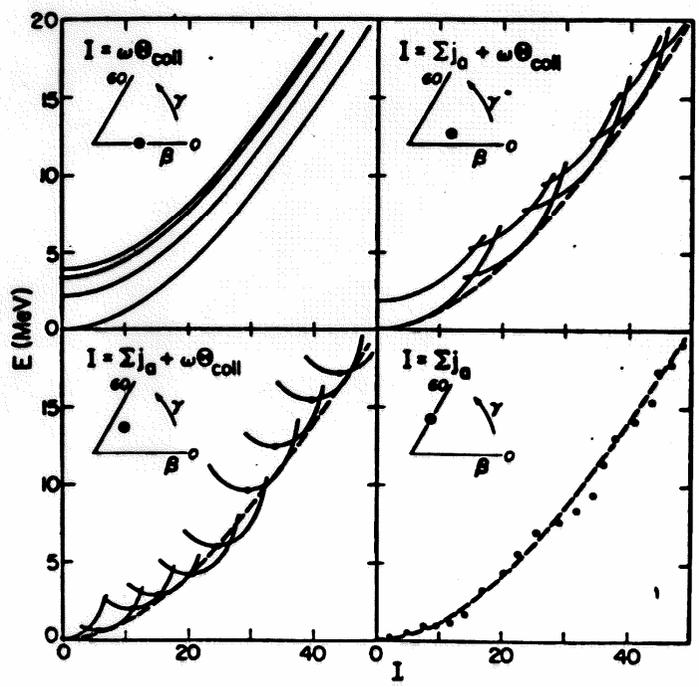
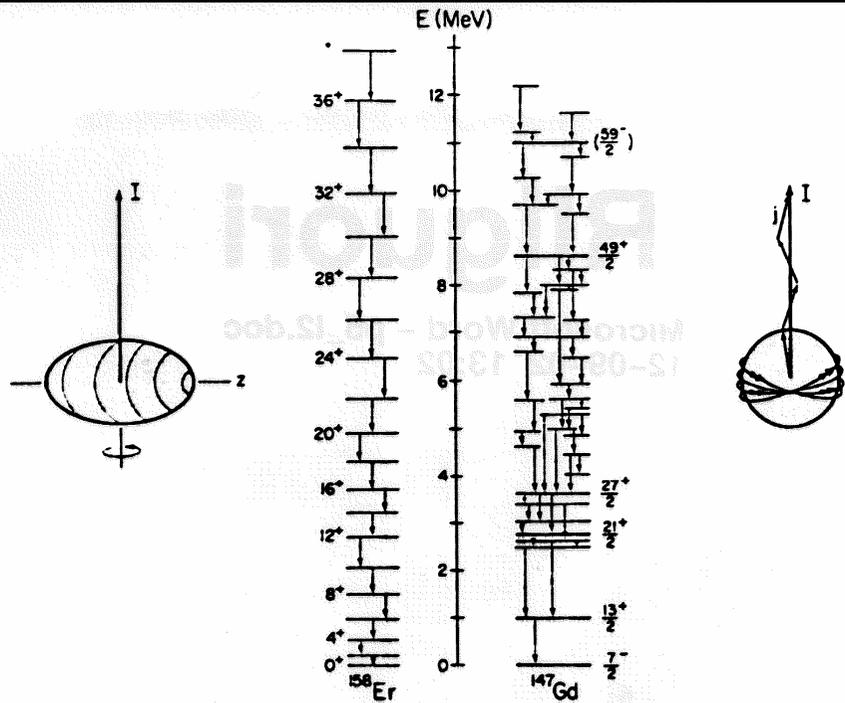


Figure 2.5 The wave function of a particle of energy $E < V_0$ encountering a barrier potential (the particle would be incident from the left in the figure). The wavelength is the same on both sides of the barrier, but the amplitude beyond the barrier is much less than the original amplitude. The particle can never be observed, inside the barrier (where it would have negative kinetic energy) but it can be observed beyond the barrier.

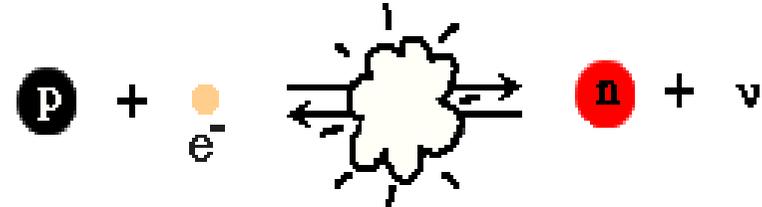


proton/neutron conversions

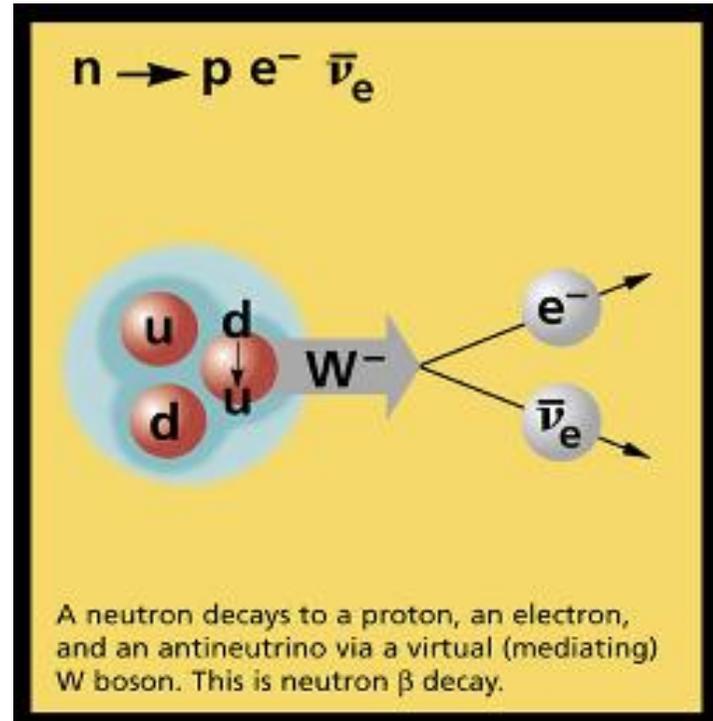
Reaction #1:



Reaction #2:



(The double arrows indicate these reactions go both ways.)



BOSONS			FERMIONS			matter constituents		
Unified Electroweak spin = 1			Leptons spin = 1/2			Quarks spin = 1/2		
Name	Mass GeV/c ²	Electric charge	Flavor	Mass GeV/c ²	Electric charge	Flavor	Approx. Mass GeV/c ²	Electric charge
γ photon	0	0	e ⁻ electron neutrino	<1×10 ⁻⁶	0	u up	0.003	2/3
W ⁻	80.4	-1	e ⁻ electron	0.000511	-1	d down	0.006	-1/3
W ⁺	80.4	+1	μ ⁻ muon neutrino	<0.0002	0	c charm	1.3	2/3
Z ⁰	91.187	0	μ ⁻ muon	0.106	-1	s strange	0.1	-1/3
			τ ⁻ tau neutrino	<0.02	0	t top	175	2/3
			τ ⁻ tau	1.7771	-1	b bottom	4.3	-1/3